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RESEARCH REPORT
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**H⁻ ion source using a localized virtual magnetic filter in the plasma electrode:
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Abstract

A new multicusp H⁻ ion source using a Localized Virtual magnetic filter of type I [Ref.6] in the plasma electrode is investigated. A multipole (MP) arrangement with a spacing of 10 mm of the magnet bars holds an extraction hole, optimizing the efficient production of high H⁻ current, and at the same time only a small electron component was co-extracted with the H⁻ ions. The local filter arrangement separates the beam electrons at a low energy. It is shown that the co-extracted total electron current is determined principally by the integrated magnetic field flux (Gcm) of the local filter with an extraction system at a constant extraction voltage. When the value of the Gcm is increased, the total electron component is reduced, while the H⁻ electrical efficiency had a broad maximum around the optimized value of the Gcm. A thicker plasma electrode should be necessary for sufficient reduction of electron current. In pure hydrogen operation, the achieved current density of H⁻ is 10 mA/cm². When Cs was seeded in a filter optimized for pure volume mode H⁻ production, the maximum H⁻ current density obtained is 51 mA/cm² and the ratio $I_{\text{ele}} / I_{\text{H}^-}$ is ~0.4 without applying a bias potential.

Keywords: hydrogen negative ion, ion source, magnetic filter, cesium deposition, NBI

I. Introduction

A small ion source with a single hole has been used to do basic study of H^+ ion source with various arrangements of a new magnetic filter. The magnetic filter originated by K.Leung et al [1] produced two regions of a source plasma (i.e., a driver region and a target region) with different electron temperature. We have made two observations from our experiments on the magnetic filter in a pure hydrogen plasma source. Firstly, the H^+ yield for the same filling pressure after the optimization of the magnetic filter tended not to show a preference for the different three types (type I, II, III) of magnet geometries [2] with a H^+ current density up to several mA/cm^2 . The electrical efficiency for H^+ ion production, however, differed for each magnet configuration. Secondly, the electrons co-extracted from the plasma decreased with an increase of the filter magnetic field strength [3]. Furthermore, from the view of Hiske's modeling for the volume plasma source [4], an optimized magnetic filter aiming for the highest extraction current density of H^+ ions should make a shallow dense plasma close to the plasma electrode with the extraction hole in the target region of the plasma.

It might be deduced that a criterion of the optimum arrangement of the magnetic filter for H^+ ion source is to aim for high current density of the extracted H^+ ions, and at the same time a small electron beam component extracted from the plasma. In most cases, if the latter is gotten, the former would be automatically gotten [5]. A virtual magnetic filter is proposed in the plasma electrode localized in the vicinity of each extraction hole. This local filter uses pairs of small permanent magnets, adjacent to each aperture as will be shown in Fig.2. We call this Localized Virtual Magnetic Filter of type I in the plasma electrode, i.e., a localized virtual magnetic filter H^+ ion source. A similar magnet arrangement has been used for electron suppression, instead of filter action, in the literature.

In this paper we will investigate/optimize the ion source characteristics of H^+ yield and the electron beam component with the local filter of type I, especially for the magnet arrangement of the multipole magnetic field. The characteristics of the Cesium seeded H^+ source will also be tested. Detailed data for the comparison of different magnetic filter arrangements of type I will be submitted for publication. The first

results of experiments for the local filter had been reported in ref.[6], and those for multiholes in ref.[7] (with a vacuum-immersed H^+ source).

II. Ion source and the local filter arrangement

The ion source consists of a multipole plasma container which has a square cross section of $26(W) \times 26(L) \text{ cm}^2$, and a small accelerator with a single extraction hole of diameter 4 mm (Fig.1). A movable back plate is fixed at 12 cm deep, creating a shallow chamber, to raise the discharge power density. Behind the backplate, a cryobottle is filled with liquid nitrogen to trap condensable impurity gases during all the phases of the study including the experiments of Cesium introduction. All six surfaces of the plasma source are covered by multipole magnets of SmCo5. They work as a discharge anode. The plasma electrode over $5 \times 12 \text{ cm}^2$ around the beam hole has a local magnetic filter, and is isolated from the anode electrically and thermally. Unless otherwise stated, the plasma electrode is connected electrically to the anode without biasing. Eight 1.5 mm-diam M shaped tungsten filaments are extended into the field-free space. The hydrogen gas pressure in the plasma source was measured in the absence of the discharge. The discharge voltage ranged from 40 ~ 80 V and the discharge current ranged up to 800 A. The pulse duration for the beam ranged from 0.1 s to 1 s.

The small accelerator system uses an accelerator consisting of six electrodes, i.e., like an accel-decel system (which is called the advanced-accelerator of the LBL group). The gaps between the electrodes are 2 mm. The first electrode or plasma electrode, along with the entire ion source, is at a negative high potential for beam acceleration. No magnets are imbedded in the extraction or the second electrode. The primary reason for the choice of the six electrode accelerator in present source, instead of a usual two electrode system with magnets, is that the fringe magnetic field produced by magnets in the fourth electrode becomes very weak at the location inside the plasma source. Modification for the magnet were made to the first electrode of the local filter, i.e. the other electrodes were unchanged. The total negative-ion current I_{H^-} was measured by a calorimeter plate at a distance of 11 or 25 cm from the accelerator. High voltage power supply drain currents, I_{HV} consist of the sum of the electron current I_{e} and

I_{H^-} . Acceleration/extraction energy was ~ 7 keV.

Cesium was seeded by a Cs oven (not shown) through the side port. The material of the plasma electrode was changed from copper to molybdenum for Cs operation. Also thin molybdenum liners were added in front of the six inside-surfaces. One gram of Cs was used for four introductions.

A magnetic geometry of 3×3 SmCo5 (Fig.2) generates a Multi-Pole magnetic field (i.e., 3×3 MP). The magnet are oriented so that the magnetization is perpendicular to the plasma electrode. The magnets are buried midway in the grooves of the 4 mm thick plasma electrode, which has an extraction hole. Four plasma electrodes which have different spacing between the magnet bars, i.e. 10 (Fig.2), 15, 20, and 40 in mm over 5×12 cm², were fabricated. The target plasma in the plasma source occupies a plasma volume localized in a region with a high magnetic field very near the plasma electrode, while the driver region occupies the most of the remainder. The line integral of the magnetic field strength along the beam axis was taken from the center of the plasma electrode to a position of twice the length of the magnet gaps.

III. Experimental results

H⁻ yield; The H^- current increases with the discharge power but saturates at a high discharge power (in Fig.3). The maximum value of the negative-ion current increases, when the discharge power as well as the gas pressure are increased. The optimum spacing for 3×3 MP became short, i.e., 10 mm, for a high discharge power plasma at a higher gas pressure, and long, i.e., 15~20 mm, for a low discharge power plasma. These characteristics are similar to ones for the rod filter, H^- source studied by K.Leung et al. A maximum H^- ion current density of 10 mA/cm² is achieved with 3Pa. The gas pressure during the discharge-on time was reduced compared with the initial gas pressure due to plasma pumping. There was a larger reduction of the pressure with the discharge power.

The extracted electron components vs the line-integrated magnetic field strength: The high voltage drain current I_{HV} (\sim the total electron current) increases linearly with the discharge power, and decreases with a decrease of the magnet gap at a given discharge power for the same gas pressure (in Fig.4). The ratio of I_{e1} / I_{H^-} becomes smaller(better) as the gap is decreased

(correspondingly strong magnetic field).

It is found from Fig.4 that the total electron component decreases with the increase of the line integrated magnetic field strength in the accelerator rather than the just the magnetic field strength. Even if the magnetic filter arrangement and/or the gap are changed, the accompanied electron behaviour tends to depend only on the value of G_{cm} . The necessary value of the line-integrated magnetic field is determined mainly by the electron energy, and the geometry of the extraction systems. The value of G_{cm} for efficient H^- production [2] should take priority in the local filter source over the value for suppressing electrons, although H^- production was optimized weakly with the magnetic field.

Characteristics of the Cesium seeded operation: Cs was introduced for the 3×3 MP with a gap of 10 mm, after a series of beam extraction experiment was executed. The amount of one Cs introduction was evaluated to be 0.1 ~ 0.25 gr. The Cs introduction was repeated, if required, until a saturation of efficiency of H^- production was observed.

When the number of Cs introductions (in Fig 5) is increased sufficiently, H^- production efficiency is increased by a factor of 3~5 compared to that for pure volume production. The H^- current increases in proportion to the discharge power, in contrast to the saturation in pure hydrogen plasma. When the discharge power (in Fig.5) and the pressure are low, the H^- production efficiency was low and saturates after small quantity of Cs introduction, i.e., one Cs introduction. For a high power discharge at high gas pressure, improvement of efficiency saturates after several depositions. The maximum achieved H^- ion current density is 51 mA/cm² [3] at a filling gas pressure of 3 Pa.

It is seen that I_{HV} is dramatically decreased with one~two Cs deposition (in Fig.6), while H^- efficiency is still low. By adding more Cs, I_{HV} current starts again to increase due to increasing H^- current. The ratio of I_{e1} / I_{H^-} of 0.3 to 0.5 at H^- ion current density of 51 mA/cm² occurs without applying a bias potential.

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Figure captions,

Fig.1 H^- ion source with local (3 x3 MP) magnetic filter of type I.

Fig.2 Localized Virtual filters in the plasma (first) electrode in the upper figure. Distribution of the transverse component of the magnetic field along the beam Z-axis in the lower figure.

Fig.3 Dependence of the H^- ion current I_{H^-} on the discharge power.

Fig.4 Dependence of HV drain current I_{HV} (where interpolation are made) on the line-integrated magnetic field strength. Different symbols correspond to I_{H^-} currents with different gas pressures of 1, 2, and 3 Pa for three discharge powers of 15, 30, and 45 kW.

Fig.5 H^- yield I_{H^-} as a function of the discharge power with the integrated number (cycles) of Cs introduction as a parameter. Curve number (N) corresponds to the data after Nth introduction of Cs. 3Pa of H_2 .

Fig.6 Dependence of the HV drain current I_{HV} on the discharge power with run as same as in Fig.5.

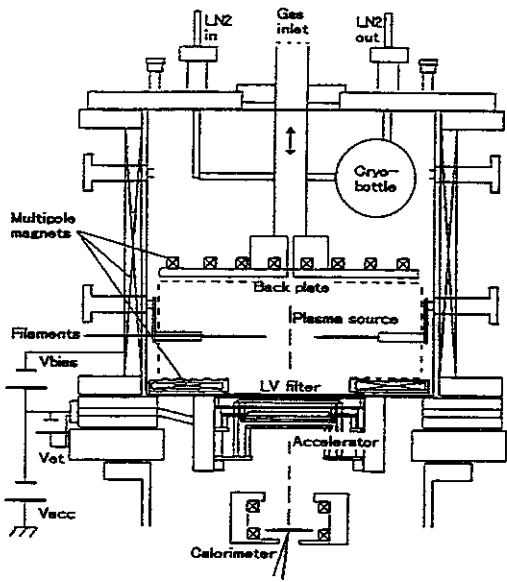


Fig. 1

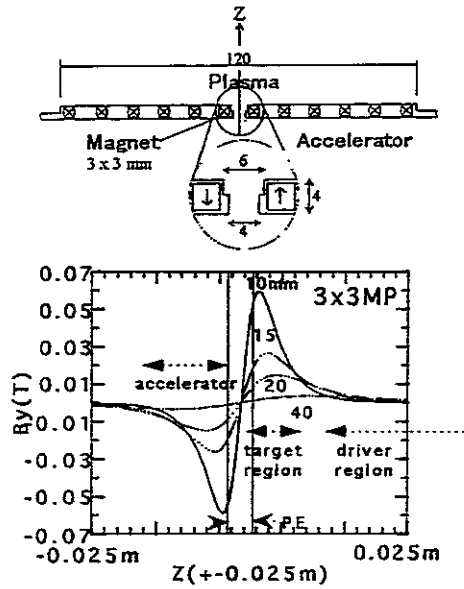


Fig. 2

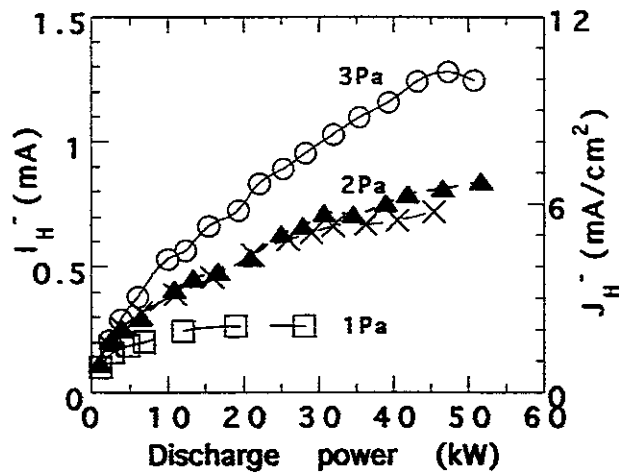


Fig. 3

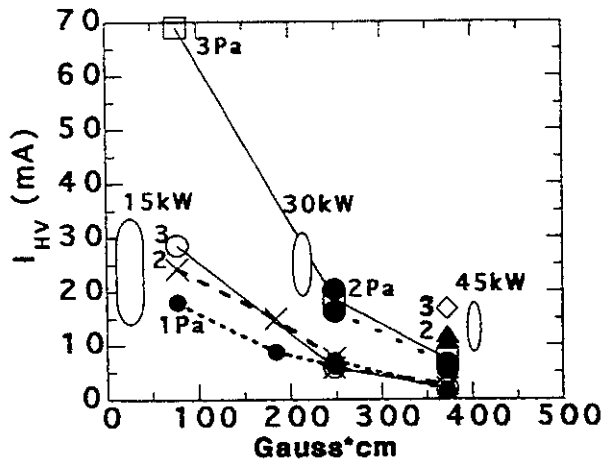


Fig. 4

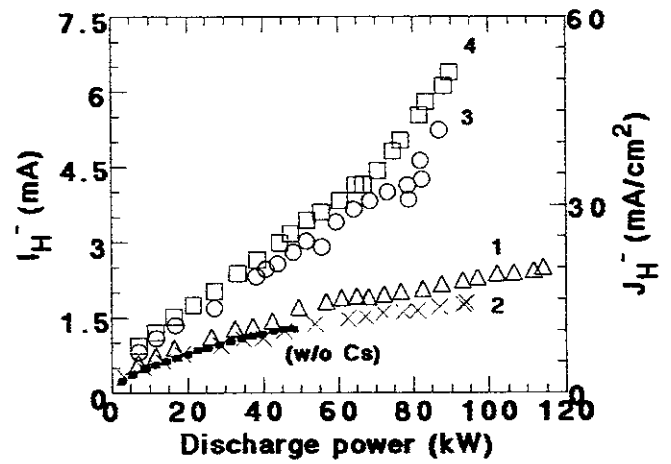


Fig. 5

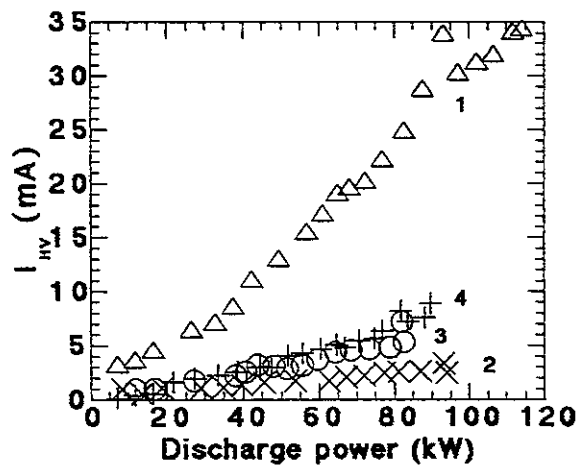


Fig. 6

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