



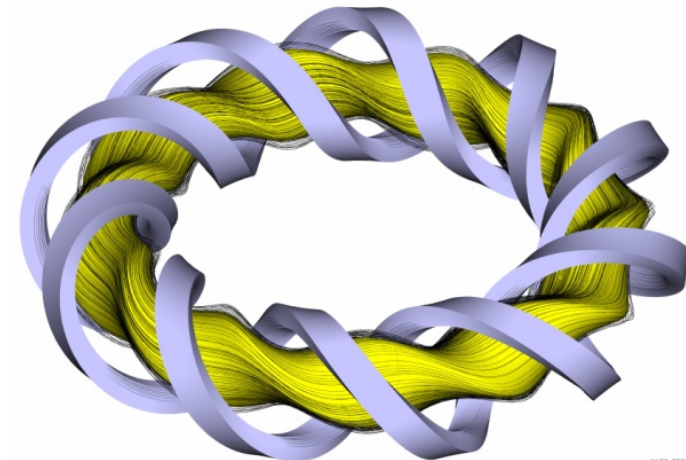
Topological Changes in Magnetic Flux Surfaces during IDB-SDC Discharge in LHD

T. Morisaki, Y. Suzuki, M. Goto, J. Miyazawa, S. Masuzaki, R. Sakamoto, Y. Narushima, M. Kobayashi, N. Ohyaabu, H. Yamada, A. Komori, O. Motojima and LHD Experimental Group

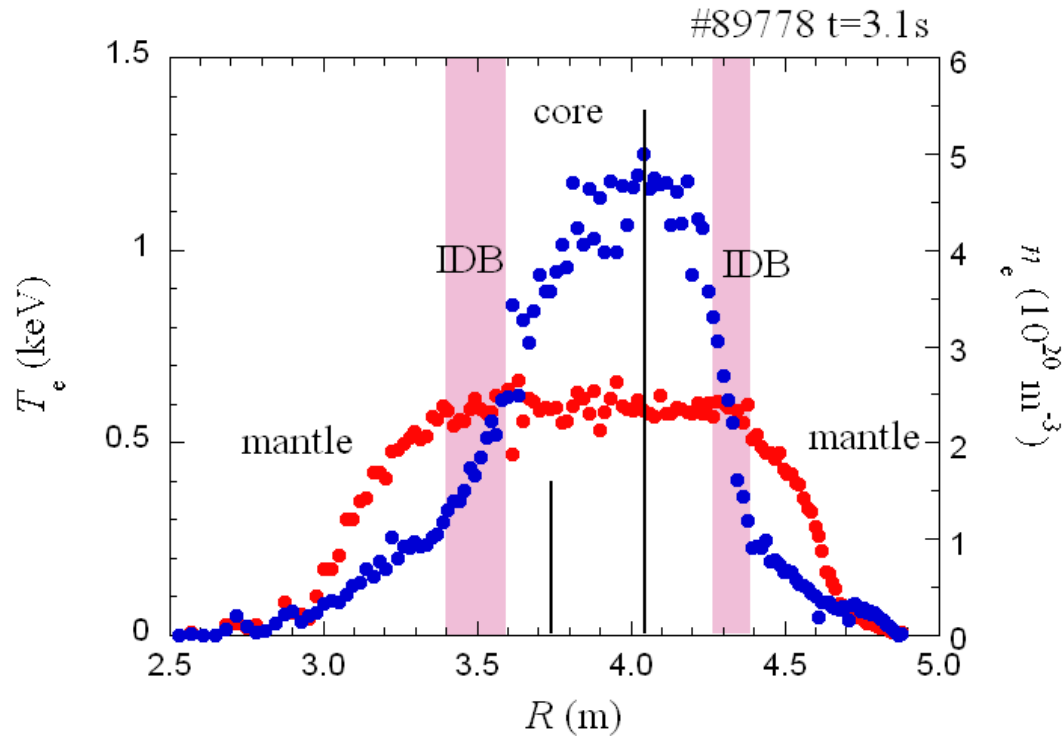
National Institute for Fusion Science

Contents

- 1) Motivation and background
- 2) Introduction to IDB-SDC discharge
- 3) Experimental numerical setup
- 4) Results
- 5) Summary



Motivation and background



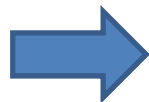
SuperDense Core plasma with Internal Diffusion Barrier (IDB-SDC plasma)

- * $n_e(0) \sim 1 \times 10^{21} \text{ m}^{-3}$
- * $T_e(0) \sim 0.4 \text{ keV}$
- * $P(0) \sim 140 \text{ kPa}$
- * $W_{\text{dia}} \sim 1 \text{ MJ}$
- * $\beta(0) \sim 4.7\%$
- * $n\tau T \sim 4.4 \times 10^{19} \text{ keVsm}^{-3}$

high pressure core region ==> - large Shafranov shift
 steep gradient at IDB ==> - spontaneous current , e.g. bootstrap
 - destabilization

on the other hand....

low pressure and low-gradient mantle



what happened to magnetic topology ?

Introduction to IDB-SDC plasma

Scenario to achieve IDB-SDC plasma

Edge plasma control with divertor and wall conditioning

- impurity control
- recycling control (profile control)

Central fuelling by pellet injection

low n_0
low edge n_e
high edge T_e
high T_e at pedestal



superdense
core

IDB-SDC discharges with highly dense core and low mantle plasmas

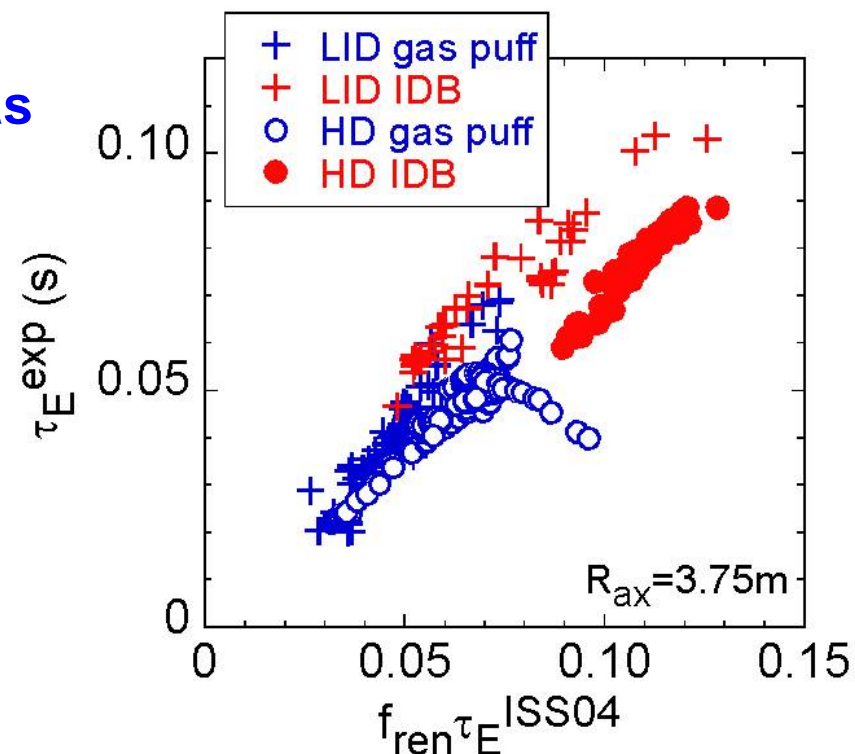
- reduce radiation loss in edge region
==> evade radiation collapse
- realize deep fuelling
- realize central deposition of heating power
- bring about steep T_e gradient



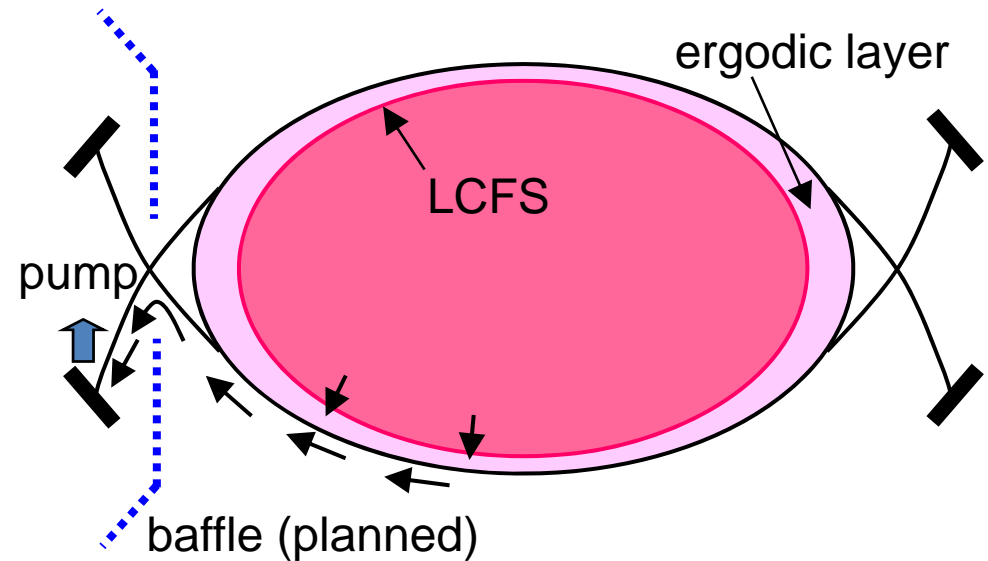
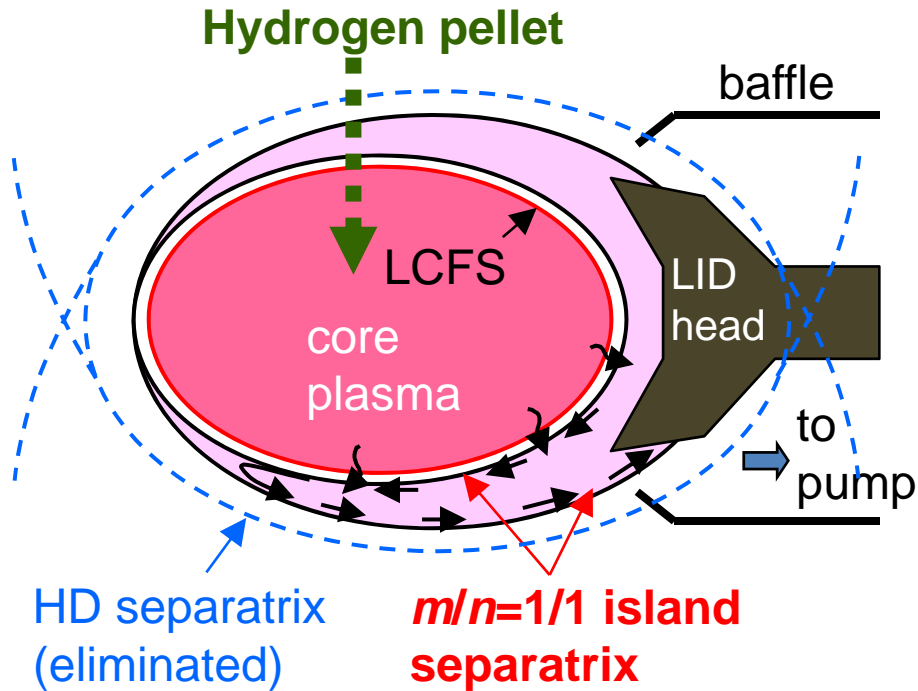
confinement improvement

keeping positive n_e dependence

$$\tau_E^{ISS04} = 0.134 a^{2.28} R^{0.64} P^{-0.61} \bar{n}_e^{-0.54} B^{0.84} t_{2/3}^{0.41}$$



Two divertor concepts in LHD



Local Island Divertor (LID)

- * utilize $m/n=1/1$ island
- * insert a divertor head **locally**
 - ==> recycling region is localized
 - ==> full closed configuration available
- * no ergodic layer.

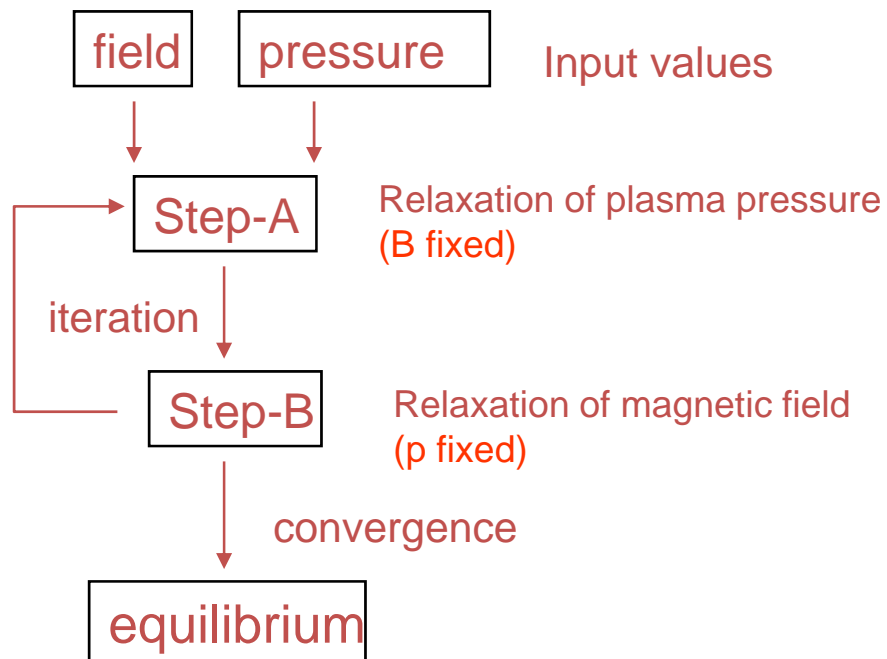
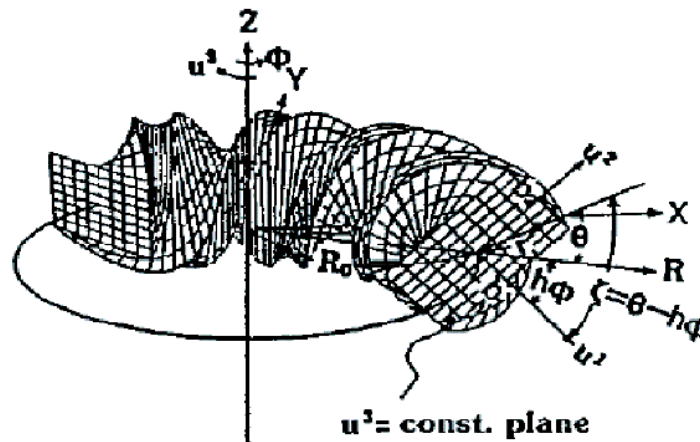
Helical Divertor (HD)

- * Intrinsic double-null type divertor
- * with thick ergodic layer
- * long legs.
- * short field lines (2-3m in LHD)
- * nonaxisymmetric

HINT2 code

3D MHD equilibrium calculation code **without assumption of nested flux surfaces**

- relaxation method (initial value problem)
- Eulerian coordinate(rotating helical coordinate



$$p^{i+1} = \bar{p} = \frac{\int_{-L_{in}}^{L_{in}} \mathcal{F} p^i \frac{dl}{B}}{\int_{-L_{in}}^{L_{in}} \frac{dl}{B}}, \quad \mathcal{F} = \begin{cases} 1 & : \text{for } L_C \geq L_{in} \\ 0 & : \text{for } L_C < L_{in} \end{cases}$$

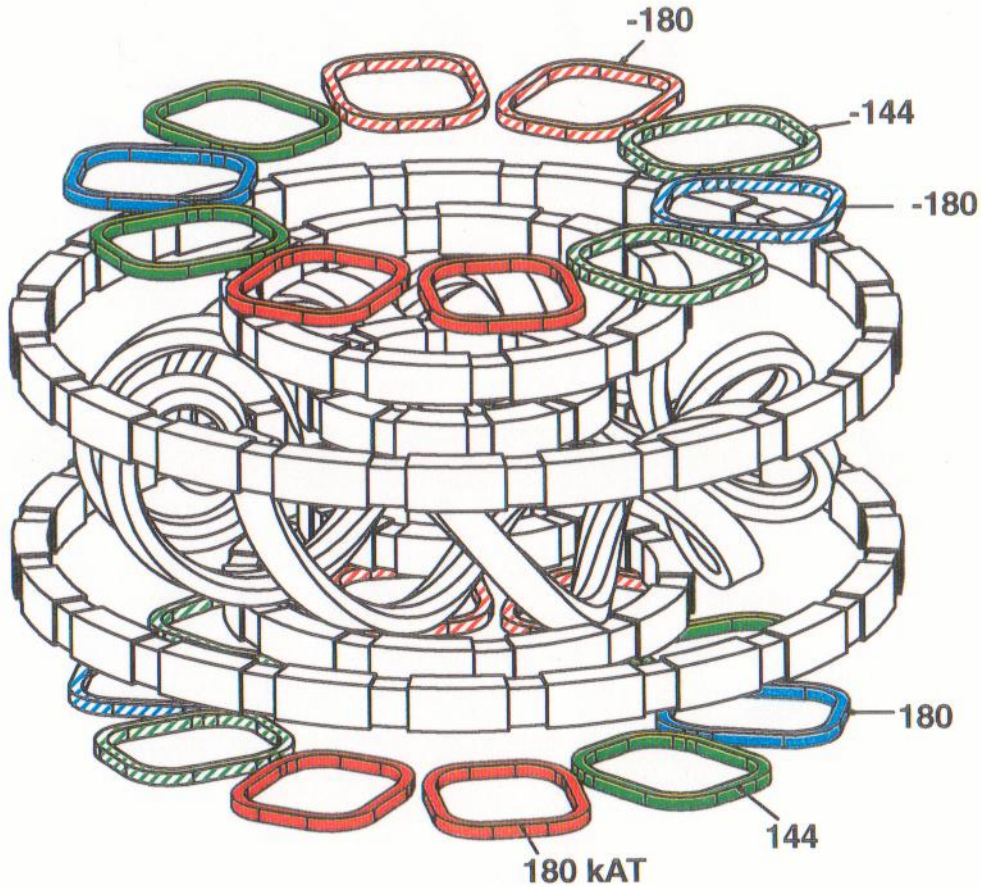
$$\frac{\partial \mathbf{v}}{\partial t} = -f_C [\nabla p - (\mathbf{j} - \mathbf{j}_0) \times \mathbf{B}]$$

$$\frac{\partial \mathbf{B}}{\partial t} = \nabla \times [\mathbf{v} \times \mathbf{B} - \eta(\mathbf{j} - \mathbf{j}_0 - \mathbf{j}_{net})]$$

$$\mathbf{j} = \nabla \times \mathbf{B}$$

$$f_C = \begin{cases} 1 & : (B \leq B_C) \\ (B_C/B)^2 & : (B > B_C) \end{cases}$$

Perturbation coils

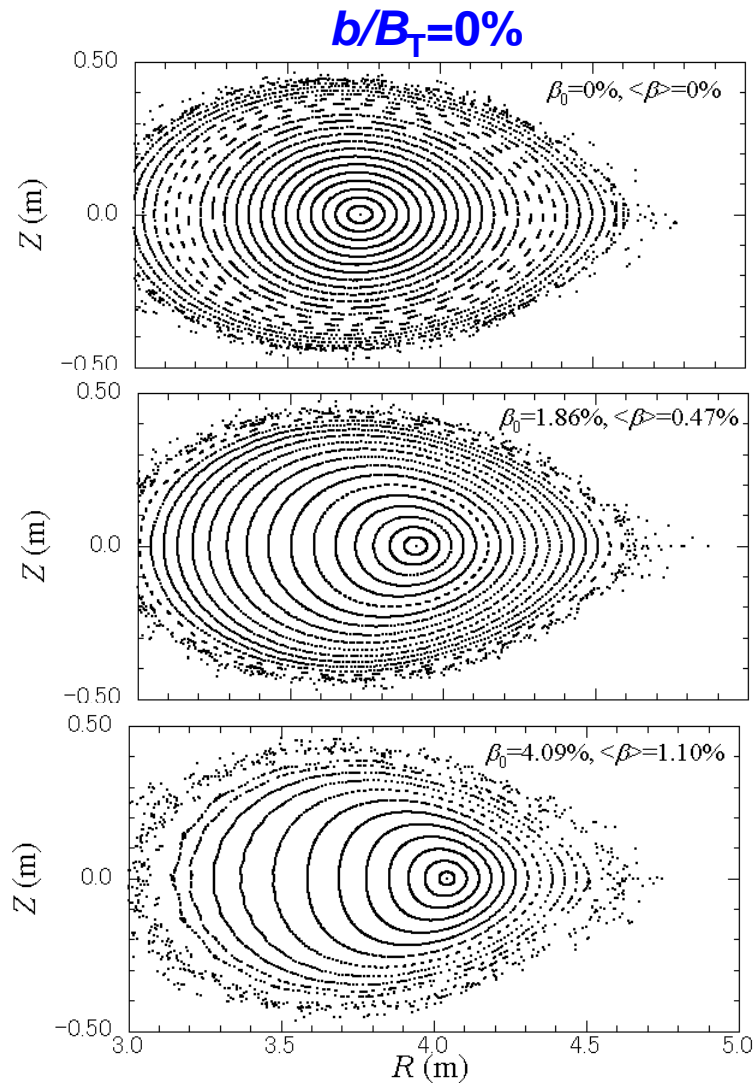


10 pairs of normal conductor coils

DC power supplies

	Steady state operation		Pulse operation	
	Current (kA)	Voltage (V)	Current (kA)	Voltage (V)
1	1.92	240	3.84	350
2	1.92	460	3.84	680
3	1.92	460	3.84	680

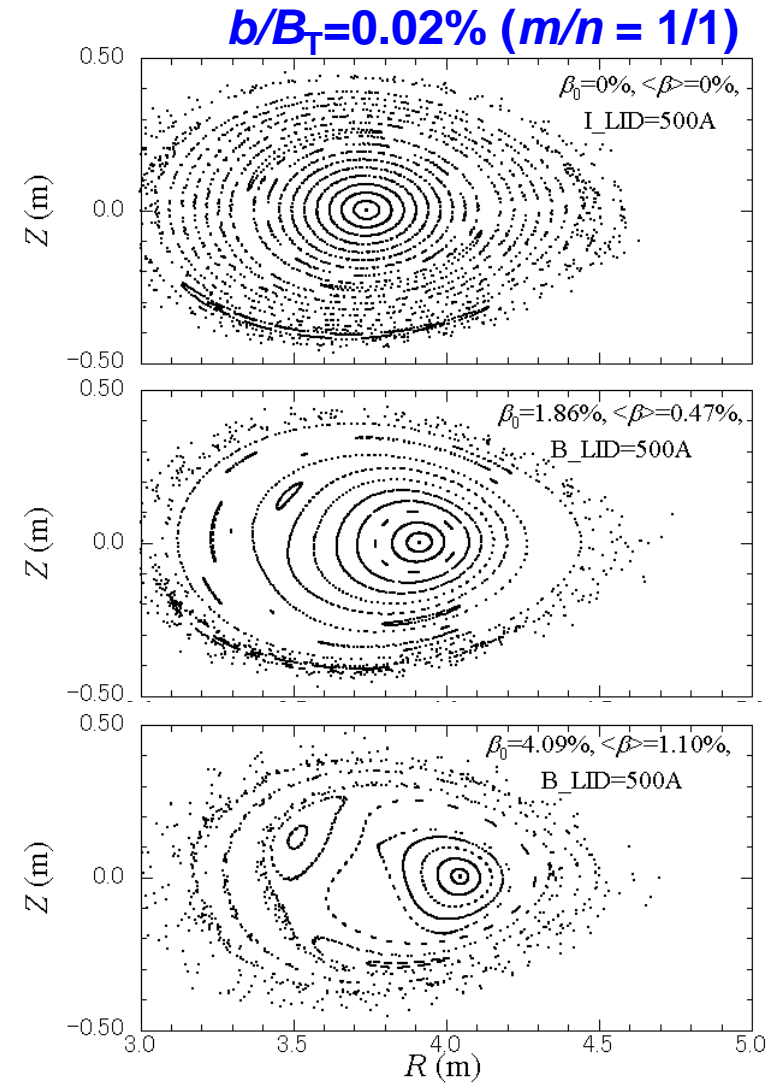
As β_0 increases (1)



$\beta_0=0\%$

$\beta_0=1.86\%$

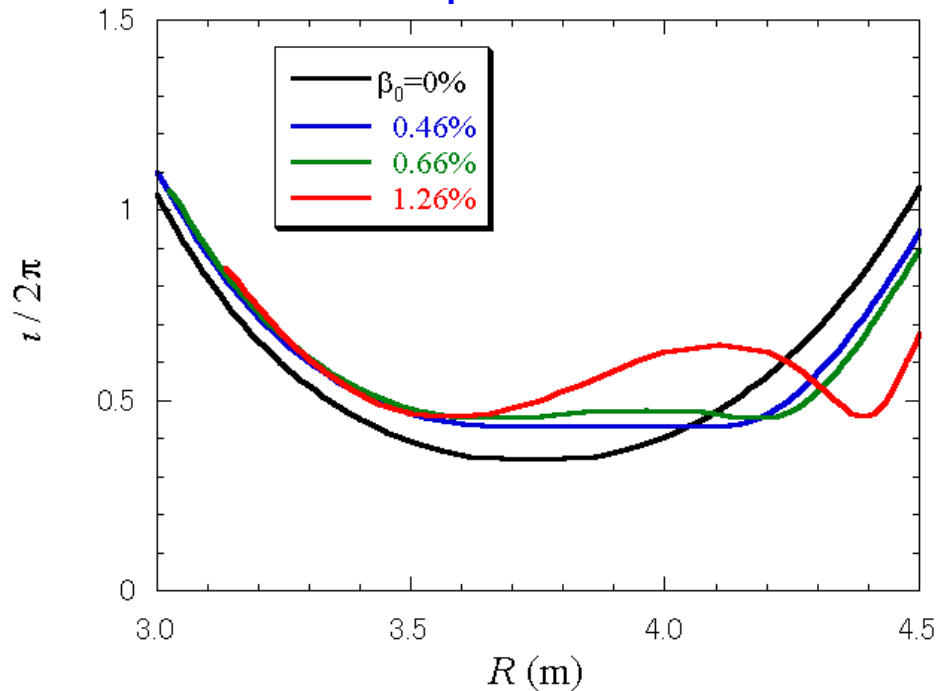
$\beta_0=4.09\%$



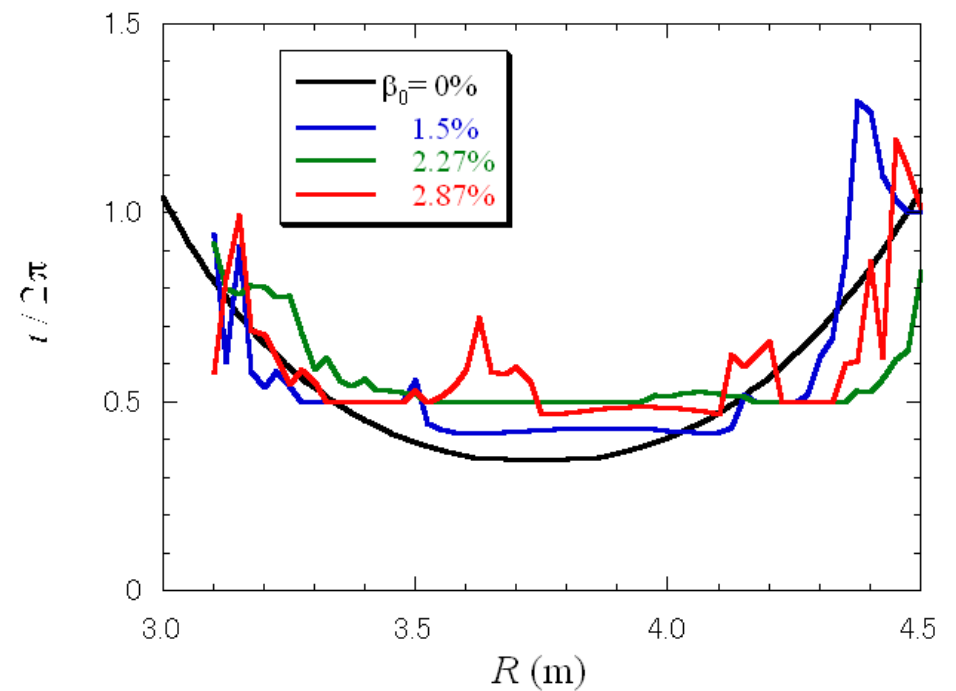
Flux surfaces are sensitive to external perturbation,
especially in high beta regime.

As β_0 increases (2)

$b/B_T=0\%$



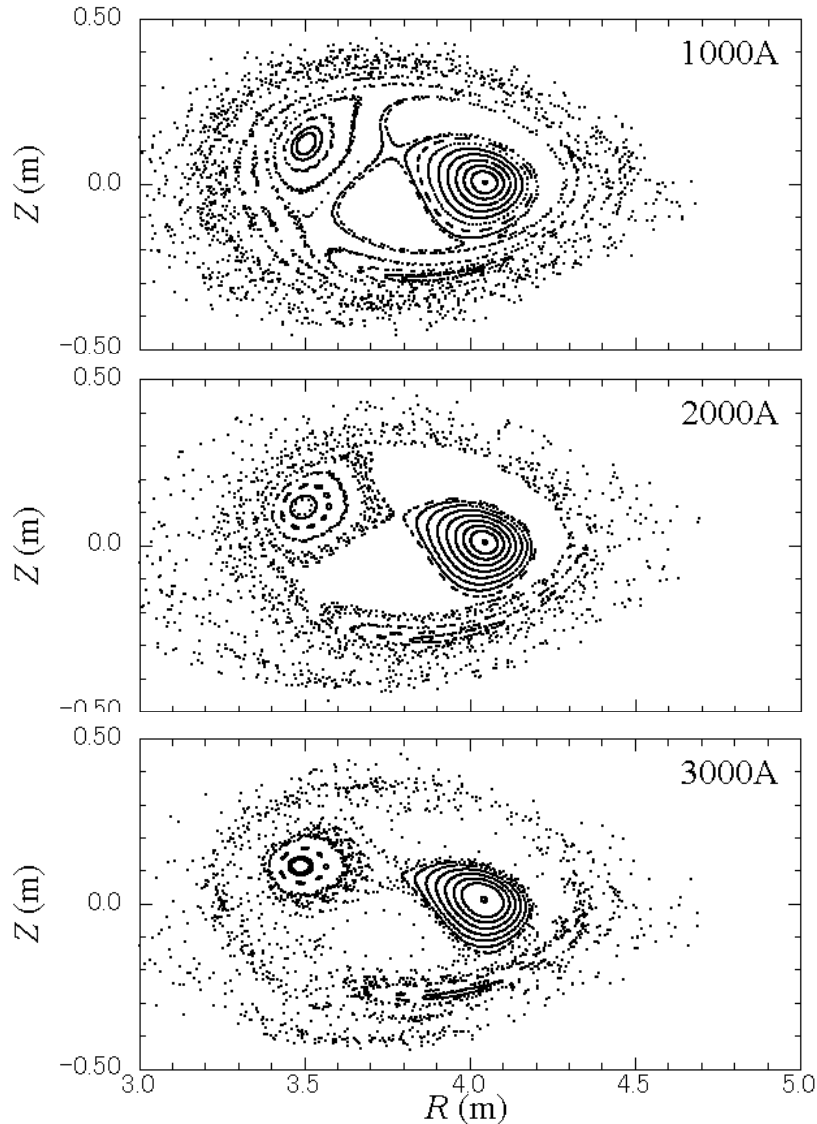
$b/B_T=0.12\%$ ($m/n = 1/1$)



- rotational transform increases with increase in beta
- shear near the center becomes low

Especially with perturbation, $m/n=2/1$ region spreads
==> wide 2/1 island grows at the core region.

As b/B_T increases



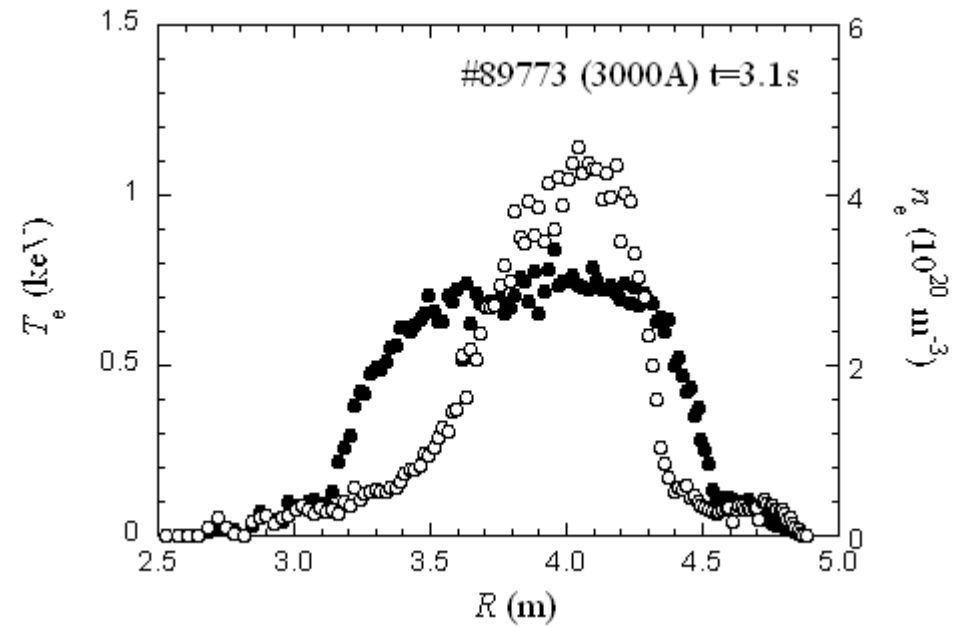
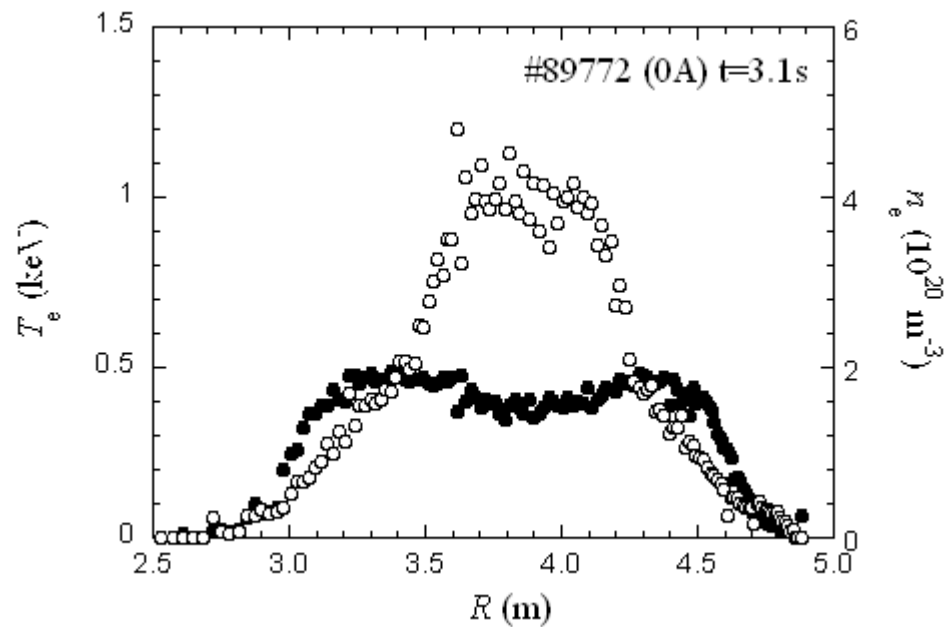
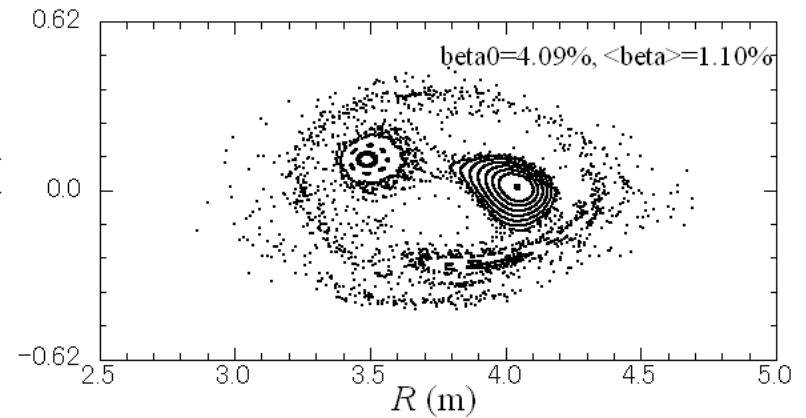
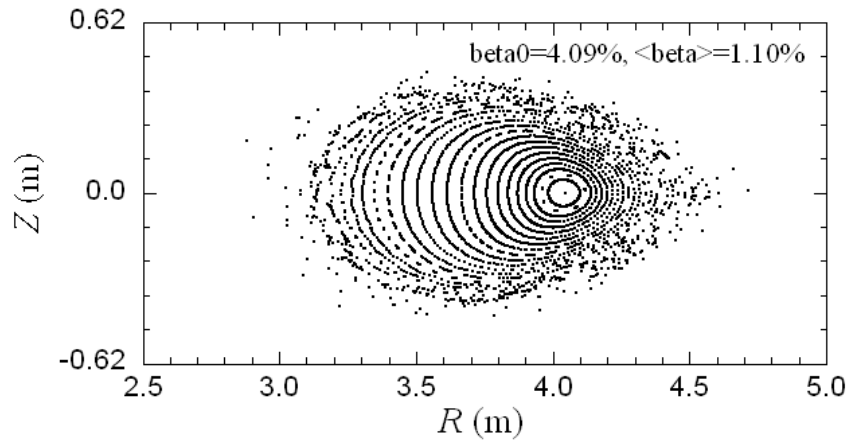
$b/B_T=0.04\%$ ($m/n = 1/1$)

$b/B_T=0.08\%$ ($m/n = 1/1$)

$b/B_T=0.12\%$ ($m/n = 1/1$)

Even a small perturbation, magnetic surfaces are modified or ergodized.

Experiment



n_e profile reflects magnetic structure, but T_e does not follow it

Kolmogorov length L_K

$$d(l) = d_0 \exp\left(\frac{l}{L_K}\right)$$

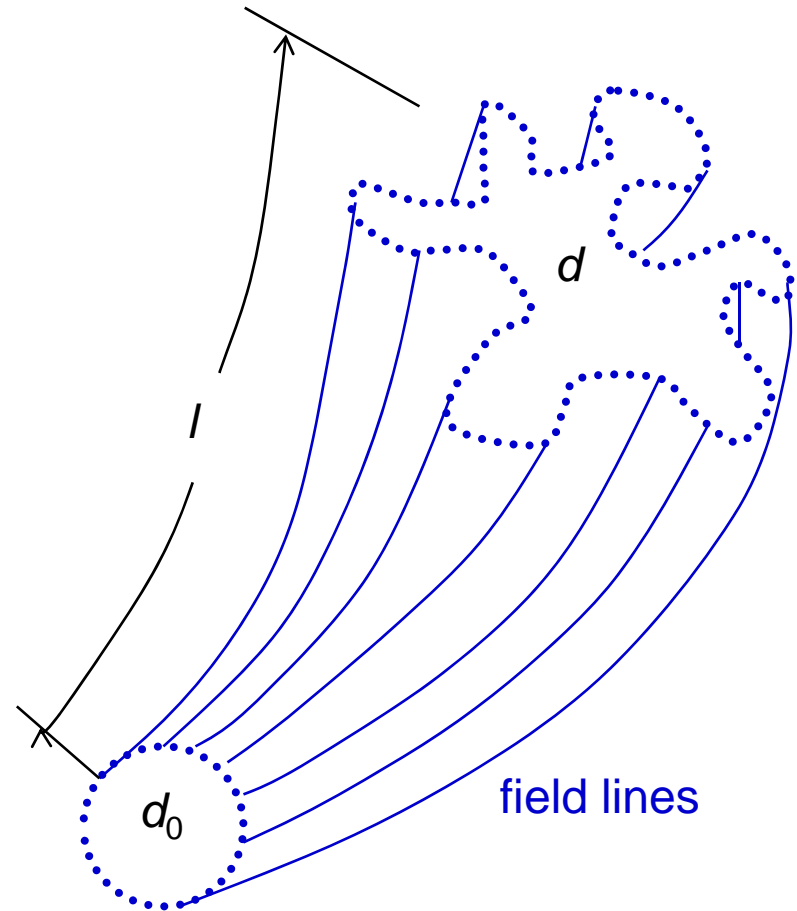
d : circumference

d_0 : initial value of circumference

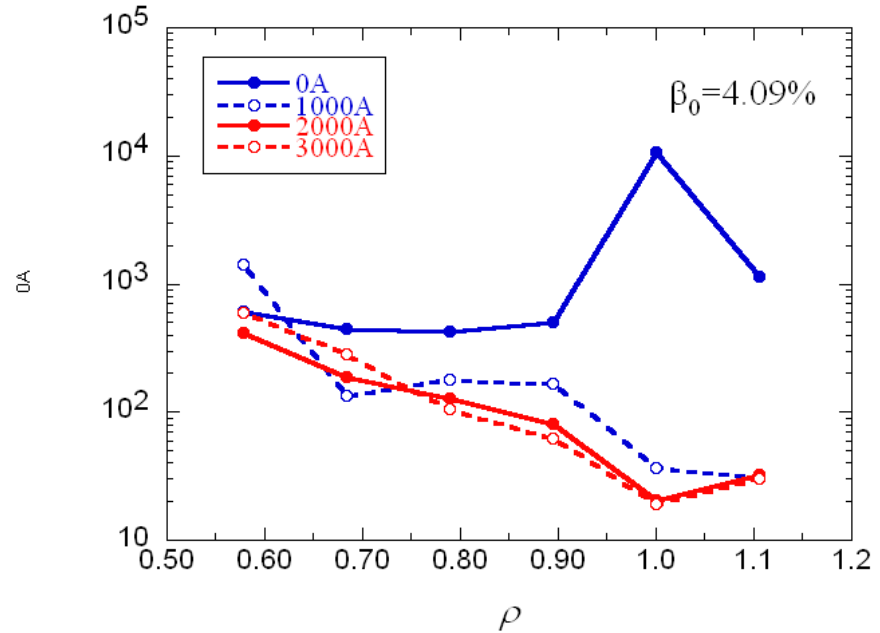
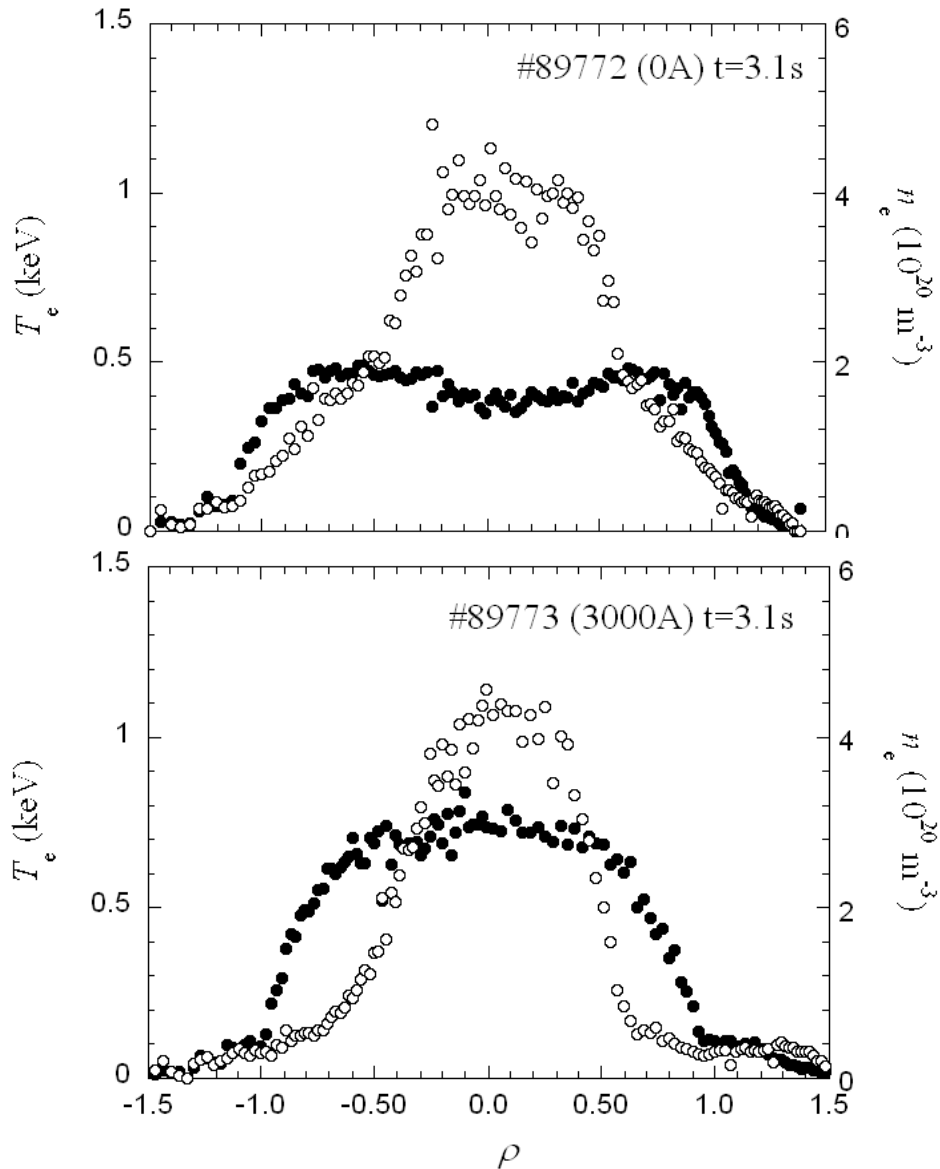
l : length of flux tube

L_K : Kolmogorov length

(e-folding length of
exponential increase of circumference)

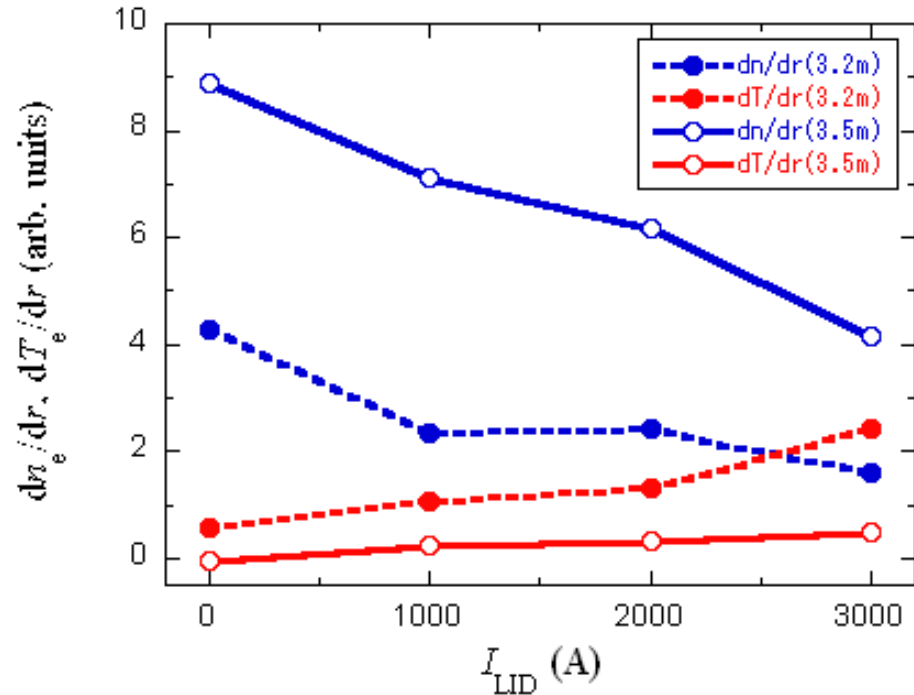
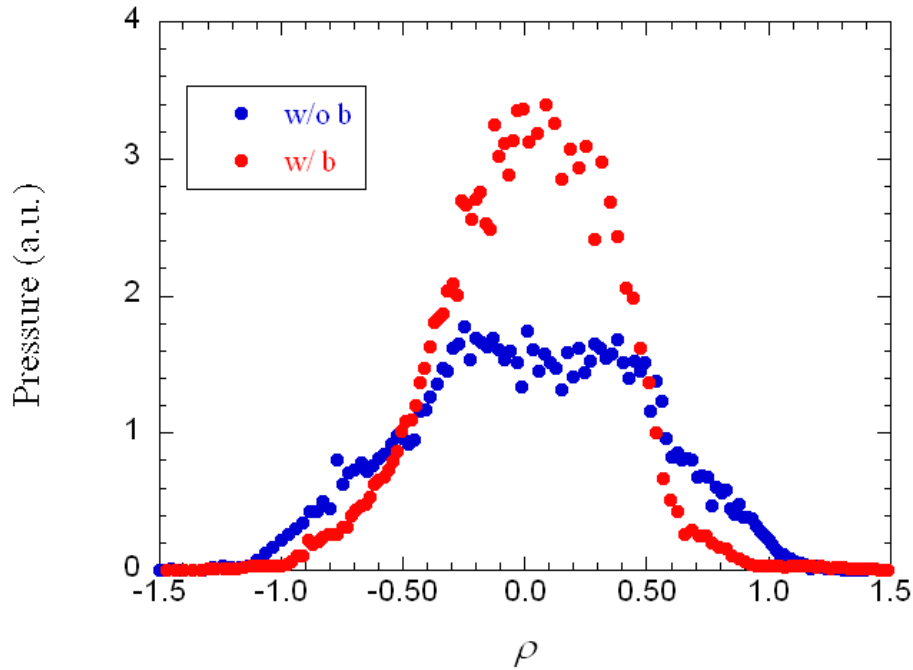


Relation between L_k and profiles



- L_k rapidly decreases when b is applied, especially in edge region
- Region where n_e is flat corresponds to L_k profile

Erogodization produces higher P_0



- higher P_0 is obtained with applying the perturbation
- gradient of n_e is well correlated with magnitude of perturbation
- higher P_0 is due to the increase of T_e
 - ==> because of the reduction of the edge density
 - ==> deep penetration of heating power

Summary

Effect of the high pressure core plasma on magnetic topology is numerically and experimentally investigated.

- Magnetic flux surfaces are strongly affected by the high pressure core plasma.
- Even a small perturbation, it enlarges magnetic islands and enhances the ergodization
- Particle transport seems to follow the magnetic structure, on the other hand, energy transport does not follow.
- Edge ergodization contributes the reduction of edge n_e , which leads to the higher central pressure.
 - ==> edge ergodization can control the core plasma performance.