# Behavior of high-pressure gasses injected to vacuum through a fast solenoid valve for supersonic cluster beam injection

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The supersonic cluster beam (SSCB) injection method is being developed as a new fueling method for the Large Helical Devise (LHD) experiment. As a first step, cluster formation at a room temperature has been investigated for various gasses using a fast solenoid valve for SSCB. Rayleigh scattering of laser light by the cluster is measured by a fast charge coupled device (CCD) camera. In the case of methane, nitrogen, and argon, clear scattering signals are observed at the high backing pressure of more than 3 - 4 MPa. In the case of hydrogen, helium, and neon, on the other hand, no scattering signal is detected at < 8 MPa. The scattering signals from argon and nitrogen clusters show approximately cubic dependence on the backing plenum pressure as expected from a model. Meanwhile, stronger pressure dependence than this has been found in the case of methane, where the scattering signal increases with the fifth power of the backing pressure at 3.2 MPa – 7 MPa, and it is further enhanced at > 7 MPa.

Keywords: fueling method, gas puffing, solenoid valve, cluster, Rayleigh scattering

#### 1. Introduction

A new fueling method of supersonic cluster beam (SSCB) injection, which is expected to be beneficial for deeper penetration of the fuel particles and higher fueling efficiency than that of gas puffing, is being developed for the Large Helical Devise (LHD) experiment. In SSCB, high-pressure hydrogen gas cooled to less than 77 K by a GM refrigerator will be injected to vacuum through a fast solenoid valve with a Laval nozzle. SSCB is an improved version of cluster jet injection (CJI) developed for HL-2M, where liquid nitrogen of 77 K is used for gas cooling [1], or, the supersonic gas injector (SGI) developed for NSTX, where a Laval nozzle is used to generate supersonic gas jet [2].

Although it is expected that SSCB will produce cluster, there is no established theory to predict the cluster size in a free jet expansion. However, it has been shown that the condition to produce cluster can be described by an empirical scaling parameter  $\Gamma^*$  that is proportional to so-called "Hagena parameter", k [3, 4],

$$\Gamma^* = k \, \frac{(d/\tan \alpha)^{0.85}}{T_0^{2.29}} P_0, \qquad (1)$$

where *d* is the nozzle diameter in  $\mu$ m,  $\alpha$  is the expansion half angle ( $\alpha = 45^{\circ}$  for sonic nozzles,  $\alpha < 45^{\circ}$  for supersonic),  $P_0$  is the backing plenum pressure in  $10^{-4}$  MPa, and  $T_0$  is the pre-expansion temperature in Kelvin.

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Massive condensation, where the cluster size exceeds 100 atoms/cluster, is generally observed for  $\Gamma^* > 1000$  [3, 4]. The parameter  $\Gamma^*$  as a function of the gas temperature is shown in Fig. 1, where  $d = 500 \ \mu\text{m}$ ,  $\alpha = 45^\circ$ , and  $P_0 = 4$  MPa are assumed. The nozzle diameter of  $d = 500 \ \mu\text{m}$  is equal to that of the valve used in this study. In this calculation, species-dependent k of 184, 3.85, 2360, 528, 185, and 1650 are used for H<sub>2</sub>, He, CH<sub>4</sub>, N<sub>2</sub>, Ne, and Ar, respectively [3]. The result implies that the gasses except helium are expected to form clusters at a room



Fig.1 Calculated results of the scaling parameter  $\Gamma^*$ , where  $d = 500 \ \mu\text{m}$ ,  $\alpha = 45^\circ$ ,  $P_0 = 4 \ \text{MPa}$ . The gasses except helium satisfy the condition of massive condensation  $\Gamma^* > 1000$  at room temperature (293 K).

temperature.

Before applying SSCB to LHD, the solenoid valve for SSCB has been tested at a room temperature in a test vacuum chamber. Various gasses shown in Fig.1 are used in the experiment to investigate the clustering condition.

# 2. Experimental Setup

The experimental setup is shown in Fig. 2. A solenoid valve of Parker-Hannifin Pulse Valve Series 99B07 with a 500 µm diameter orifice is used. This valve is equipped with a tapered nozzle. The available backing pressure is up to 8 MPa. This valve is set inside the vacuum chamber. The pressure in the vacuum chamber is measured by a pressure gauge of MKS Baratron capacitance manometer (MODEL#617A) set at the opposite side of the valve. When the valve is open, the gas flows from left to right in Fig. 2. Various gasses of H<sub>2</sub>, He, CH<sub>4</sub>, N<sub>2</sub>, Ne, and Ar are used in the experiment. A semiconductor laser of NEOARK LDP2-6535A with 650 nm standard wavelength and 35 mW power is set inside the chamber to perpendicularly intersect the gas flow. A beam dump is set at the opposite side of the laser and the valve is rolled by black tape in such a way that the stray light is suppressed. The distance between the valve exit and the laser chord is variable from 3.5 mm to 4.0 mm. The CCD camera of  $1280 \times 1024$  pixels is arranged in the direction perpendicular to both the gas flow and the laser beam. An example CCD image is shown in Fig. 3.

The total Rayleigh scattering signal  $S_{\rm RS}$  is proportional to the product of the scattering cross section  $\sigma$ and the number density of clusters  $n_{\rm c}$ . The cross section  $\sigma$ is proportional to the square of the averaged cluster size  $N_{\rm c}$ which is defined by the averaged number of atoms per cluster.  $n_{\rm c}$  is approximately given by the monomer density before becoming cluster,  $n_0$ , divided by  $N_{\rm c}$ , i.e.,  $n_{\rm c} \approx n_0/N_{\rm c}$ . The scattering signal  $S_{\rm RS}$  is proportional to  $P_0N_{\rm c}$  since the monomer density is proportional to the backing plenum pressure  $P_0$ . Farges et al. showed that  $N_{\rm c} \propto P_0^{-1.8-2.1}$ , assuming a multilayer icosahedral model [5, 6]. This means that the scattered light signal  $S_{\rm RS}$  should vary as



Fig. 2 Schematic of the experimental setup. The distance from the solenoid valve to the baratron pressure gauge is 4.2 m. Inside the chamber is pumped to less than 10<sup>-4</sup> Pa. The laser is set inside the chamber.



Solenoid valve

Scattering light

Fig. 3 The scattering light image detected by the CCD camera in the case of  $CH_4$ . The backing pressure is 8.0 MPa, and the exposure time is 10 ms. The laser beam direction (*y*) is perpendicular to the gas flow (*x*).

below,

$$S_{RS} \propto P_0^{2.8-3.1}$$
 (2)

#### 3. Results

Figure 4 shows the temporal behavior of the

exposure time 1 ms



Fig. 4 Typical temporal behavior of the scattering signal intensity. The working gas is Ar and the backing pressure is 7.1 MPa.

scattering signal intensity in the case of Ar. The valve is opened from 0 - 120 ms. The camera exposure time is fixed to 1 ms. As seen in the figure, the scattering light signal intensity is approximately constant during the valve open. Hereinafter, the valve opening time of 40 ms and the CCD camera trigger timing of 35 ms are fixed.

Typical scattering signal profiles in the case of CH<sub>4</sub> are shown in Fig. 5. The direction y in Fig. 5 is parallel to the laser light (see Fig. 3). The backing plenum pressure  $P_0$  is varied from 6.0 MPa to 8.0 MPa. While the exposure time of the CCD camera is fixed to 10 ms. Nearly symmetric profiles are also observed for N<sub>2</sub> and Ar.

Maxima of scattering signals are plotted in Fig. 6. The scattering signal increases with  $\sim P_0^{2.8}$  for Ar (Fig. 6(a)) and  $\sim P_0^{3.2}$  for N<sub>2</sub> (not shown) at room temperature. These results are similar to the expectation of Eq. (2) and the results of Ref [3]. However, for CH<sub>4</sub> at room temperature (Fig. 6(b)), it is found that the backing pressure dependence is stronger than expected, i.e.,  $S_{\rm RS} \propto P_0^{4.8}$  at  $P_0 < 7$  MPa and  $S_{\rm RS} \propto P_0^{8.6}$  at  $P_0 > 7$  MPa. This result is different from  $S_{\rm RS} \propto P_0^{2.8-3.1}$  (Eq.(2)).

The scattering signal is detected when the backing plenum pressure is above 3.2 MPa, 3.0 MPa, and 4.0 MPa for CH<sub>4</sub>, Ar, and N<sub>2</sub>, respectively. In the case of H<sub>2</sub>, He, and Ne, no scattering signal is detected up to 8.0 MPa of the backing plenum pressure. Although both hydrogen and



Fig. 5 Typical scatter signal profiles in the case of CH<sub>4</sub>. The direction y is parallel to the laser light. The backing pressure  $P_0$  is scanned from 6.0 MPa to 8.0MPa while the gas puff pulse length of 40 ms and the CCD camera trigger timing of 35 ms and the exposure time of 10 ms are fixed.

neon satisfy the condition of  $\Gamma^* > 1000$ , no cluster is detected. When the first scattering signal is detected by the CCD camera,  $\Gamma^*$  is 31000, 22000, and 9000 for CH<sub>4</sub>, Ar, and N<sub>2</sub>, respectively. These are much higher than the condition of massive condensation,  $\Gamma^* > 1000$  reported in Refs. [3, 4].









Fig. 6 Peak scattered signal as a function of the backing pressure for (a) argon at 293 K, and for (b) methane at 293 K.

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# 4. Summary

Clustering condition for various gasses at a room temperature has been investigated using the fast solenoid valve for SSCB. The Rayleigh scattering signal is detected by the CCD camera when the backing plenum pressure is above 3.2 MPa, 3.0 MPa, and 4.0 MPa for CH<sub>4</sub>, Ar, and N<sub>2</sub>, which correspond to 31000, 22000, and 9000 of  $\Gamma^*$ , respectively. In the case of H<sub>2</sub>, He, and Ne, no scattering signal is detected. Nearly symmetric shapes of the scattering signal profile in CH<sub>4</sub>, Ar, and N<sub>2</sub> are observed by the CCD camera. It has been found that the scattering signal intensity dependence on the backing pressure is similar to the expectation of Eq. (2) and the results of Ref [3] for N<sub>2</sub> and Ar. In the case of CH<sub>4</sub>, stronger backing pressure dependence is observed, i.e.,  $S_{\rm RS} \propto P_0^{4.8}$  below 7 MPa and  $S_{\rm RS} \propto P_0^{8.6}$  at  $P_0 > 7$  MPa. These are different from  $S_{\rm RS} \propto P_0^{2.8-3.1}$  (Eq.(2)) expected from the result of Farges et al.[5, 6]. Farges et al. estimated this relation assuming a multilayer icosahedral model for Ar cluster. This model seems to be reasonable also for N<sub>2</sub>, which shows similar pressure dependence as Ar. However, new structure model is required to determine the cluster size of CH<sub>4</sub>, which shows stronger backing pressure dependence than Ar and N<sub>2</sub>.

In this work, no signal is detected in the case of  $H_2$ and Ne, although the massive condensation condition is satisfied. Possible causes of this might be ; the noise level was larger than the scattering signal, or the laser power was too low. In order to observe the scattering signals of hydrogen and neon clusters, it is necessary to improve the experimental setup.

In SSCB, high-pressure gas cooled to less than 77 K will be injected to the fusion plasma through a fast solenoid valve with a Laval nozzle. At a low temperature below 77 K,  $\Gamma^*$  for H<sub>2</sub> increases to the similar level as those of Ar and CH<sub>4</sub> at a room temperature (see Fig. 1) where clear Rayleigh scattering signals are observed in this study. Therefore, it is expected that hydrogen cluster beam will be easily formed below 77 K.

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