Sputtering Yield Formula for B₄C Irradiated with Monoenergetic Ions at Normal Incidence

T. Ono, T. Kawamura, K. Ishii and Y. Yamamura

(Received - Mar. 29, 1996)

NIFS-DATA-34 Apr. 1996

RESEARCH REPORT NIFS-DATA Series

This report was prepared as a preprint of compilation of evaluated atomic, molecular, plasma-wall interaction, or nuclear data for fusion research, performed as a collaboration research of the Data and Planning Center, the National Institute for Fusion Science (NIFS) of Japan. This document is intended for future publication in a journal or data book after some rearrangements of its contents.

Inquiries about copyright and reproduction should be addressed to the Research Information Center, National Institute for Fusion Science, Nagoya 464-01, Japan.
Erratum: eq. (6) on p. 4 is to be replaced by

\[ \alpha^* = 0.249(M_k/M_1)^{0.56} + 0.0035(M_k/M_1)^{1.5} \]
\[ = 0.088(M_k/M_1)^{0.15} + 0.165(M_k/M_1) \]

\[ M_1 \leq M_k \]
\[ M_1 \geq M_k \]  (6)
Sputtering Yield Formula for B₄C Irradiated with Monoenergetic Ions at Normal Incidence

T. Ono¹, T. Kawamura², K. Ishii³ and Y. Yamamura³

¹ Nagoya University, Nagoya 464-01, Japan
² National Institute for Fusion Science, Nagoya 464-01, Japan
³ Okayama University of Science, Ridaicho, Okayama 700, Japan

Abstract

To fill a lack in sputtering yield data for a B₄C material which may be a promising plasma-facing material in fusion devices, yields of H⁺, D⁺, T⁺, He⁺, B⁺, C⁺, Ne⁺, Ar⁺ and Kr⁺ ions were calculated for this multi-component material for normal incidence with a computer simulation code ACAT. A fitting formula of sputtering yield for a B₄C was proposed based on an empirical formula for monoatomic target materials at normal incidence. By fitting the formula to the calculated data, best-fit values of the parameters included in it were derived for the material. Good agreement was found between the formula and the data. Thus, the formula proposed for the multi-component material provides a way to estimate the erosion of a B₄C material irradiated with above ions at normal incidence. Preferential sputtering for this material was also mentioned briefly.

KEYWORDS: physical sputtering, boron carbide, fitting formula, normal incidence, preferential sputtering
1. Introduction

In fusion devices the areas of the materials facing a plasma have to receive large fluxes of energy and particles. For a magnetically confined plasma, these fluxes are concentrated on divertor plates or limiters, so that these materials may be seriously damaged and eroded. Impurities, which are the constituents of the materials, are then generated through erosion. Impurity reduction is an important problem for fusion plasmas, since impurities introduced then into a core plasma cause core plasma cooling by radiation and dilution of the fuel. To prevent accumulation of impurities in a core plasma, erosion of plasma-facing materials in a fusion device should be very low.

Boronization, a recently developed wall conditioning, is a thin boron coating onto plasma-facing surfaces using a boron-containing gas mixture. Boronization executed in major tokamaks has proved to be effective to reduce metallic impurities and oxygen as well.\(^1\)\textsuperscript{-6} The thickness of the boron layers used was usually limited to an order of several 100nm. However, such thin layers are not sufficient to keep lower the impurity level and recycling hydrogen for many high-power discharges. Hence, a few hundred micrometers thick B\(_4\)C layers have been prepared and investigated on their hydrogen retention characteristics.\(^7\) It was shown that B\(_4\)C has a lower release temperature of implanted deuterium than graphite which is a commonly used plasma-facing material in fusion devices. In addition, B\(_4\)C was installed in the tokamaks as divertor tiles or coatings.\(^8\)\textsuperscript{,9} By the combined effects of boronization with decaborane and the boron burst from the B\(_4\)C tiles, oxygen impurity was reduced to the noise detection level.\(^8\) The B\(_4\)C coatings also showed excellent durability under a real tokamak divertor plasma.\(^9\) Thus, B\(_4\)C is thought to be one of the potential candidates for plasma-facing materials in fusion devices.

Sputtering is one of the main processes of erosion of a material exposed to ion bombardment. Sputtering yield depends strongly on the incident energy of a projectile. Hence, it is important to know energy dependence of ion-induced sputtering yield of plasma ions for B\(_4\)C. The experimental data available now are not enough for the purpose. To fill a lack in the yield data for this material, sputtering yield was calculated for nine kinds of ions at normal incidence. The ions are H\(^+\), D\(^+\), T\(^+\), He\(^+\), B\(^+\), C\(^+\), Ne\(^+\), Ar\(^+\) and Kr\(^+\). The calculations were done for the multi-component target with a Monte Carlo simulation code ACAT\(^10\) in a wide range of incident energy. The calculated data of sputtering yield will be compared with the experimental data. Then, based on a Yamamura and Tawara formula\(^11\) that represents sputtering yield for monoatomic target materials at normal incidence, a fitting formula of sputtering yield for a B\(_4\)C material will be proposed. By fitting the formula to above calculated data, four unknown parameters included in the formula will be determined. In particular, sputtering threshold energy, which is one of these parameters, will be
presented in a functional form.

Preferential sputtering for this material will also be referred to briefly.

2. Fitting Formula of Sputtering Yield for Multi-Component Materials

Based on the Yamamura and Tawara formula for monoatomic target materials, an empirical fitting formula for multi-component materials $Y(E)$ is proposed:

$$Y(E) = 0.042 \frac{Q <\alpha^* S_n(E)>}{<U>} \frac{1}{1 + \Gamma \langle ke^{0.33}\rangle} \left[ 1 - \sqrt[3]{\frac{E_{th}}{E}} \right]^3,$$  \hspace{1cm} (1)

where the factor $\Gamma$ has the following form:

$$\Gamma = \frac{W}{1 + (M_1/\gamma)^3},$$  \hspace{1cm} (2)

where $M_1$ is the mass of a projectile in u; $E$ is the incident energy of a projectile; $E_{th}$ is the sputtering threshold energy; the numerical factor is expressed in units of $\text{Å}^{-2}$; $Q$, $W$, $E_{th}$ and $s$ are the fitting parameters of a target material: among those, $Q$ and $W$ depend on a target material, while $E_{th}$ is a function of a mass ratio of a projectile and a mean mass of atoms in a target material: the latter mass is defined as

$$<M_2> = \sum_k c_k M_k,$$  \hspace{1cm} (3)

where $c_k$ and $M_k$ are the constituent ratio and the mass of a $k$-atom of a multi-component material: the power $s$ in eq.(1) is slightly dependent on both a projectile and a target material: the physical quantities other than $Q$, $W$, $E_{th}$ and $s$ in eqs.(1) and (2) are derived uniquely from the ones for a monoatomic material: $<U>$ is the effective surface binding energy of a multi-component material and it is defined as

$$<U> = \sum_k c_k U_k,$$  \hspace{1cm} (4)

where $U_k$ means the surface binding energy of a pure $k$-material: the term $<\alpha^* S_n>$ is calculated as

$$<\alpha^* S_n(E)> = \sum_k c_k \alpha^*(M_k/M_1)S_{nk}(E),$$  \hspace{1cm} (5)

where the best-fit values of $\alpha^*$ are given for monoatomic materials as a function of a mass ratio $M_k/M_1$ in the form of \(^{11}\)}
\[ \alpha^* = 0.249 \left( \frac{M_1}{M_1} \right)^{0.56} + 0.0035 \left( \frac{M_1}{M_1} \right)^{1.5} \]
\[ = 0.088 \left( \frac{M_1}{M_1} \right)^{0.15} + 0.165 \left( \frac{M_1}{M_1} \right) \quad M_1 \geq M_k \]
\[ M_1 \leq M_k \]

\[ S_{\text{n}(E)} \] is the nuclear stopping cross-section of a k-atom and it is expressed as

\[ S_{\text{n}(E)} = 84.78 \frac{Z_1 Z_k}{(Z_1^{2/3} + Z_k^{2/3})^{1/2}} \frac{M_1}{M_1 + M_k} s_{\text{n}}^{\text{TF}}(\varepsilon_k) \]

in units of eV Å^2/atom: Z_k is the atomic number of a k-atom: \( s_{\text{n}}^{\text{TF}}(\varepsilon_k) \) is the reduced nuclear stopping cross-section and it is approximated\(^{12}\) by

\[ s_{\text{n}}^{\text{TF}}(\varepsilon_k) = \frac{3.441 \sqrt{\varepsilon_k} \ln(\varepsilon_k + 2.718)}{1 + 6.355 \sqrt{\varepsilon_k} + \varepsilon_k (6.882 \sqrt{\varepsilon_k} + 1.708)} \]

where \( \varepsilon_k \) is a reduced energy and it is given by

\[ \varepsilon_k = \frac{0.03255}{Z_1 Z_k (Z_1^{2/3} + Z_k^{2/3})^{1/2}} \frac{M_k}{M_1 + M_k} E \text{ (eV)} \]

The term \( < k_e e^{0.3} > \) is given as

\[ < k_e e^{0.3} > = \sum_k c_k k_{ek} \varepsilon_k^{0.3} \]

where \( k_{ek} \) is the Lindhard electronic stopping coefficient\(^{13}\) of a k-atom and it is defined as

\[ k_{ek} = 0.079 \frac{(M_1 + M_k)^{3/2}}{M_1^{3/2} M_k^{1/2}} \frac{Z_1^{2/3} Z_k^{1/2}}{(Z_1^{2/3} + Z_k^{2/3})^{3/4}} \]

3. Calculated Yield Data and Fitting Procedure

Sputtering yield of a B_4C material was calculated for monoenergetic ions with normal incidence in the energy range of a few tens to 10^4 eV or higher. The calculations were done for the multi-component target with a Monte Carlo simulation code ACAT\(^{10}\) that simulates atomic collisions in an amorphous target material based on a binary collision approximation. The density of a B_4C material considered is 2.52[g/cm^3]. The effective surface binding energy for an i-atom in a multi-
component material \( U_i \) is given by\(^{14,15}\)

\[
U_i = \sum_k c_k U_{ik},
\]

where

\[
U_{ik} = \frac{1}{2}(U_{ii} + U_{ik}),
\]

where \( U_{ii} \) is the nearest neighbor \( i-i \) bond strength and similarly for \( U_{kk} \), and \( U_{ik} \); \( U_{ii} \) and \( U_{kk} \) are approximated by the heats of sublimation of pure elemental \( i \)- and \( j \)-materials. For self-sputtering for \( \text{B}^+ \) and \( \text{C}^+ \) ions, reflection of the incident ions is not included in the yield values. The calculated data marked with stars are shown in Figs. 1-9. The data marked with "A" in these figures are those obtained by experiments\(^{16-18}\) and compiled by Eckstein et al.\(^{19}\) Good agreement is seen between the calculated and the experimental data for \( \text{H}^+ \), \( \text{D}^+ \) and \( \text{C}^+ \) ions if the data scattering is taken into account at the same energies for the experimental data. Slight differences between the ACAT data and the experiments, however, exist for \( \text{He}^+ \) and \( \text{Ne}^+ \) ions.

From eq.(2), \( \Gamma \) are expected to become much smaller than unity for large \( M_1 \), considering that the values of \( W \) are 4.39 and 1.70 for \( \text{B} \) and \( \text{C} \).\(^{11}\) From eqs (9)-(11), the term \( \langle k_e E^{0.3} \rangle \) can be shown easily to be also much smaller than unity for large \( M_1 \) in the energy range concerned. Because of these facts, the term \( \Gamma \langle k_e E^{0.3} \rangle \) in eq.(1) can be neglected for large \( M_1 \). Thus, the best-fit values of \( Q \), \( E_{th} \) and \( s \) are obtained first by fitting the formula to the yield data for a large \( M_1 \). \( \Gamma \) approaches \( W \) in eq.(2) for small \( M_1 \), i.e., \( M_1 \) of hydrogen or deuterium. The term \( \Gamma \langle k_e E^{0.3} \rangle \) affects the formula in the high energy for small \( M_1 \). Secondly, with the use of these facts and the value of \( Q \) calculated above, the best-fit values of \( W \), \( E_{th} \) and \( s \) are calculated by fitting the formula to the data for hydrogen or deuterium. Finally, with the values of \( Q \) and \( W \) derived above, the best-fit values of \( E_{th} \) and \( s \) are derived by fitting the formula to the data for intermediately large \( M_1 \).

4. Results

The calculated data of sputtering yield distribute systematically over a wide range of incident
energy for all the ions considered, while the experimental data are scattered for H\(^+\) and D\(^+\) ions and are lacking for T\(^+\) and B\(^+\) ions. Hence, fitting of the yield formula was done with the calculated data. The result is shown with the thick solid lines in Figs.1-9. The power \( s \) in eq.(1) for B\(_4\)C has a best-fit functional form of

\[
s = \frac{10M_1}{1+4M_1} \quad .
\]

(14)

The best-fit values of \( Q \) and \( W \) derived by these fits are listed in Table 1, together with the other physical quantities included in the empirical formula. Through these fittings good agreement was achieved between the yield formula and the calculated data. The symbols in Fig.10 represent the best-fit values of \( E_{th} \) divided by \( <U> \). Using these data, a best-fit functional form of \( E_{th} \) was adopted:

\[
\frac{E_{th}}{<U>} = \frac{12.12 - 5.35(<M_2>/<M_1>) + 2.2(<M_2>/<M_1>)^2 \gamma \quad M_1 \geq <M_2>}{1 + 3.7(M_1/<M_2>) + 4.29(M_1/<M_2>)^2 \gamma \quad M_1 \leq <M_2>}
\]

(15)

where \( <M_2> \) is the mean mass of atoms in a B\(_4\)C material and it is expressed as

\[
<M_2> = \sum_k c_k M_k \quad .
\]

(16)

\( \gamma \) corresponds to the energy transfer factor for an elastic collision and it is defined as

\[
\gamma = \frac{4M_1<M_2>}{(M_1 + <M_2>)^2} \quad .
\]

(17)

The curve shown in Fig.10 is drawn using eq.(15). From this figure, good agreement is clear between eq.(15) and above-mentioned data. The thin solid lines in Figs.1-9 mean yield curves with threshold energies calculated with eq.(15). Good agreement is seen again between the ACAT data and the empirical formula. This result indicates that sputtering yield with a B\(_4\)C material may be estimated for unlisted ions with the present empirical formula.

Preferential sputtering occurring for a B\(_4\)C material is illustrated with the case of H\(^+\) ion
incidence as shown in Fig.11. The circles indicate a ratio of a partial sputtering yield of boron and the total sputtering yield which was shown in Fig.1, while the squares show that of a partial sputtering yield of carbon and the total sputtering yield. This figure indicates that carbon atoms are not sputtered for incident energies lower than the sputtering threshold energy of carbon atoms. Such strong preferential sputtering will drastically modify the surface composition of a B₄C material when that material is exposed to low-temperature plasma ions for a long time.

5. Conclusion

To fill a lack in the sputtering yield data for B₄C, sputtering yields were calculated for this multi-component material with a computer-simulation code ACAT for normal incidence of H⁺, D⁺, T⁺, He⁺, B⁺, C⁺, Ne⁺, Ar⁺ and Kr⁺ ions. Fairly good agreement was found between the calculated and the experimental data. A fitting formula of sputtering yield for a multi-component material of a B₄C was made based on an empirical formula for monoatomic target materials at normal incidence. By fitting it to the calculated data, best-fit values of the four parameters included in the formula were derived. From comparison with the data, the formula agreed well with all the data. Thus, it provides a way to estimate the erosion of a B₄C material irradiated with above ions at normal incidence. There is no ambiguity in the formula for a projectile, because the parameters included in the formula are functions of the physical quantities of a B₄C and its constituent materials and of a mass ratio of a projectile and a mean mass of atoms in the target material. This indicates that sputtering yield of the material may be estimated for unlisted ions with the present formula. Fitting of the formula for the other multi-component materials is in progress.

It was also pointed out that strong preferential sputtering takes place for this material at incident energies lower than about 10² eV. A detailed study of this phenomenon for the material, in which changes of the surface composition due to ion doses are taken into account, is also in progress.
References


11) Y.Yamamura and H.Tawara : to be published in At. Data and Nucl. Data Tables.


Figure Captions

Fig.1: Energy dependence of sputtering yield of B\textsubscript{4}C with H\textsuperscript{+}. The thick and thin solid lines mean yield curves with the best-fit threshold energy and that calculated by eq.(15), respectively. 
\(<M_2> = 11.05, \frac{<M_2>}{<M_1>} = 10.96, Q = 1.70, <U> = 6.09\text{eV}, W = 1.75, s = 2.00.\)

Fig.2: Energy dependence of sputtering yield of B\textsubscript{4}C with D\textsuperscript{+}. The thick and thin solid lines mean yield curves with the best-fit threshold energy and that calculated by eq.(15), respectively. 
\(<M_2> = 11.05, \frac{<M_2>}{<M_1>} = 5.49, Q = 1.70, <U> = 6.09\text{eV}, W = 1.75, s = 2.22.\)

Fig.3: Energy dependence of sputtering yield of B\textsubscript{4}C with T\textsuperscript{+}. The thick and thin solid lines mean yield curves with the best-fit threshold energy and that calculated by eq.(15), respectively. 
\(<M_2> = 11.05, \frac{<M_2>}{<M_1>} = 3.66, Q = 1.70, <U> = 6.09\text{eV}, W = 1.75, s = 2.31.\)

Fig.4: Energy dependence of sputtering yield of B\textsubscript{4}C with He\textsuperscript{+}. The thick and thin solid lines mean yield curves with the best-fit threshold energy and that calculated by eq.(15), respectively. 
\(<M_2> = 11.05, \frac{<M_2>}{<M_1>} = 2.76, Q = 1.70, <U> = 6.09\text{eV}, W = 1.75, s = 2.35.\)

Fig.5: Energy dependence of sputtering yield of B\textsubscript{4}C with B\textsuperscript{+}. The thick and thin solid lines mean yield curves with the best-fit threshold energy and that calculated by eq.(15), respectively. 
\(<M_2> = 11.05, \frac{<M_2>}{<M_1>} = 1.02, Q = 1.70, <U> = 6.09\text{eV}, W = 1.75, s = 2.44.\)

Fig.6: Energy dependence of sputtering yield of B\textsubscript{4}C with C\textsuperscript{+}. The thick and thin solid lines mean yield curves with the best-fit threshold energy and that calculated by eq.(15), respectively. 
\(<M_2> = 11.05, \frac{<M_2>}{<M_1>} = 0.92, Q = 1.70, <U> = 6.09\text{eV}, W = 1.75, s = 2.45.\)

Fig.7: Energy dependence of sputtering yield of B\textsubscript{4}C with Ne\textsuperscript{+}. The thick and thin solid lines mean yield curves with the best-fit threshold energy and that calculated by eq.(15), respectively. 
\(<M_2> = 11.05, \frac{<M_2>}{<M_1>} = 0.55, Q = 1.70, <U> = 6.09\text{eV}, W = 1.75, s = 2.47.\)

Fig.8: Energy dependence of sputtering yield of B\textsubscript{4}C with Ar\textsuperscript{+}. The thick and thin solid lines mean yield curves with the best-fit threshold energy and that calculated by eq.(15), respectively.
\(<M_2> = 11.05, \frac{<M_2>}{<M_1>} = 0.28, \ Q = 1.70, \ <U> = 6.09\text{eV}, \ W = 1.75, \ s = 2.48 \)

Fig.9: Energy dependence of sputtering yield of B\(_4\)C with Kr\(^+\). The thick and thin solid lines mean yield curves with the best-fit threshold energy and that calculated by eq.(15), respectively.
\(<M_2> = 11.05, \frac{<M_2>}{<M_1>} = 0.13, \ Q = 1.70, \ <U> = 6.09\text{eV}, \ W = 1.75, \ s = 2.49.\)

Fig.10: Relative sputtering threshold energy of B\(_4\)C as a function of a mass ratio \(<M_2>/<M_1>\). The symbols correspond with the best-fit values of \(E_{th}\) adopted for the yield fittings.

Fig.11: Energy dependence of ratios of partial sputtering yields of B and C and total sputtering yield of B\(_4\)C. The circles correspond with the former ratio and the squares with the latter one.

Table caption

Table 1: Physical parameters included in the empirical formula.
<table>
<thead>
<tr>
<th>( \Omega )</th>
<th>1.70</th>
</tr>
</thead>
<tbody>
<tr>
<td>( W )</td>
<td>1.75</td>
</tr>
<tr>
<td>( &lt;U&gt; )</td>
<td>6.09 eV</td>
</tr>
<tr>
<td>( &lt;M_z&gt; )</td>
<td>11.05</td>
</tr>
</tbody>
</table>
Fig. 1

$H^+ \rightarrow B_4C$

Sputtering yield (atoms/ion)

Energy (eV)


*: ACAT
Fig. 2

D$^+ \rightarrow B_4C

Sputtering yield (atoms/ion)

Energy (eV)


\*: ACAT
Fig. 3
He$^+$ → B$_4$C

Sputtering yield (atoms/ion)

Energy (eV)

★ : ACAT

Fig. 4
Fig. 5
Fig. 6
Fig. 7

Ne⁺ → B₄C

Sputtering yield (atoms/ion)

Energy (eV)

★ : ACAT
Fig. 8
Fig. 9
Fig. 10
Fig. 11
Publication List of NIFS-DATA Series

NIFS-DATA-1  Y. Yamamura, T. Takiguchi and H. Tawara,
Data Compilation of Angular Distributions of Sputtered Atoms; Jan. 1990

NIFS-DATA-2  T. Kato, J. Lang and K. E. Berrington,
Intensity Ratios of Emission Lines from OV Ions for Temperature and Density Diagnostics; Mar. 1990 [At Data and Nucl Data Tables 44(1990)133]

NIFS-DATA-3  T. Kaneko,
Partial Electronic Straggling Cross Sections of Atoms for Protons; Mar. 1990

NIFS-DATA-4  T. Fujimoto, K. Sawada and K. Takahata,
Cross Section for Production of Excited Hydrogen Atoms Following Dissociative Excitation of Molecular Hydrogen by Electron Impact; Mar. 1990

NIFS-DATA-5  H. Tawara,

NIFS-DATA-6  H. Tawara, Y. Itikawa, H. Nishimura, H. Tanaka and Y. Nakamura,

NIFS-DATA-7  H.Tawara,
Bibliography on Electron Transfer Processes in Ion-Ion/Atom/Molecule Collisions—Updated 1990--; Aug. 1990

Excitation Collision Strengths, Cross Sections and Rate Coefficients for OV, SiXI, FeXXIII, MoXXXIX by Electron Impact (1s²2s²-1s²2s2p-1s²2p² Transitions) Dec.1990

NIFS-DATA-9  T.Kaneko,
Partial and Total Electronic Stopping Cross Sections of Atoms and Solids for Protons; Dec. 1990

NIFS-DATA-10 K.Shima, N.Kuno, M.Yamanouchi and H.Tawara,
Equilibrium Charge Fraction of Ions of Z=4-92 (0.02-6 MeV/u) and Z=4-20 (Up to 40 MeV/u) Emerging from a Carbon Foil; Jan.1991 [AT.Data and Nucl. Data Tables 51(1992)173]
T. Kaneko, T. Nishihara, T. Taguchi, K. Nakagawa, M. Murakami, M. Hosono, S. Matsushima, K. Hayase, M. Moriya, Y. Matsukuma, K. Miura and Hiro Tawara,

Hiro Tawara,
*Total and Partial Cross Sections of Electron Transfer Processes for Be\(\text{II}^+\) and B\(\text{II}^+\) Ions in Collisions with H, H\(_2\) and He Gas Targets - Status in 1991-;* June 1991

*Partial and Total Electronic Stopping Cross Sections of Atoms for a Singly Charged Helium Ion: Part II;* Aug. 1991

T. Kato, K. Masai and M. Arnaud,
*Comparison of Ionization Rate Coefficients of Ions from Hydrogen through Nickel;* Sep. 1991

T. Kato, Y. Itikawa and K. Sakimoto,
*Compilation of Excitation Cross Sections for He Atoms by Electron Impact;* Mar. 1992

*Atomic Processes Relevant to Polarization Plasma Spectroscopy;* Apr. 1992

H. Tawara,
*Electron Stripping Cross Sections for Light Impurity Ions in Colliding with Atomic Hydrogens Relevant to Fusion Research;* Apr. 1992

T. Kato,

Hiro Tawara,
*Atomic and Molecular Data for H\(_2\)O, CO&CO\(_2\) Relevant to Edge Plasma Impurities*, Oct. 1992

Hiro. Tawara,
*Bibliography on Electron Transfer Processes in Ion-Ion/Atom/Molecule Collisions -Updated 1993-;* Apr. 1993
NIFS-DATA-21 J. Dubau and T. Kato,
Dielectronic Recombination Rate Coefficients to the Excited States of C I from C II; Aug. 1994

NIFS-DATA-22 T. Kawamura, T. Ono, Y. Yamamura,
Simulation Calculations of Physical Sputtering and Reflection Coefficient of Plasma-Irradiated Carbon Surface; Aug. 1994

NIFS-DATA-23 Y. Yamamura and H. Tawara,
Energy Dependence of Ion-Induced Sputtering Yields from Monoatomic Solids at Normal Incidence; Mar. 1995

NIFS-DATA-24 T. Kato, U. Safronova, A. Shlyaptseva, M. Cornille, J. Dubau,
Comparison of the Satellite Lines of H-like and He-like Spectra; Apr. 1995

NIFS-DATA-25 H. Tawara,
Roles of Atomic and Molecular Processes in Fusion Plasma Researches - from the cradle (plasma production) to the grave (after-burning) -; May 1995

NIFS-DATA-26 N. Toshima and H. Tawara
Excitation, Ionization, and Electron Capture Cross Sections of Atomic Hydrogen in Collisions with Multiply Charged Ions; July 1995

NIFS-DATA-27 V.P. Shevelko, H. Tawara and E. Salzborn,
Multiple-Ionization Cross Sections of Atoms and Positive Ions by Electron Impact; July 1995

NIFS-DATA-28 V.P. Shevelko and H. Tawara,
Cross Sections for Electron-Impact Induced Transitions Between Excited States in He: n, n'=2,3 and 4; Aug. 1995

NIFS-DATA-29 U.I. Safronova, M.S. Safronova and T. Kato,
Cross Sections and Rate Coefficients for Excitation of \Delta n = 1 Transitions in Li-like Ions with 6<Z<42; Sep. 1995

NIFS-DATA-30 T. Nishikawa, T. Kawachi, K. Nishihara and T. Fujimoto,
Recommended Atomic Data for Collisionsal-Radiative Model of Li-like Ions and Gain Calculation for Li-like Al Ions in the Recombining Plasma; Sep. 1995

NIFS-DATA-31 Y. Yamamura, K. Sakaoka and H. Tawara,
Computer Simulation and Data Compilation of Sputtering Yield by Hydrogen Isotopes (\(^{1}\text{H}^+, \^{2}\text{D}^+, \^{3}\text{T}^+) and Helium (\(^{4}\text{He}^+) Ion Impact from Monoatomic Solids at Normal Incidence; Oct. 1995
NIFS-DATA-32  T. Kato, U. Safronova and M. Ohira,
*Dielectronic Recombination Rate Coefficients to the Excited States of CII from CIII*; Feb. 1996

NIFS-DATA-33  K.J. Snowdon and H. Tawara,

NIFS-DATA-34  T. Ono, T. Kawamura, K. Ishii and Y. Yamamura
*Sputtering Yield Formula for B4C Irradiated with Monoenergetic Ions at Normal Incidence*; Apr. 1996