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Anomalous Crater Formation in Pulsed-Laser-Illuminated Aluminum Slab and Debris Distribution

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Abstract

We have succeeded for the first time to simulate dynamic phase transition from metal to vapor. The crater size created by 650 mJ laser in 8nsec agrees well with the simulation including an elastic-plastic effect and its depth is 100 μ m which corresponds to anomalously large cutting speed of 10⁶ cm/sec for 8ns pulse. However, the simulation demonstrates that evaporation is caused by remnant hot gas during several 100 ns and occurs in the direction of 75 degree in an intermittent and unstable manner leading to a unique explanation on the angular distribution of debris.

Keywords : Laser, Phase Transition, Cutting, Simulation

Recently there has been a growing interest in laser generated debris from a solid target related to x-ray sources for lithography and microscope¹, femtosecond laser-matter interaction, the prepulse inprint problem in laser fusion, industrial application and so on. However, a numerical scheme to attack such dynamic processes has been beyond the capability of conventional one. In order to simulate these processes, we need a scheme that can simultaneously treat the dynamic evaporation process from metal to liquid and vapor, surface tension, elastic-plastic effect and so on. The major difficulty is to treat a very sharp interface between metal and vapor numerically. The Lagrangean representation is a convenient way to treat a sharp interface but can not be used for the case when the topology changes : the metal surface is melted and flies away as a droplet from the surface. The Eulerian representation with a fixed grid system is another choice to be considered but the numerical diffusion did not permit the use of it at sharp interface. The second problem is to simultaneously treat multi-materials adjacent each other which have different acoustic impedance like liquid and gas.

Recently we have proposed a unique scheme² to overcome these difficulties and succeeded for the first time to simulate dynamic phase transition from metal to liquid and gas. The code has been applied to various problems like laser-material processing³, Shoemaker-Levy 9⁴, vapor explosion⁵ and so on.

In this letter, we demonstrate that the code can give a unique physical insight into number of interesting features such as crater formation mechanism and angular distribution of debris in an

experiment performed at the Institute of Laser Engineering, Osaka University regarding the x-ray source development⁶. The experiment was performed in a configuration shown in Fig.1 : a YAG laser of 650 mJ in 8nsec is used to obliquely illuminate an aluminum slab target with an angle of 45 degree to the target normal. The same target was used for many shots by rotating and translating it around the axis so that the fresh surface is illuminated in each shot. Debris collectors of glass plates are located 3 cm from the target to measure the angular and size distribution. The diameter of a laser heated plasma is 200 μm observed with an x-ray pinhole camera filtered with a 0.2 μm thick Aluminum. Figure 2 shows an overview of the typical crater obtained in the experiment and simulation. The diameter of the crater is almost the same as that of x-ray picture and the depth is about 100 μm . This crater depth seems to be anomalous because the cutting speed is 100 $\mu\text{m}/8\text{nsec} \sim 10^6 \text{ cm/sec}$ if this crater should have been created during laser pulse. Since the speeds of sound wave and elastic wave inside aluminum are order of 10^5 cm/sec , the cutting speed is much larger than these speeds. Is this speed physically possible? This question is the starting point of the present work and its answer will be given below.

Figure 3 shows the time development of density contour from hydrodynamic calculation including thermal conduction, viscosity, elastic-plastic effect⁷, surface tension⁸, equation of state (EOS)⁹ and laser energy deposition. It is interesting to note that the crater is not formed during the laser pulse, but it develops gradually in the time scale of several 100 ns well after the laser pulse ended. The very high

temperature plasma more than a few tens of eV produced by the 8 ns laser pulse and most of them expands from the target. However, some of them still stay near the target for long time after the laser pulse because of recoil force from expanded plasmas and act as heat source to melt aluminum metal in the time scale of several 100 ns. Note that the energy required to make 100 μm crater is only 10 mJ and is only a few percent of laser energy. Therefore, only a small portion of energy remained after bulk expansion is enough to create the crater.

When the plasma temperature becomes less than melting temperature around 290 nsec (the time is measured from the laser peak), the stress of aluminum whose strength is 0.248 Mbar and yield strength is 2.2976 kbar is recovered and no distortion occurs after that time. This yield stress is quite important to determine the final crater size. Without yield strength, the crater develops further even after 490 ns and the depth becomes more than 300 μm although less difference is seen at the beginning around 90 ns.

The plasma heated crater formation leads to other interesting phenomena. Since the plasma acts not only as heat source but also as pressure source, the dynamic expansion of evaporated material at later time is strongly modified. Since a high pressure region is just in front of the evaporation surface, the vapor is forced to flow bypassing through a narrow channel between the metal surface and this pressure source. Therefore, the vapor preferentially flows toward a circumference with a large angle to the target normal. This effect is the exactly the same as that obtained in the experiment. Figure 4 shows a distribution of

debris from the targets. The histogram is the experimental result and it was drawn from 2000 shots accumulated. The distribution was obtained using an interference microscope by measuring the thickness of aluminum piled-up on the glass plate with an accuracy of $0.01 \mu\text{m}$ and then getting volume combined with the area measured by a microscope. The estimated error in the distribution is a few tens of percent and is small enough to compare with the simulation result. Clearly there exist two peaks around 0 and 75 degree. As in Fig.3, the plasma expands directed normally to the target at early time $t < 90 \text{ ns}$ and this expansion is a bulk part of the laser-heated plasma. As already stated, this expansion causes recoil force to the hot plasma surrounding aluminum surface. This main part of the expansion creates a peak at 0 degree. The triangles in Fig.4 show the distribution calculated from the time integration of mass flow ru up to $t = 90 \text{ ns}$. At an early stage $t < 90 \text{ ns}$, no peak appears around 75 degree. On the other hand, the expansion at later stage $t > 90 \text{ nsec}$ is limited to the sideward direction as stated before and creates the peak at 75 degree. Therefore accumulated distribution up to 490 ns shown by circles in Fig.4 increases mainly at 75 degree.

Although the simulation succeeded to replicate the two peaks at 0 and 75 degree, it failed to replicate a narrow distribution at 0 degree, where the simulation result shows a broad distribution. This broad distribution already started in the distribution accumulated up to 90 ns and perhaps is mainly due to an early time behavior of expansion during laser heating. We are not sure whether this difference is caused by some unknown mechanism to collimate the hot plasma expansion

or angular dependence of reflectivity at the collector plate. In the experimental result, we have assumed 100 % adhesion to the plate.

Simulation also predicts further interesting behavior. The expansion at $t < 40$ nsec is quite uniform because its temperature is quite high \sim a few tens of eV. The experiment supports this result and the debris around 0 degree is very fine and indistinguishable with an optical microscope. On the other hand, the simulation result at $t = 290$ nsec shows some filamentary streams flowing from the surface. The experiment also supports this result and the debris at 75 degree consists of 1 to 20 μm sized particles. Since the simulation is two-dimensional axisymmetric, we cannot estimate the particle size but we can suggest the origin of the filaments.

The evaporated gas at later time must flow through a very narrow channel to the circumference and therefore this narrow channel creates thermal insulation layer between metal surface and heat source. Then heat flow is instantaneously suppressed reducing evaporation. If the evaporation is suppressed, the narrow channel disappears and the heat source directly contact the metal surface. Thus the heat flow from heat source to metal surface is recovered causing again evaporation and the formation of narrow channel of thermal insulation layer. Repeating this process, the intermittent evaporation occurs and the filamentary streams appear as shown in Fig.3. Although the detailed time evolution is not seen from Fig.3, the animation picture shows that, during flight to vacuum, condensation occurs from low density clouds to well defined filaments seen in Fig.3 at 290 ns.

The phenomena analyzed here implies the very important effect in short-pulse, high-intensity laser-matter interaction. Even if the laser pulse is short and is not sufficiently long to create the crater, the remaining high temperature plasma can play a role to create a large crater if the laser intensity is sufficiently high ($\sim 10^{12}$ W/cm² in the present case). Most importantly, we now have a very useful tool to analyze the dynamic phase transition and related instabilities.

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- ⁹ The EOS is modified from that used in ref.3 with Clausius-Clapeyron’s law considering latent heat in the region of two phases.

Figure Captions

Fig.1 : Experimental configuration. An aluminum slab target is placed in a vacuum chamber of 5×10^{-5} Torr and the debris collector is located 3cm from the target.

Fig.2 : Typical shape of the crater obtained (left) from the experiment and (right) from the simulation, where the left half of each figure is aluminum.

Fig.3 : Top, middle and bottom figures show density contours at 40 ns, 90 ns, and 290 ns from the laser peak. In the figures, laser light is incident from the right and logarithmic density is shown : purple, green, yellow, red show the density of $2.7 \cdot 10^{-2}$, 10^{-3} , background density ($<10^{-4}$) g/cm^3 , respectively. The vertical length of the figure is $800 \mu\text{m}$.

Fig.4 : Debris distribution. 0 degree corresponds to the target normal. The histogram shows the experimental result, while circles and triangles show the accumulated mass from the simulation at 490 ns and 90 ns, respectively.

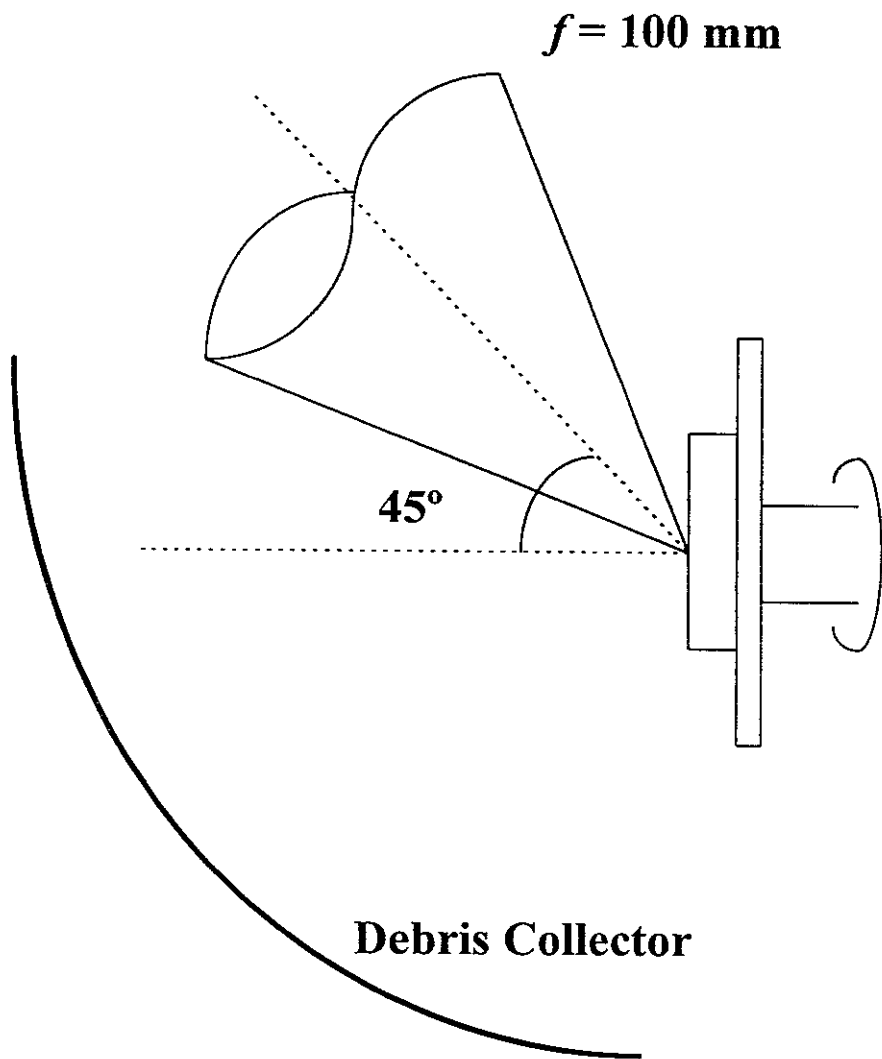


Fig.1

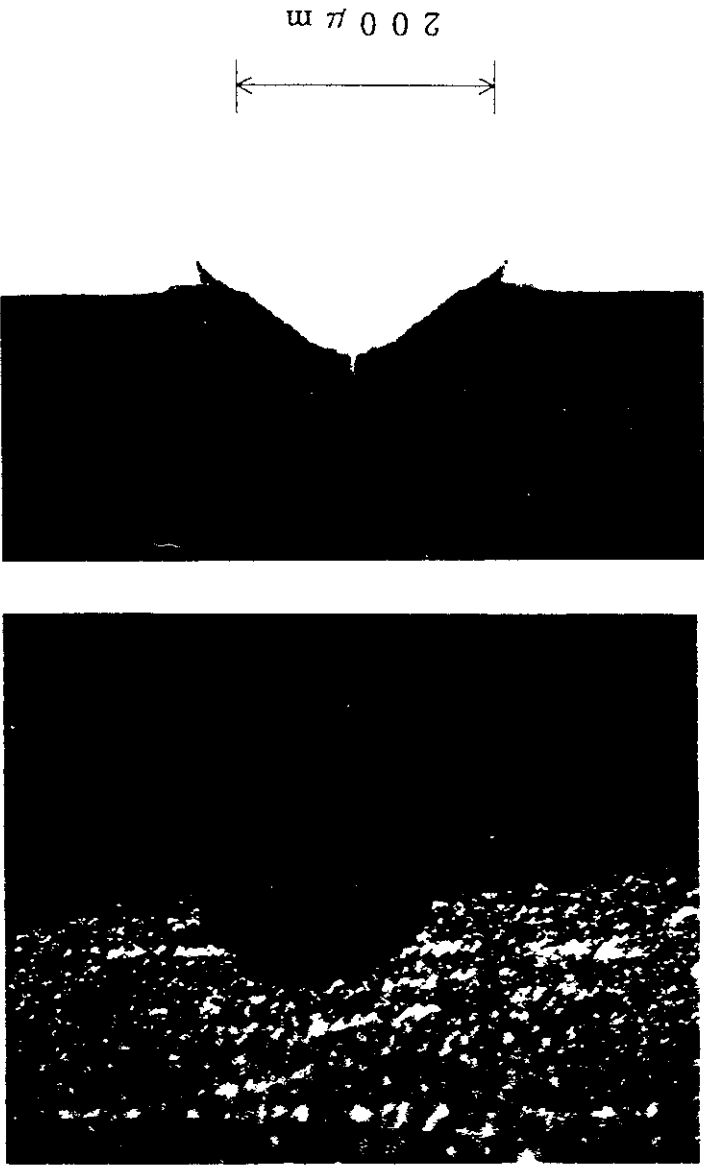


Fig.2

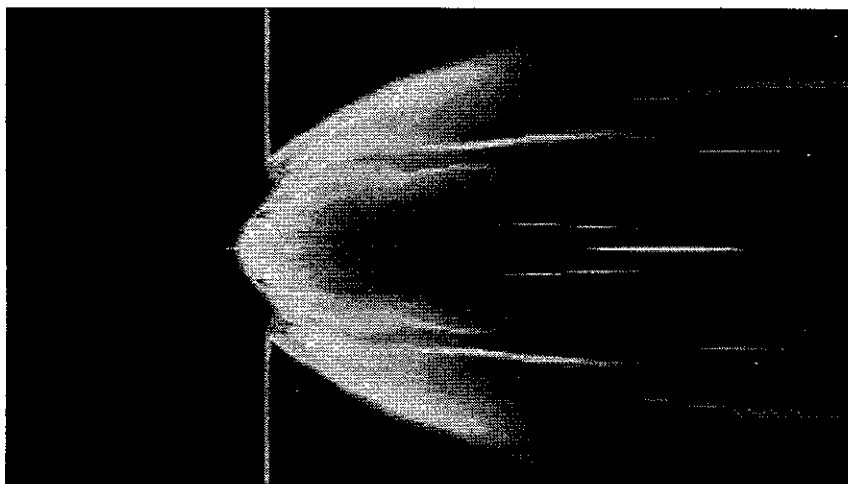
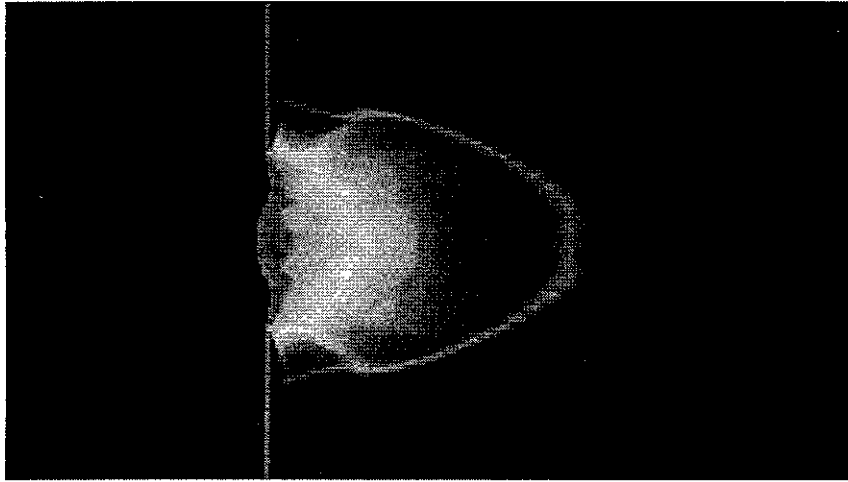


Fig.3

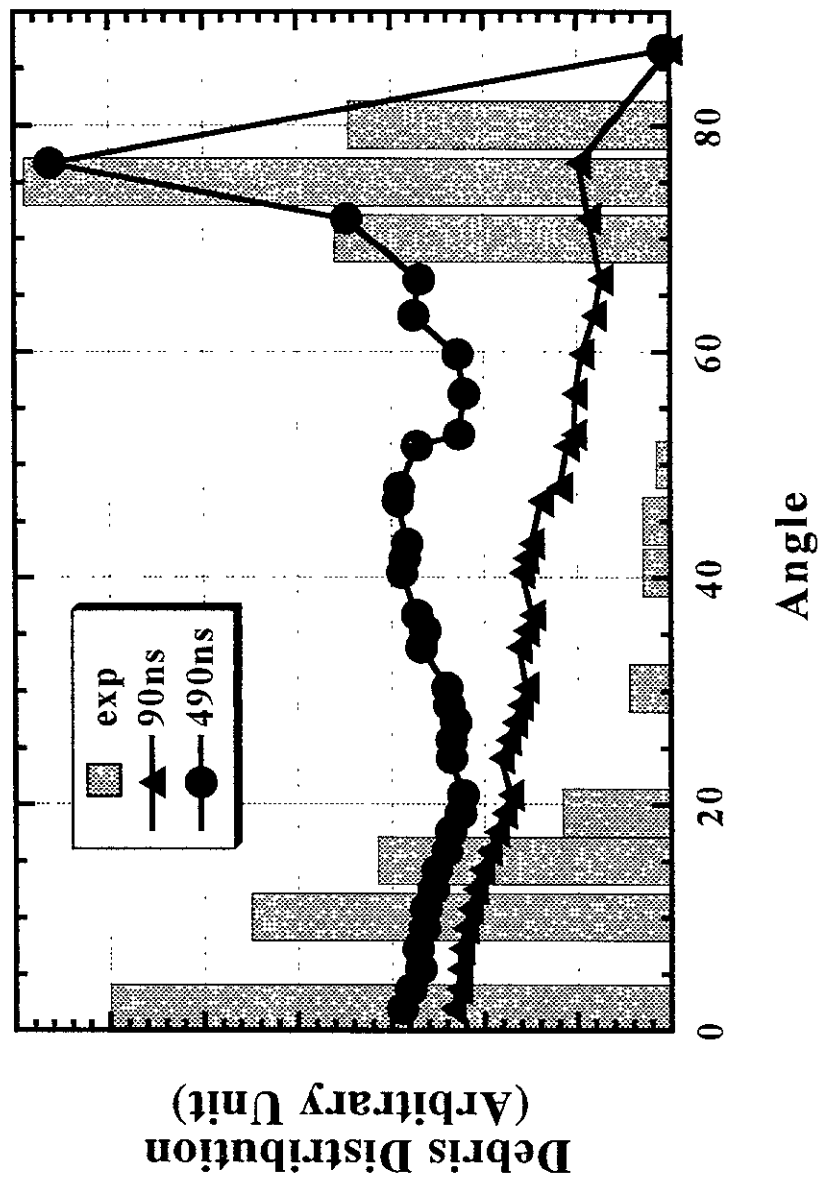


Fig.4

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