

**NATIONAL INSTITUTE FOR FUSION SCIENCE****Dielectronic Recombination of Be-like Fe Ion****K. Moribayashi and T. Kato**

(Received - Jan. 16, 1996)

NIFS-DATA-36

Apr. 1996

**RESEARCH REPORT  
NIFS-DATA Series**

This report was prepared as a preprint of compilation of evaluated atomic, molecular, plasma-wall interaction, or nuclear data for fusion research, performed as a collaboration research of the Data and Planning Center, the National Institute for Fusion Science (NIFS) of Japan. This document is intended for future publication in a journal or data book after some rearrangements of its contents.

Inquiries about copyright and reproduction should be addressed to the Research Information Center, National Institute for Fusion Science, Nagoya 464-01, Japan.

# Dielectronic recombination of Be-like Fe ion

Kengo Moribayashi and Takako Kato

National Institute for Fusion Science, Chikusaku Nagoya 464, Japan

## Abstract

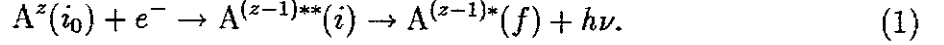
Energy level( $E$ ), radiative transition probability( $Ar$ ), and autoionization rate( $Aa$ ) for Be-like  $\text{Fe}^{22+}$  ion are calculated with use of Cowan's code. Using these atomic data, the dielectronic recombination rate coefficient( $\alpha$ ) to the excited states and the intensity factor( $Qd$ ) of the dielectronic satellite lines have been calculated. The doubly excited states  $1s^23lnl'$  as well as the  $1s^22pnl$  of  $\text{Fe}^{22+}$  ion are considered. The results are given in tables and figures. The  $n$ - and  $l$ -dependence for  $Ar$ ,  $Aa$ , and  $\alpha$  is studied. With use of it,  $Aa$  and  $Ar$  at large  $n$  are extrapolated. The dielectronic recombination processes from the  $1s^22pnl$  and those from the  $1s^23lnl'$  dominate at low and at high temperature, respectively. The qualitative different behaviors for  $E$ ,  $Ar$ , and  $\alpha$  between Be-like ions and He-like ions are discussed with use of atomic nuclear charge scaling.

## Keywords

dielectronic recombination rate coefficient, Beryllium-like ion, Fe XXIII ion, energy level, radiative transition probability, autoionization rate, atomic nuclear charge scaling, autoionization state, Cowan's code, dielectronic satellite spectra

## 1. Introduction

The dielectronic recombination(DR) process plays an important role in high temperature plasma such as tokamak plasma and solar corona[1, 2]. The process is composed by the two step transitions, the dielectronic capture and the radiative decay,



Here  $A^{(z-1)**}$  expresses the autoionization state of  $A^{(z-1)}$  ions. Dielectronic recombination rate coefficient( $\alpha$ ) from the initial state ( $i_0$ ) to the final state ( $f$ ) in Eq.(1), which provides us the population of the excited states, is represented by

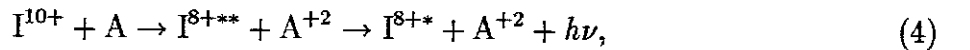
$$\alpha(i_0, f) = \frac{C(T_e)}{g_0} \sum_i \frac{g(i)Aa(i, i_0)Ar(i \rightarrow f)}{\sum_{\text{all } i'_0} Aa(i, i'_0) + \sum_{\text{all } f'} Ar(i \rightarrow f')} \exp\left(-\frac{\Delta E_i}{kT_e}\right) \quad (2)$$

with

$$C(T_e) = \frac{1}{2} \left( \frac{h^2}{2\pi m k T_e} \right)^{3/2}. \quad (3)$$

Here  $Ar$ ,  $Aa$ ,  $T_e$ ,  $g_0$  and  $\Delta E_i$ ; represent radiative transition probability, autoionization rate, electron temperature, statistical weight of the initial state ( $i_0$ ), and the energy difference between the autoionization state ( $i$ ) and the initial state ( $i_0$ ), respectively. Therefore in order to estimate  $\alpha$ , the atomic data such as the energy levels,  $Ar$  and  $Aa$  are required.

The  $Ar$  and  $Aa$  values are also useful for understanding double-electron charge transfer processes by ion-atom collision as well as the DR processes. Recently the double-electron charge transfer processes between  $I^{10+}$  ion and multi-electron atoms (A) such as Ne, Ar, Kr, Xe were measured[3, 4, 5]. The following process,



is similar to the DR process as in Eq.(1). The second step is the radiative transition process, which is the same as that in the DR process.

The spectral lines of  $L$ -shell ions have been often measured (eg.[6, 7]) and the dielectronic recombination of the  $Fe^{22+}$  ion have been often studied[8-11]. However, as long as

we know, there are few papers which give us enough data of the dielectronic recombination rate coefficient to each final bound state. These data are necessary to estimate the population of the excited states by a collisional radiative model[12]. Our purpose is to calculate the data for the dielectronic recombination to the excited states of  $\text{Fe}^{22+}$  ion.

We treat the  $1s^23l_1nl_2$  states, which are important for the DR process around the 1 keV temperature, as well as the  $1s^22l_1nl_2$  states. Here  $l_1(l_2)$  expresses the angular momentum of inner (outer) electron. The calculations are executed with the use of Cowan's code[13, 14] in which configuration interaction(CI) method is employed. Full number of configurations must be considered for small principal quantum number  $n$ . However, as the value of  $n$  increases, the number of configuration can be restricted because electron correlation becomes weaker. After we test the convergence of  $Ar$  and  $Aa$  for the number of the states, we selected the configurations necessary for the calculation. We calculated the  $Ar$  and  $Aa$  values up to  $n = 20$ . Further extrapolation is used for  $n > 20$ .

The aims of this paper are (i) to estimate  $\alpha$  for each final state as well as the intensity of the dielectronic satellite spectra, and (ii) to understand the atomic nuclear charge dependence as will be mentioned in Sec.2 for  $Ar$ ,  $Aa$ , and  $\alpha$ .

## 2. Atomic nuclear charge scaling for radiative transition probability and autoionization rate

In this section, we describe  $Z$ -scaling for  $Ar$  and  $Aa$  values following the discussion in Ref.[15], where  $Z$  represents the atomic effective nuclear charge of ions, in order to understand the behavior of  $\text{Fe}^{22+}$  ion and the phenomena of the dielectronic recombination more clearly. The  $Z$  value is the quantity removed the number of the inner electrons from the real atomic nuclear charge. For example, in the case of  $1s^22pnl$  state in  $\text{Fe}^{22+}$  ion, it is 24 for the  $2p$  electron and 23 for the  $nl$  electron, respectively.

$Ar$  for the process from state A to state B is given by

$$Ar = N\sigma^3S. \quad (5)$$

Here  $N$  and  $\sigma$  are the normalization constant and the energy difference between the states A and B, respectively and  $S$  is the line-strength given by

$$S = | \langle \psi_B | \vec{r} | \psi_A \rangle |^2, \quad (6)$$

where  $\psi_{A(B)}$  is the wave function of the state A(B). The  $Z$  value of  $\text{Fe}^{22+}$  ion is so large that the independent particle model is suitable for them in discussing qualitatively. Firstly  $S$  decreases according to  $Z^{-2}$  because the atomic radius  $\langle r \rangle \propto Z^{-1}$ . On the other hand, in order to estimate the  $Z$ -scaling for  $\sigma$ ,  $Z$ -expansion method[16] is useful. With the use of this method, the energy for large  $Z$  can be written as

$$E \cong Z^2(E_0 + E_1/Z), \quad (7)$$

because second and over order expansion can be neglected for large  $Z$ . Here  $E_{0(1)}$  is the 0(1)-th expansion coefficient. The  $E_0$  values are the same for the state A and the state B in the case when the principal quantum numbers  $n$  are the same between the states A and B. Therefore  $\sigma$  can be approximated by

$$\begin{aligned} \sigma &\propto Z^2 && \text{for } \Delta n \neq 0 \\ \sigma &\propto Z && \text{for } \Delta n = 0, \end{aligned} \quad (8)$$

where  $\Delta n$  expresses the difference between the principal quantum number of the transition electron in the state A and that in the state B. From Eqs.(5) and (8),  $Ar$  can be given by

$$\begin{aligned} Ar &\propto Z^4 && \text{for } \Delta n \neq 0 \\ Ar &\propto Z && \text{for } \Delta n = 0. \end{aligned} \quad (9)$$

Then,  $Aa$  is given by

$$Aa \sim \frac{2\pi}{\hbar} | \langle \Psi_{j_1} \Psi_{j_2} | \frac{1}{r_{12}} | \Psi_{j_3} \Psi_i \rangle |^2, \quad (10)$$

where the states  $j$  and  $i$  express discrete and continuum states, respectively and  $r_{12}$  is the distance between two electrons. In Cowan's code[13, 14], a free electron state is assumed to be separable from the bound state. Then the wave function of a free electron, which is of the same form as that for a bound electron, is given by

$$\Psi_{i, \ell m_\ell m_s}(\vec{r}) = \frac{1}{r} P_{\ell l}(r) Y_{\ell m_\ell}(\theta, \phi) \Sigma_{m_s}(s_z). \quad (11)$$

with

$$P_{el} \sim \left(\frac{2}{\pi p_c}\right)^{1/2} \frac{1}{r} \sin(p_c r + \delta(r)). \quad (12)$$

Here  $p_c$  equals to  $\epsilon^{1/2}$  where  $\epsilon$  is the free electron energy. That is,  $p_c$  increases according to

$$\begin{aligned} p_c &\propto Z & \text{for } \Delta n_c \neq 0 \\ p_c &\propto Z^{1/2} & \text{for } \Delta n_c = 0. \end{aligned} \quad (13)$$

Namely,

$$\begin{aligned} \Psi_i &\propto Z^{1/2} & \text{for } \Delta n_c \neq 0 \\ \Psi_i &\propto Z^{3/4} & \text{for } \Delta n_c = 0 \end{aligned} \quad (14)$$

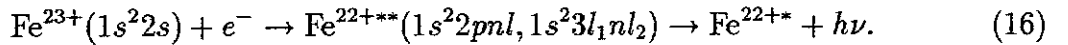
because  $\langle r^{-1} \rangle \propto Z$ . Then  $\Psi_j$  and  $\langle r_{12}^{-1} \rangle$  increase according to  $Z^{3/2}$  and  $Z$ , respectively and  $dr_1^3 dr_2^3$  decreases according to  $Z^{-6}$ . Therefore  $Aa$  becomes

$$\begin{aligned} Aa &\propto Z^0 & \text{for } \Delta n_c \neq 0 \\ Aa &\propto Z^{1/2} & \text{for } \Delta n_c = 0. \end{aligned} \quad (15)$$

Here  $\Delta n_c$  expresses the difference of the principal quantum number between the states  $j_1$  and  $j_3$ . After all it is found that the  $Z$ -dependence for  $Aa$  is weaker than that for  $Ar$ .

### 3. Atomic data :Energy level, radiative transition probability, and autoionization rate

We investigate the following dielectronic recombination process of  $\text{Fe}^{22+}$  ion,



In this section, we present the atomic data of  $Aa$  and  $Ar$  calculated with use of Cowan's code[13, 14]. We use the extrapolation formula of  $Ar$  and  $Aa$  for  $n > 20$  based on  $n$ -dependence for  $Ar$  and  $Aa$  in order to save the CPU time and memory. It is noticed that the autoionization states at large  $n$  contribute to the dielectronic recombination significantly at high temperature plasma.

Table I lists the energy levels and weighted transition probabilities summed over all the final states ( $\sum gAr$ ) for the  $1s^2 2snl$  and  $1s^2 2pnl$  ( $n = 2 \sim 5$ ) states of  $\text{Fe}^{22+}$  ion. Here  $g$  is the statistical weight of the upper state. The energy levels given by Shirai et.al.[17] are also listed. For  $n = 2$  and 3, the differences between the present energy level and the others are less than 5%. On the other hand, for  $n = 4$  and 5, the agreement becomes better, that is, the errors are less than 1%. In our table, the two states with the same designation exist, for example,  $2p3d \ ^3P_2$  ( $\ ^3D_2$ ). This is because the  $2p3d \ ^3P_2$  state mixes with  $2p3d \ ^3D_2$  one strongly. Therefore this level can be named as  $2p3d \ ^3D_2$  as well.

Table II lists  $Aa$  and  $Ar$  values for  $n = 11, 13,$  and  $20$  levels. The ionization energy from the ground state  $1s^2 2s^2$  of  $\text{Fe}^{22+}$  ion is about  $15730000 \text{ cm}^{-1}$  ( $1950\text{eV}$ ). The autoionization takes place only when the energy for the doubly excited states is above the threshold energy. For the  $1s^2 2pnl$  ( $n < 11$ ) states, the energies are less than that of the threshold. Therefore,  $1s^2 2pnl$  ( $n < 11$ ) states are not autoionization states although these states are the doubly excited states. The energy difference between the threshold  $1s^2 2s$  and the  $1s^2 2snl$  states increases according to  $Z^2$  because the threshold corresponds to the  $1s^2 2snl$  ( $n \rightarrow \infty$ ). On the other hand, the energy difference between the  $1s^2 2snl$  and  $1s^2 2pnl$  is proportional to  $Z$  as mentioned in Eq.(8) in Sec.2. Therefore, the  $\Delta E/I_p$  value for  $1s^2 2pnl$  becomes smaller as  $Z$  increases where  $I_p$  is the ionization potential of  $1s^2 2s^2$ . Consequently for the ions with larger  $Z$ , the  $1s^2 2pnl$  state at the larger  $n$  values locates below the threshold. It should be noted that this feature is not seen in He-like and Li-like ions since there is no  $\Delta n=0$  transition. In these ions, the threshold corresponds to the  $1s$  or  $1s^2$  states. Then, the transition for dielectronic recombination is only the  $np \rightarrow 1s$ , whose energy difference increases according to  $Z^2$ . Therefore,  $\Delta E$  increases according to  $Z^2$ . In the doubly excited  $1s^2 3l_1 n l_2$  states,  $\Delta E$  increases according to  $Z^2$  because only the  $3l_1 \rightarrow 2l'$  ( $\Delta n \neq 0$ ) transition takes place.

For the  $1s^2 2pnl$  ( $n = 11$ ) states, as is seen in Table II, both autoionization states and bound states exist. This is due to the fact that the two thresholds,  $1s^2 2p \ ^2P_{1/2}$  and  $1s^2 2p \ ^2P_{3/2}$  exist. As  $n$  increases, the  $1s^2 2pnl$  states reach either  $1s^2 2p \ ^2P_{1/2}$  or  $1s^2 2p \ ^2P_{3/2}$  state. The energy difference between them is about  $130000\text{cm}^{-1}$ , which is almost the same

as that between the autoionization states and the bound states for  $1s^22p11l$ . From  $n = 13$ , all of the energies begin to be above the threshold, although the  $\Delta E$  values are very small (60000 and 185000  $\text{cm}^{-1}$ ). At  $n = 20$ , the  $\Delta E$  values become about 260000  $\text{cm}^{-1}$  (32.2eV) and 385000  $\text{cm}^{-1}$  (47.7 eV). Although they are far below the second threshold of  $1s^22p$  states whose energy is 392000  $\text{cm}^{-1}$  (48.6 eV) for  $1s^22p^2P_{1/2}$  state and 520000 $\text{cm}^{-1}$  (64.5 eV) for the  $1s^22p^2P_{3/2}$  state measured from the  $1s^22s$  state.

Table III lists  $Ar$  from the  $1s^22snl$  and  $1s^22pnl$  ( $n \leq 5$ ) to the  $1s^22s^2$ ,  $1s^22p^2$ , and  $1s^22s2p$  states. The  $Ar$  values given by Shirai et.al.[17] are also shown. Our results agree very well with them within 10 % except for a few transitions where  $n$  is small values, that is,  $n=2$  or 3. In the table, the radiative transitions with large values take place mainly for the 4 cases,  $\{\Delta J = \Delta L = \Delta l_1$  (or  $\Delta l_2$ )=1 $\}$ ,  $\{\Delta J = 0, \Delta L = \Delta l_1$  (or  $\Delta l_2$ )=1 $\}$ ,  $\{\Delta L = 0, \Delta J = \Delta l_1$  (or  $\Delta l_2$ )=1 $\}$ , and  $\{\Delta J = \Delta L = 0, \Delta l_1$  (or  $\Delta l_2$ )=1 $\}$ . Here  $\Delta X$  expresses the difference of  $X$  values between the lower state and upper states. For the 4 cases,  $\Delta X=-1$  dose not exist. The plus sign and 0 of  $\Delta X$  mean that the  $X$  value of upper state is larger than or equals to that of the lower state.

Figs.1 show  $gAr$  from the  $1s^22pnl$  of  $\text{Fe}^{22+}$  ion as a function of  $n$ . Here the  $Ar$  values are given by

$$gAr = \sum_{S'L'J'} \sum_{SLJ} (2J+1) Ar(1s^22pnl^{2S+1}L_J \rightarrow 1s^22l''n'l'^{2S'+1}L'_{J'}). \quad (17)$$

The  $Ar$  values for the transition from  $1s^22pnl$  to the low excited states such as  $1s^22p2p$ ,  $1s^22p3p$ , are more than or comparable to those for the  $1s^22pnl \rightarrow 1s^22snl$  transitions even at  $n = 20$ . In the case of small  $Z$  such as  $\text{C}^{2+}$  ion, the latter is much larger than the former at  $n = 20$ [19]. This comes from the fact that the former and the latter increase according to  $Z^4$  and  $Z$  as mentioned in Eq.(9) in Sec.2, respectively. In the case of  $\text{Fe}^{22+}$  ion, we can use the following extrapolation for  $n > 20$  because the  $Ar$  values for the  $1s^22pnl \rightarrow 1s^22snl$  and those for the  $1s^22pnl \rightarrow 1s^22pn'l'$  processes keep constant and decrease according to  $n^{-3}$  as  $n$  increases, respectively.

$$Ar(1s^22pnl^S L_J \rightarrow 1s^22snl^{S'} L'_{J'}) \cong Ar(1s^22p20l^S L_J \rightarrow 1s^22s20l^{S'} L'_{J'}) \quad (18)$$



and

$$Ar(1s^2 2pnl \ ^S L_J \rightarrow 1s^2 2pn'l' \ ^{S'} L'_J) \cong Ar(1s^2 2p20l \ ^S L_J \rightarrow 1s^2 2pn'l' \ ^{S'} L'_J) \times (20/n)^3. \quad (19)$$

Fig.2 shows  $Aa$  values given by

$$Aa = \sum_{SLJ} Aa(1s^2 2pnl \ ^S L_J, i_0) \quad (20)$$

from the  $1s^2 2pnl$  ( $n \geq 13$ ) of  $\text{Fe}^{22+}$  as a function of  $n$ . As mentioned before, since the energies of the  $1s^2 2pnl$  ( $n < 13$ ) are below the ionization threshold,  $Aa$  is plotted only for  $n \geq 13$  in Fig.2. The  $Aa$  values for  $n > 20$  are also estimated by scaling the values at  $n = 20$  as

$$Aa(1s^2 2pnl \ ^S L_J, i_0) \cong Aa(1s^2 2p20l \ ^S L_J, i_0) \times (20/n)^3 \quad (21)$$

because the  $Aa$  values are expected to decrease according to  $n^{-3}$ .

Figs.3 show the  $Ar$  values, the same form as in Fig.1 from the  $1s^2 3l_1 n l_2$  to the  $1s^2 2l'_1 n' l'_2$  state. Even for the small value of  $n$ , the  $Ar$  values for the transitions  $1s^2 3snl \rightarrow 1s^2 2pnl$  ( $n < 11$ ),  $1s^2 3pnl \rightarrow 1s^2 2snl$ , and  $1s^2 3dnl \rightarrow 1s^2 2pnl$  ( $n < 11$ ) are much larger than those for the other transitions such as  $1s^2 3pnl \rightarrow 1s^2 2l'3l''$ . The  $Ar$  values for all the transitions from the  $1s^2 3l_1 n l_2$  state to the bound states increase according to  $Z^4$  because of the  $3l_1 \rightarrow 2l'_1$  transitions. As a result, these  $Ar$  values for the  $1s^2 3l_1 n l_2 \rightarrow 1s^2 2l'_1 n l'_2$ , which keep constant for  $n$ , are much larger than those for the  $1s^2 3l_1 n l_2 \rightarrow 1s^2 2l'_1 n' l'_2$  at large  $n$ .

Figs.4 show the  $Aa$  values of the  $1s^2 3l_1 n l_2$ . For  $1s^2 3l_1 n l_2$  states, there are three autoionization processes to  $1s^2 2scl$ ,  $1s^2 2pcl$ , and  $1s^2 3scl$  where  $cl$  represents the free electron. However, since the energy for the  $1s^2 3l_1 n l_2$  at  $n = 20$  is below the threshold  $1s^2 3s$  state, the processes to the  $1s^2 3scl$  are ignored. The  $Aa$  value from the  $1s^2 3snl$  to the  $1s^2 2s$  is more than or comparable to that to the  $1s^2 2p$ . On the other hand, in the processes from the  $1s^2 3pnl$  and  $1s^2 3dnl$ , the  $Aa$  values to  $1s^2 2p$  are much larger than those to  $1s^2 2s$ . Since the  $Aa$  values are smaller in comparison with  $Ar$  for  $1s^2 3l_1 n l_2$  states for  $n > 20$ , the processes for  $n > 20$  give little contribution to dielectronic recombination.

Fig.5 shows  $l$ -distribution of the  $Ar$  value for  $1s^2 2pnl \rightarrow 1s^2 2l''n'l'$  given by

$$Ar = \sum_{n'l''} \sum_{S'L'J'} \sum_{SLJ} Ar(1s^2 2pnl \ ^S L_J \rightarrow 1s^2 2l''n'l' \ ^S L'_{J'}). \quad (22)$$

for  $n=11, 13, 20$ . Each curve for  $n=11, 13,$  and  $20$  shows the similar features. Namely the  $Ar$  values have a peak at  $l = 2$  and decrease gradually for large  $l$  ( $l \geq 4$ ) values. The transitions  $1s^2 2pnl \rightarrow 1s^2 2pn'l'$  give a main contribution to the dielectronic recombination at low temperature and the  $Ar$  values for these transitions decrease according to  $n^{-3}$ . The  $Ar$  values for the  $1s^2 2pnl \rightarrow 1s^2 2snl$  keep constant for both  $l$  and  $n$  as is seen in Ref.[18, 19]. Fig.6 shows the  $Aa$  values given in Eq.(20) for  $n=15$  and  $20$  as a function of  $l$ . We find a peak at  $l = 1$  and a dip at  $l=3$ . For  $l \geq 5$ , the  $Aa$  values decrease rapidly. Since  $\alpha$  is proportional to  $Aa$  in the case of  $Ar \gg Aa$ , we can neglect the contribution from large  $l$  ( $l \geq 7$ ) for high  $n$  states.

#### 4. Dielectronic satellite spectra

The dielectronic satellite lines (e.g.  $2snl - 3pnl$ ) of Be-like ions appear near the excitation lines (e.g.  $2s - 3p$ ) of Li-like ions. Table IV shows wavelength ( $\lambda$ ), radiative transition probabilities summed over all the final states ( $\sum gAr$ ), autoionization rate ( $Aa$ ), and the intensity factor ( $Qd$ ) for the dielectronic satellite lines given by

$$Qd(i_0, i, f) = \frac{g(i)Aa(i, i_0)Ar(i \rightarrow 2l''n'l' \ ^S L'_{J'})}{\sum_{all \ i'_0} Aa(i, i'_0) + \sum_{all \ f'} Ar(i \rightarrow f')}. \quad (23)$$

Here the  $1s^2 2p11l, 1s^2 2p13l, 1s^2 3l_1 3l_2,$  and  $1s^2 3l_1 4l_2$  states are selected as autoionization states ( $i$ ) and the  $\lambda$  values are taken from 6 to  $1000\text{\AA}$ . From this table, we found that the  $Qd$  values become large in the same case where the radiative transitions with large value take place, that is,  $\{\Delta J = \Delta L = \Delta l_1 \text{ (or } \Delta l_2)=1\}, \{\Delta J = 0, \Delta L = \Delta l_1 \text{ (or } \Delta l_2)=1\}, \{\Delta L = 0, \Delta J = \Delta l_1 \text{ (or } \Delta l_2)=1\},$  and  $\{\Delta J = \Delta L = 0, \Delta l_1 \text{ (or } \Delta l_2)=1\}$  as listed in Table III in Sec.3. The  $Qd$  values in the case of  $\text{Fe}^{22+}$  ion are much larger than those in  $\text{C}^{2+}$  ion[18]. The  $Qd$  values in the  $\text{C}^{2+}$  ion are less than  $10^{11}(1/s)$ , on the other hand, those in the  $\text{Fe}^{22+}$  ion reach a maximum value of  $10^{13}(1/s)$ . This means that the dielectronic satellite spectra is more important in the  $\text{Fe}^{22+}$  ion. In particular, the  $Qd$  values from the

$3l - 2l'$  transitions around  $11\text{\AA}$  are larger.

With use of the  $Qd$  values, roughly speaking, satellite spectral intensity for the transition can be estimated as the order of  $10^{-12}\text{cm}^3/\text{s}$  at  $T_e=300\text{eV}$  and that of  $10^{-13}\text{cm}^3/\text{s}$  at  $T_e=1\text{keV}$ , respectively. On the other hand, the excitation rate coefficient of  $2s-3d$  and  $2s-3p$  lines in Li-like  $\text{Fe}^{23+}$  ion which provide us to estimate the line intensity are about  $10^{-13}\text{cm}^3/\text{s}$  and  $10^{-11}\text{cm}^3/\text{s}$  at  $T_e=300\text{eV}$  and  $1\text{keV}$ , respectively[20]. Then the excitation rate coefficients of  $2s-3l$  lines in Be-like  $\text{Fe}^{22+}$  ion are also about  $10^{-13}\text{cm}^3/\text{s}$  and  $10^{-11}\text{cm}^3/\text{s}$  at  $T_e=300\text{eV}$  and  $1\text{keV}$ , respectively[12]. Namely the dielectronic satellite spectra is important at only the low temperature below  $T_e=300\text{eV}$ .

## 5. Dielectronic recombination rate coefficient to the excited states

The DR rate coefficient given in Eq.(2) is important to estimate the population densities of the excited states through the dielectronic recombination processes. We summed up the transitions  $i \rightarrow f$  in Eq.(2) for the autoionizing states  $i$  ( $1s^2n_1l_1n_2l_2 \ 2^{S+1}L_J$ ). With the use of the numerical value  $C(T_e) = 3.3 \times 10^{-24}(I_H/T_e)^{3/2}$  and  $g_0 = 2$  in Eq.(2), we obtain the DR rate coefficient to the excited state  $1s^22l'_1n'l'_2$  as

$$\alpha_d(1s^22s, 1s^22l'_1n'l'_2 \ 2^{S+1}L'_{J'}) = \frac{3.3 \times 10^{-24}}{2} \left(\frac{I_H}{T_e}\right)^{3/2} \sum_{n_1} \sum_{n_2} \sum_{l_2} \sum_{l_1LSJ} Q_d(1s^22s, 1s^2n_1l_1n_2l_2 \ 2^{S+1}L_J, 1s^22l'_1n'l'_2 \ 2^{S+1}L'_{J'}) \exp\left(-\frac{\Delta E_i}{kT_e}\right) \quad (24)$$

where  $I_H$  is the ionization potential of a hydrogen atom and  $n_1$  ( $n_2$ ) is the principal quantum number of the inner (outer) electron.

For the  $\Delta E_n$  value of  $1s^22pnl$  for  $n > 20$ , the following extrapolation equation is used,

$$\Delta E_n = \Delta E_\infty - \frac{a}{n^2}, \quad (25)$$

with

$$a = 20^2 \times (\Delta E_\infty - \Delta E_{20}) \quad (26)$$

because the energy at  $n = 20$  is far from the threshold. Here  $\Delta E_\infty$  ( $=48.6$  or  $64.5\text{eV}$ ) and  $\Delta E_{20}$  ( $=32.2$  or  $47.7\text{eV}$ ) represent the  $\Delta E_i$  value of the threshold  $1s^22p$  and that at

the  $1s^2 2p 20l$  states, respectively. The detailed explanation about  $\Delta E$  was done in Sec.3.

In order to obtain the values  $\alpha_d$  to the excited state  $1s^2 2l'_1 n' l'_2$ , two kind of transitions,  $1s^2 2pnl \rightarrow 1s^2 2l'_1 n' l'_2$  and  $1s^2 3l_1 n l_2 \rightarrow 1s^2 2l'_1 n l_2$  are taken into account in our calculation. For the  $1s^2 2pnl \rightarrow 1s^2 2l'_1 n' l'_2$  transition, the sum in Eq.(24) up to  $n=11-500$ ,  $l=0-8$  and for the  $1s^2 3l_1 n l_2 \rightarrow 1s^2 2l'_1 n' l'_2$  transition, the sum up to  $n = 3-20$ ,  $l = 0-5$  are taken, respectively. The  $\alpha_d$  values for each excited state are given in Table V.

## 6. Total dielectronic recombination rate coefficients

The total DR rate coefficient is obtained by the sum of all the levels,

$$\alpha_d^t = \sum_{n'=2}^{20} \alpha_{n'} + \sum_{n'=21}^{500} \alpha_{n'}, \quad (27)$$

where

$$\alpha_{n'} = \sum_{l'_1 l'_2} \sum_{S' L' J'} \alpha_d(1s^2 2s, 1s^2 2l'_1 l'_2 \quad {}^{2S'+1} L'_{J'}). \quad (28)$$

The second term in the right hand side of Eq.(27) is calculated by extrapolation as discussed in Sec.3 using the values at  $n = 20$ .

In Fig.7, the total dielectronic recombination rate coefficient is shown with other theoretical results[8, 9, 11]. Our results are smaller than those by McLaughlin & Hahn[8] and Romanic[11] by 20% and larger than those by Roszman's[9] by 20% at 1keV. The processes from the  $1s^2 2pnl$  and those from  $1s^2 3l_1 n l_2$  dominate at low temperature( $T_e < 200\text{eV}$ ) and high temperature( $T_e > 500\text{eV}$ ), respectively. The small hollow is seen around  $T_e = 300\text{eV}$  just where the processes from the  $1s^2 2pnl$  contribute comparable to those from the  $1s^2 3l_1 n l_2$ . In the  $\text{C}^{2+}$  ion, all of the  $\alpha$  values from the  $1s^2 3l_1 n l_2$  are always smaller than those from the  $1s^2 2pnl$  states[18]. The temperature where  $\alpha$  has the maximum is determined by the temperature near  $\Delta E/T_e \sim 1$ . As mentioned before, the  $\Delta E/I_p$  values for the  $1s^2 2pnl$  and those for the  $1s^2 3l_1 n l_2$  become smaller and keep constant as  $Z$  increase, respectively. Namely the difference of the peak position of  $\alpha$  values in  $T_e$  for the process from the  $1s^2 2pnl$  and that from the  $1s^2 3l_1 n l_2$  increases as  $Z$  increases. Therefore, in  $\text{Fe}^{22+}$  ion which has large  $Z$  value, the process from the  $1s^2 3l_1 n l_2$  states dominates at high  $T_e$ . For the He-like and Li-like ion, since all of the  $\Delta E$  values increase according to

$Z^2$ , all of the peaks shift according to  $Z^2$ . As a result, almost the similar figures are seen between He-like  $C^{4+}$  and  $Fe^{24+}$  ions[11].

Fig.8 shows  $\alpha_{n'}$  given by Eq.(28) for  $n'=2\sim 6$  as a function of  $T_e$ . At low  $T_e$ , the processes to  $n' = 2$  or 3 dominate. On the other hand, as  $T_e$  increases, the processes to  $n' = 4 \sim 6$  become comparable to that to  $n' = 3$ . Further the  $\alpha$  value to  $n' = 2$  decreases rapidly. This comes from the fact that the processes from  $1s^23l_1n'l_2$  to  $n' = 2$  contribute little to dielectronic recombination.

Figs.9(a)-(d) show  $f_\alpha$  as a function of  $n'$  which is principal quantum number of the final bound state. Here  $f_\alpha$  gives the quantity summed up  $Qd$  over the autoionization  $n$  states, that is,

$$f_\alpha(n') = \sum_{n,l,L,S,J} \sum_{n',l',L',S',J'} Qd(i_0, nl^{2S+1}L_J, 2l'n'l'^{2S'+1}L'_{J'}). \quad (29)$$

With the use of the  $f_\alpha(n')$ , the  $n'$ -dependence for  $\alpha_{n'}$  can be understood roughly, though the  $f_\alpha$  is independent of  $T_e$ . For the  $n' < 11$  states, the contribution from transitions  $1s^22pnl \rightarrow 1s^22pn'l'$  is significant, since  $Ar$  values for these transitions are much larger than that of the  $1s^22pnl \rightarrow 1s^22snl$  transition. However, for  $n' > 11$ , the  $1s^22pn'l'$  states are autoionizing states and only the  $1s^22pnl \rightarrow 1s^22snl$  transition remains. Therefore, big gaps around  $n' = 11$  are found in Fig.9(a). For  $11 \leq n' \leq 100$ ,  $f_\alpha$  keeps constant. On the other hand, for  $n' \geq 200$ ,  $f_\alpha$  decreases according to  $n^{-3}$ . This is due to the fact that  $f_\alpha$  is proportional to  $Ar$  for  $11 \leq n' \leq 100$  and  $Aa$  for  $n' \geq 200$  because of  $Aa \gg Ar$  and  $Ar \gg Aa$ , respectively. For  $1s^23pnl_2-1s^22snl_2$  transition,  $f_\alpha$  decreases according to  $n'^{-3}$  after around  $n'=10$  as is seen from Fig.9(b). Furthermore, the  $f_\alpha$  values for  $n \geq 13$  equal to 0 in the  $1s^23snl_2-1s^22pn_2l$ , and  $1s^23dnl_2-1s^22pnl_2$  transitions (see Figs.9(c) and (d)), because  $1s^22pnl(n \geq 13)$  states are autoionization states. Therefore,  $f_\alpha$  for  $n > 20$  is ignored in our calculation for  $1s^23l_1n'l_2$  levels.

Fig.10 shows ratio of  $\sum_{n''=2}^{n'} \alpha_{n''}$  to  $\sum_{n''=2}^{500} \alpha_{n''}$  as a function of  $n'$ . Here the  $\alpha_{n''}$  values are given by Eq.(28). The ratios at  $n' = 10$  are about 0.7 and 0.85 for  $T_e=100\text{eV}$  and  $1\text{keV}$ , respectively. These values are much larger than that in the case of the  $C^{2+}$  ion. In the  $C^{2+}$  ion, they are less than 0.1[18].

## 6. Conclusion

We calculate the radiative transition probability ( $Ar$ ), autoionization rate ( $Aa$ ), the intensity factor of the dielectronic satellite spectra ( $Qd$ ) and the dielectronic recombination rate coefficient ( $\alpha_d$ ) of  $\text{Fe}^{22+}$  ion with use of Cowan's code[13, 14]. We treat not only the  $1s^22l_1nl_2$  states but also  $1s^23l_1nl_2$  states. We confirm the  $n$ -dependence for  $Ar$  and  $Aa$  which are expected theoretically. Namely the  $Ar$  values for the  $1s^22pnl \rightarrow 1s^22pn'l'(n' = 2 \sim 4)$  and  $1s^22pnl \rightarrow 1s^22snl$  processes decrease according to  $n^{-3}$  and keep constant as  $n$  increase, respectively. And the values  $Aa$  decrease according to  $n^{-3}$ . We used these dependence for the extrapolation for  $n > 20$ . The  $Ar$  values for the  $1s^22pnl \rightarrow 1s^22pn'l'(n' = 2 \sim 4)$  processes are much larger than those for the  $1s^22pnl \rightarrow 1s^22snl$  processes even at large  $n$  for  $\text{Fe}^{22+}$ . In the case of  $\text{C}^{2+}$ , the differences are not so large. This comes from the fact that the former and the latter increase according to  $Z^4$  and  $Z$ , respectively.

Around  $n = 20$ , the  $Aa$  value from the  $1s^22pnl$  levels is much larger than  $Ar$  values, while the  $Aa$  value from the  $1s^23l_1nl_2$  is much smaller than the  $Ar$  value. The  $Aa$  values decrease rapidly as  $l$  increases for large  $l$ . On the other hand, the  $Ar$  values keep almost constant for large  $l$ . Then we can neglect the dielectronic recombination processes with  $l \geq 7$  for large  $n$ .

The dielectronic satellite spectra is important only in the low temperature below  $T_e=300\text{eV}$ . As for the dielectronic recombination, at low  $T_e(\leq 200\text{eV})$ , the processes from the  $1s^22pnl$  states dominate. On the other hand, at high  $T_e(\geq 500\text{eV})$ , only the processes from the  $1s^23l_1nl_2$  give a significant contribute to the dielectronic recombination. There is a different qualitative feature between  $\alpha$  from the  $1s^22pnl$  states and that from the  $1s^23l_1nl_2$  states through  $Z$ -scaling. This comes from the fact that the  $\Delta E$  values for the former ( $\Delta n = 0$  transition) and those for the latter ( $\Delta n \neq 0$  transition) increase as  $Z$  and  $Z^2$  as  $Z$  increases, respectively. For He-like ions, only one peak is found since there are only  $\Delta n \neq 0$  transitions. Finally the  $\alpha$  values for each final bound state are calculated

and listed.

## Acknowledgment

We wish to thank Dr. U.I.Safronova for her useful discussions and guidance about the Cowan's code.

## References

- [1] A.Burgess, *Astrophys. J.*, **139** 776(1964), **141** 1588(1965).
- [2] M.J. Seaton and P.J. Storey, "Di-electronic recombination" in "atomic processes and applications" edited by P.G. Burke and B.L. Moisewitsch, North-Holland 1976 p.133.
- [3] I.Yamada, F.J. Currell, A.Danjo, M.Kimura, A.Matsumoto, N.Nakamura, S.Ohtani, H.A. Sakaue, M.Sakurai, H.Tawara, H.Watanabe, and M.Yoshino, *J.Phys.B*, **28** L9(1995).
- [4] M. Kimura, N. Nakamura, H. Watanabe, I. Yamada, A.Danjo, K.Hosaka, A.Matsumoto, S.Ohtani, H.A. Sakaue, M.Sakurai, H.Tawara and M.Yoshino, *J.Phys.B*, **28** L643(1995).
- [5] N.Nakamura, F.J. Currell. A.Danjo, M.Kimura, A.Mastumoto, S.Ohtani, H.A.Sakaue, M.Sakurai, H.Tawara, H.Watanabe, I.Yamada, and M.Yoshino, *J.Phys.B*, **28** 2959(1995).
- [6] D.L. McKenzie, P.B.Landecker, U.Feldman, and G.A. Doschek, *Astrophys. J.*, **289** 849(1985).

- [7] C.C. Klepper, J.T. Hogan, S.J. Tobin, R.C. Isler, D.Guilhem, W.R. Hess, and P.Monier-Garbet, *J.Nucl. Mate.*, **220-222** 521(1994).
- [8] D.J. McLaughlin and Y. Hahn, *Phys.Rev.A*, **29** 712(1984).
- [9] L.J. Roszman, *Phys.Rev.A*, **35** 2122(1987).
- [10] D.C. Griffin and M.S. Pindzola, *Phys.Rev.A*, **35** 2821(1987).
- [11] C.J. Romanik, *Astrophys. J.*, **330** 1022(1988).
- [12] I.Murakami and T.Kato, *Phsica Scripta*, to be published(1996).
- [13] R.D. Cowan, *J.Opt.Soc.Am.*, **58** 808(1968).
- [14] R.D. Cowan, "The theory of atomic structure and spectra", University of California Press, 1981.
- [15] Y.Hahn, *Adv.Atom and Mol. Phys.*, **21** 123(1985).
- [16] L.A. Vainshtein and U.I. Safronova, *Physica Scripta*, **31** 519(1985).
- [17] T.Shirai, Y.Funatake, K.Mori, J.Sugar, W.L.Wiese, and T.Nakai, *J. Phys. Chem. Ref. Data*, **19** 174(1990).
- [18] T.Kato, U.I.Safronova, and M.Ohira, submitted to NIFS-DATA(1996).
- [19] K.Moribayashi and T.Kato, in preparation for publication.
- [20] H.L. Zhang, D.H. Sampson, and C.J. Fontes, *Atom. Data and Nucl. Data Tables*, **31**, **44** (1990).



## Figure Captions

**Fig.1**  $gAr$  given in Eq.(17) from the  $1s^2 2pnl_2$  to the  $1s^2 ln'l'$  as a function of  $n$ : (a) $l_2 = s$ , (b) $l_2 = p$ , (c) $l_2 = d$ , and (d) $l_2 = f$ : The configurations indicated mean final states  $1s^2 2ln'l'$ .

**Fig.2**  $Aa$  given in Eq.(20) from the  $1s^2 2pnl_2$  as a function of  $n$ .

**Fig.3**  $gAr$  given in Eq.(17) from the  $1s^2 3l_1 nl_2$  to the  $1s^2 2l'_1 n'l_2$  as a function of  $n$ : (a) $l_1 = s, l_2 = p$ , (b) $l_1 = s, l_2 = d$ , (c) $l_1 = p, l_2 = p$ , (d) $l_1 = p, l_2 = d$ , (e) $l_1 = d, l_2 = p$ , and (f) $l_1 = d, l_2 = d$ .

**Fig.4**  $Aa$  given in Eq.(20) from the  $1s^2 3l_1 nl_2$  as a function of  $n$ : (a) $l_1 = s, l_2 = p$ , (b) $l_1 = s, l_2 = d$ , (c) $l_1 = p, l_2 = p$ , (d) $l_1 = p, l_2 = d$ , (e) $l_1 = d, l_2 = p$ , and (f) $l_1 = d, l_2 = d$ .

**Fig.5**  $Ar$  given in Eq.(22) from the  $2pnl_2$  as a function of  $l_2$ .

**Fig.6**  $Aa$  given by Eq.(20) from the  $2p20l_2$  as a function of  $l_2$ .

**Fig.7**  $\alpha_2^t$  given by Eq.(27) as a function of  $T_e$ . Solid line: present result, dotted line: Romanik[11] dashed-dot line: Roszman[9], and \*: Mc & Hahn[8].

**Fig.8**  $\alpha_{n'}$  given by Eq.(28) to the final  $n'$  states as a function of  $T_e$ .

**Fig.9**  $f_\alpha$  given in Eq.(29) as a function as  $n'$ :the processes for (a) $1s^2 2pnl \rightarrow 1s^2 2l_1 n'l'$ , (b) $1s^2 3pnl \rightarrow 1s^2 2sn'l'$ , (c) $1s^2 3snl \rightarrow 1s^2 2pn'l'$ , and (d) $1s^2 3dnl \rightarrow 1s^2 2pn'l'$ .

**Fig.10** Ratio of  $\sum_{n''=2}^{n'} \alpha_{n''}$  to  $\sum_{n''=2}^{\infty} \alpha_{n''}$  as a function of  $n'$ .

**Table I** Comparison between our energy level (in units of  $1000\text{cm}^{-1}$ ) for  $1s^2 2l_1 n l_2$  ( $n = 2 \sim 5$ ) and those by others[17]. The first to third columns give the state, present results, and the other results.  $\sum gAr$  (in units of  $1/s$ ) is also give in fourth column.

state	energy ( $1000\text{cm}^{-1}$ )		$\sum gAr(1/s)$
	present work	Shirai et.al.[17]	
$1s^2 2s^2 \ ^1S_0$	0.0000		0.0000D+00
$1s^2 2s2p \ ^3P_0$	347.3849		0.0000D+00
$1s^2 2s2p \ ^3P_1$	378.2174	379.130	0.1882D+09
$1s^2 2s2p \ ^3P_2$	469.4410	471.780	0.0000D+00
$1s^2 2s2p \ ^1P_1$	723.9891	752.840	0.5403D+11
$1s^2 2p^2 \ ^3P_0$	952.8893	956.100	0.1235D+11
$1s^2 2p^2 \ ^3P_1$	1023.7057	1027.200	0.4541D+11
$1s^2 2p^2 \ ^3P_2$	1068.6712	1071.700	0.6588D+11
$1s^2 2p^2 \ ^1D_2$	1195.4688	1204.420	0.5322D+11
$1s^2 2p^2 \ ^1S_0$	1389.0091	1423.000	0.2999D+11
$1s^2 2s3s \ ^3S_1$	8906.0916		0.1109D+14
$1s^2 2s3s \ ^1S_0$	8967.8524		0.1341D+13
$1s^2 2s3p \ ^3P_0$	9071.1666		0.4021D+11
$1s^2 2s3p \ ^3P_1$	9071.8007		0.1395D+14
$1s^2 2s3p \ ^1P_1$	9103.9508	9107.000	0.2719D+14
$1s^2 2s3p \ ^3P_2$	9105.7396		0.3593D+12
$1s^2 2s3d \ ^3D_1$	9194.9285	9199.000	0.7122D+14
$1s^2 2s3d \ ^3D_2$	9199.5042	9209.000	0.1156D+15
$1s^2 2s3d \ ^3D_3$	9206.9365	9212.000	0.1560D+15
$1s^2 2s3d \ ^1D_2$	9264.5227	9273.000	0.8418D+14
$1s^2 2p3s \ ^3P_0$	9344.1329	9295.000	0.2404D+13
$1s^2 2p3s \ ^3P_1$	9360.1533		0.8793D+13
$1s^2 2p3p \ ^3D_1$	9456.9341	9455.000	0.1422D+14
$1s^2 2p3s \ ^3P_2$	9467.5091		0.1576D+14
$1s^2 2p3s \ ^1P_1$	9508.1716	9470.000	0.1135D+14
$1s^2 2p3p \ ^3D_2$	9520.2011	9524.000	0.2581D+14
$1s^2 2p3p \ ^3S_1(^1P_1)$	9521.0900		0.3794D+14
$1s^2 2p3p \ ^3P_0$	9539.4390		0.6826D+13
$1s^2 2p3d \ ^3F_2$	9582.1103		0.1346D+14
$1s^2 2p3p \ ^3P_1$	9612.0438		0.1998D+14
$1s^2 2p3p \ ^3D_3$	9618.9271	9624.000	0.3435D+14
$1s^2 2p3d \ ^3F_3$	9623.3964		0.5323D+14
$1s^2 2p3d \ ^3P_2(^3D_2)$	9633.4096	9753.000	0.1900D+15
$1s^2 2p3p \ ^3S_1(^1P_1)$	9633.7634		0.3794D+14
$1s^2 2p3p \ ^3P_2$	9636.2749		0.3616D+14
$1s^2 2p3d \ ^3D_1$	9650.0550	9637.000	0.8162D+14
$1s^2 2p3p \ ^1D_2$	9695.9224		0.4238D+14
$1s^2 2p3d \ ^3F_4$	9715.3896		0.2279D+11
$1s^2 2p3d \ ^1D_2$	9723.0885	9638.000	0.7766D+14
$1s^2 2p3d \ ^3D_3$	9746.0909	9749.000	0.1723D+15
$1s^2 2p3d \ ^3P_2(^3D_2)$	9763.8260	9753.000	0.1900D+15

Table I (continued)

state	energy (1000cm <sup>-1</sup> )		$\sum gAr(1/s)$
	present work	Shirai et.al.[17]	
1s <sup>2</sup> 2p3d <sup>3</sup> P <sub>1</sub>	9763.9267		0.6008D+14
1s <sup>2</sup> 2p3p <sup>1</sup> S <sub>0</sub>	9764.6273		0.6828D+13
1s <sup>2</sup> 2p3d <sup>3</sup> P <sub>0</sub>	9765.9251		0.1690D+14
1s <sup>2</sup> 2p3d <sup>1</sup> F <sub>3</sub>	9818.8896	9830.000	0.2580D+15
1s <sup>2</sup> 2p3d <sup>1</sup> P <sub>1</sub>	9827.9633	9828.000	0.6827D+14
1s <sup>2</sup> 2s4s <sup>3</sup> S <sub>1</sub>	11956.9647		0.5799D+13
1s <sup>2</sup> 2s4s <sup>1</sup> S <sub>0</sub>	11978.3784	11981.000	0.1460D+13
1s <sup>2</sup> 2s4p <sup>3</sup> P <sub>0</sub>	12022.0180		0.9760D+12
1s <sup>2</sup> 2s4p <sup>3</sup> P <sub>1</sub>	12023.6211	12024.000	0.6389D+13
1s <sup>2</sup> 2s4p <sup>3</sup> P <sub>2</sub>	12036.5661		0.4957D+13
1s <sup>2</sup> 2s4p <sup>1</sup> P <sub>1</sub>	12040.8051	12044.000	0.1786D+14
1s <sup>2</sup> 2s4d <sup>3</sup> D <sub>1</sub>	12073.3000	12073.000	0.2791D+14
1s <sup>2</sup> 2s4d <sup>3</sup> D <sub>2</sub>	12074.9435	12075.000	0.4592D+14
1s <sup>2</sup> 2s4d <sup>3</sup> D <sub>3</sub>	12077.9368	12081.000	0.6291D+14
1s <sup>2</sup> 2s4d <sup>1</sup> D <sub>2</sub>	12095.8155	12098.000	0.4060D+14
1s <sup>2</sup> 2s4f <sup>3</sup> F <sub>2</sub>	12099.3689		0.2018D+14
1s <sup>2</sup> 2s4f <sup>3</sup> F <sub>3</sub>	12100.1046		0.2817D+14
1s <sup>2</sup> 2s4f <sup>3</sup> F <sub>4</sub>	12101.5594		0.3603D+14
1s <sup>2</sup> 2s4f <sup>1</sup> F <sub>3</sub>	12105.7731		0.2749D+14
1s <sup>2</sup> 2p4s <sup>3</sup> P <sub>0</sub>	12365.6176		0.1448D+13
1s <sup>2</sup> 2p4s <sup>3</sup> P <sub>1</sub>	12369.9492		0.5162D+13
1s <sup>2</sup> 2p4p <sup>3</sup> D <sub>1</sub>	12414.5827		0.1099D+14
1s <sup>2</sup> 2p4p <sup>3</sup> P <sub>1</sub>	12440.2134		0.1180D+14
1s <sup>2</sup> 2p4p <sup>3</sup> D <sub>2</sub>	12441.8823	12443.000	0.1930D+14
1s <sup>2</sup> 2p4p <sup>3</sup> P <sub>0</sub>	12443.3447		0.3898D+13
1s <sup>2</sup> 2p4d <sup>3</sup> F <sub>2</sub>	12465.6063		0.1562D+14
1s <sup>2</sup> 2p4d <sup>3</sup> P <sub>2</sub> ( <sup>3</sup> D <sub>2</sub> )	12480.3829		0.8534D+14
1s <sup>2</sup> 2p4d <sup>3</sup> F <sub>3</sub>	12484.5020	12484.000	0.4913D+14
1s <sup>2</sup> 2p4d <sup>3</sup> D <sub>1</sub>	12489.0985		0.3243D+14
1s <sup>2</sup> 2p4s <sup>3</sup> P <sub>2</sub>	12493.8382		0.9762D+13
1s <sup>2</sup> 2p4f <sup>3</sup> G <sub>3</sub>	12494.0227		0.2807D+14
1s <sup>2</sup> 2p4f <sup>3</sup> F <sub>2</sub> ( <sup>3</sup> D <sub>2</sub> )	12497.2232		0.3988D+14
1s <sup>2</sup> 2p4f <sup>3</sup> D <sub>3</sub>	12498.1299		0.5580D+14
1s <sup>2</sup> 2p4f <sup>3</sup> G <sub>4</sub>	12498.2095		0.3503D+14
1s <sup>2</sup> 2p4s <sup>1</sup> P <sub>1</sub>	12499.7779		0.6132D+13
1s <sup>2</sup> 2p4p <sup>1</sup> P <sub>1</sub>	12548.4640		0.1201D+14
1s <sup>2</sup> 2p4p <sup>3</sup> P <sub>2</sub>	12554.4155		0.1971D+14
1s <sup>2</sup> 2p4p <sup>3</sup> D <sub>3</sub>	12555.5399	12560.000	0.2543D+14
1s <sup>2</sup> 2p4p <sup>3</sup> S <sub>1</sub>	12559.4162		0.1203D+14
1s <sup>2</sup> 2p4p <sup>1</sup> D <sub>2</sub>	12578.5215		0.2117D+14
1s <sup>2</sup> 2p4d <sup>1</sup> D <sub>2</sub>	12594.8100	12597.000	0.3124D+14

Table I (continued)

state	energy (1000cm <sup>-1</sup> )		$\sum gAr(1/s)$
	present work	Shirai et.al.[17]	
$1s^2 2p4d \ ^3F_4$	12594.9332		0.1832D+14
$1s^2 2p4p \ ^1S_0$	12601.6971		0.3730D+13
$1s^2 2p4d \ ^3D_3$	12601.9783	12603.000	0.5983D+14
$1s^2 2p4d \ ^3P_1$	12608.7309	12615.000	0.2675D+14
$1s^2 2p4d \ ^3P_2(^3D_2)$	12609.0068	12614.000	0.8534D+14
$1s^2 2p4d \ ^3P_0$	12609.5607		0.7670D+13
$1s^2 2p4f \ ^1F_3$	12616.2456		0.2814D+14
$1s^2 2p4f \ ^3F_4$	12618.7211		0.3575D+14
$1s^2 2p4f \ ^3F_2(^3D_2)$	12622.1187		0.3988D+14
$1s^2 2p4f \ ^3D_3$	12622.8548		0.5580D+14
$1s^2 2p4f \ ^3G_5$	12623.8101		0.4369D+14
$1s^2 2p4f \ ^1G_4$	12625.6836		0.3445D+14
$1s^2 2p4d \ ^1F_3$	12626.3489	12631.000	0.9687D+14
$1s^2 2p4f \ ^3D_1$	12628.2555		0.1206D+14
$1s^2 2p4d \ ^1P_1$	12630.7654		0.2905D+14
$1s^2 2p4f \ ^1D_2$	12632.1174		0.1977D+14
$1s^2 2s5s \ ^3S_1$	13340.1472		0.3262D+13
$1s^2 2s5s \ ^1S_0$	13350.6150		0.1051D+13
$1s^2 2s5p \ ^3P_0$	13372.3330		0.7769D+12
$1s^2 2s5p \ ^3P_1$	13373.2644		0.3910D+13
$1s^2 2s5p \ ^3P_2$	13379.7439		0.3914D+13
$1s^2 2s5p \ ^1P_1$	13382.6334	13380.000	0.1040D+14
$1s^2 2s5d \ ^3D_1$	13398.2552	13369.000	0.1405D+14
$1s^2 2s5d \ ^3D_2$	13399.0603	13400.000	0.2321D+14
$1s^2 2s5d \ ^3D_3$	13400.5873	13404.000	0.3192D+14
$1s^2 2s5d \ ^1D_2$	13408.6652		0.2198D+14
$1s^2 2s5f \ ^3F_2$	13411.0161		0.1049D+14
$1s^2 2s5f \ ^3F_3$	13411.4109		0.1465D+14
$1s^2 2s5f \ ^3F_4$	13412.1359		0.1876D+14
$1s^2 2s5g \ ^3G_3$	13414.5857		0.8510D+13
$1s^2 2s5g \ ^3G_4$	13414.6910		0.1094D+14
$1s^2 2s5f \ ^1F_3$	13414.8637		0.1442D+14
$1s^2 2s5g \ ^3G_5$	13415.2810		0.1334D+14
$1s^2 2s5g \ ^1G_4$	13415.3497		0.1088D+14
$1s^2 2p5s \ ^3P_0$	13740.5573		0.9003D+12
$1s^2 2p5s \ ^3P_1$	13742.3442		0.3206D+13
$1s^2 2p5p \ ^3D_1$	13765.8425		0.6810D+13
$1s^2 2p5p \ ^3P_1$	13778.3225		0.7066D+13
$1s^2 2p5p \ ^3P_0$	13779.4794		0.2189D+13
$1s^2 2p5p \ ^3D_2$	13779.5178		0.2300D+14
$1s^2 2p5d \ ^3F_2$	13791.6814		0.9643D+13
$1s^2 2p5d \ ^3P_2(^3D_2)$	13799.8628		0.4601D+14

Table I (continued)

state	energy (1000cm <sup>-1</sup> )		$\sum gAr(1/s)$
	present work	Shirai et.al.[17]	
$1s^2 2p5d \ ^3F_3$	13801.3344	13804.000	0.2930D+14
$1s^2 2p5d \ ^3D_1$	13802.7551		0.1650D+14
$1s^2 2p5f \ ^3G_3$	13805.5873		0.1458D+14
$1s^2 2p5f \ ^3F_2$	13807.6145		0.2076D+14
$1s^2 2p5f \ ^3D_3(^3F_3)$	13807.8383		0.2904D+14
$1s^2 2p5f \ ^3G_4$	13808.0317		0.1819D+14
$1s^2 2p5g \ ^3H_4$	13808.7443		0.1082D+14
$1s^2 2p5g \ ^3G_3(^1F_3)$	13809.1446		0.1725D+14
$1s^2 2p5g \ ^3H_5$	13809.4533		0.1316D+14
$1s^2 2p5g \ ^3F_4(^1G_4)$	13809.7474		0.2182D+14
$1s^2 2p5s \ ^3P_2$	13866.1620		0.4483D+13
$1s^2 2p5s \ ^1P_1$	13869.5662		0.3772D+13
$1s^2 2p5p \ ^1P_1$	13895.2625		0.7190D+13
$1s^2 2p5p \ ^3D_2$	13897.6771		0.2300D+14
$1s^2 2p5p \ ^3D_3$	13899.3086	13904.000	0.1554D+14
$1s^2 2p5p \ ^3S_1$	13900.5980		0.7470D+13
$1s^2 2p5p \ ^1D_2$	13909.5561		0.1195D+14
$1s^2 2p5d \ ^1D_2$	13918.8817	13922.000	0.1645D+14
$1s^2 2p5d \ ^3F_4$	13919.2935		0.1278D+14
$1s^2 2p5p \ ^1S_0$	13919.6196		0.2009D+13
$1s^2 2p5d \ ^3D_3$	13922.0994	13929.000	0.2949D+14
$1s^2 2p5d \ ^3P_1$	13925.1138		0.1401D+14
$1s^2 2p5d \ ^3P_2(^3D_2)$	13925.3917		0.4601D+14
$1s^2 2p5d \ ^3P_0$	13925.4502		0.4071D+13
$1s^2 2p5f \ ^1F_3$	13929.7345		0.1456D+14
$1s^2 2p5f \ ^3F_4$	13931.0177		0.1854D+14
$1s^2 2p5f \ ^3F_2$	13932.6088		0.2076D+14
$1s^2 2p5f \ ^3D_3(^3F_3)$	13932.8966		0.2904D+14
$1s^2 2p5f \ ^3G_5$	13933.1734		0.2276D+14
$1s^2 2p5g \ ^3G_4$	13933.2690		0.1092D+14
$1s^2 2p5d \ ^1F_3$	13933.3831	13945.000	0.4784D+14
$1s^2 2p5g \ ^3G_5$	13933.8926		0.1330D+14
$1s^2 2p5g \ ^3G_3(^1F_3)$	13934.3049		0.1725D+14
$1s^2 2p5f \ ^1G_4$	13934.6683		0.1787D+14
$1s^2 2p5g \ ^3F_4(^1G_4)$	13934.8866		0.2182D+14
$1s^2 2p5g \ ^1H_5$	13935.1630		0.1310D+14
$1s^2 2p5f \ ^3D_1$	13935.3870		0.6307D+13
$1s^2 2p5d \ ^1P_1$	13935.5699		0.1524D+14
$1s^2 2p5g \ ^3H_6$	13935.6891		0.1548D+14
$1s^2 2p5g \ ^3F_2$	13936.4973		0.6070D+13
$1s^2 2p5g \ ^3F_3$	13937.1839		0.8520D+13
$1s^2 2p5f \ ^1D_2$	13937.6619		0.1038D+14

**Table II** Energy level(in units of  $1000\text{cm}^{-1}$ ),  $\sum gAr$  (in units of  $1/\text{s}$ ), and  $Aa$  (in units of  $10^{13}/\text{s}$ ) for  $1s^2 2pn l_2$  ( $n=11, 13, 20$ ) and  $1s^2 3l_1 n l_2$  ( $n=3, 4$ ). For  $1s^2 3l_1 n l_2$ ,  $\sum Aa$  is given in fourth column.

state	energy( $1000\text{cm}^{-1}$ )	$\sum gAr(1/\text{s})$	$Aa(10^{13}/\text{s})$
$1s^2 2p11s$ $^3P_0$	15647.2331	0.1245D+12	0.00000D+00
$1s^2 2p11s$ $^3P_1$	15647.3615	0.4036D+12	0.00000D+00
$1s^2 2p11p$ $^3D_1$	15649.6018	0.1451D+13	0.00000D+00
$1s^2 2p11p$ $^3P_0$	15650.6947	0.3323D+12	0.00000D+00
$1s^2 2p11p$ $^3P_1(^1P_1)$	15650.6990	0.2925D+13	0.00000D+00
$1s^2 2p11p$ $^3D_2(^1D_2)$	15650.7981	0.3965D+13	0.00000D+00
$1s^2 2p11d$ $^3F_2$	15651.9678	0.1151D+13	0.00000D+00
$1s^2 2p11d$ $^3P_2(^3D_2)$	15652.6700	0.4387D+13	0.00000D+00
$1s^2 2p11d$ $^3F_3$	15652.7976	0.2812D+13	0.00000D+00
$1s^2 2p11d$ $^3D_1$	15652.8674	0.1477D+13	0.00000D+00
$1s^2 2p11f$ $^3G_3$	15653.1884	0.1464D+13	0.00000D+00
$1s^2 2p11f$ $^3F_2(^3D_2)$	15653.3755	0.2151D+13	0.00000D+00
$1s^2 2p11f$ $^3D_3(^3F_3)$	15653.3901	0.2873D+13	0.00000D+00
$1s^2 2p11f$ $^3G_4$	15653.4433	0.1770D+13	0.00000D+00
$1s^2 2p11g$ $^3H_4$	15653.5150	0.1022D+13	0.00000D+00
$1s^2 2p11g$ $^3G_3(^1F_3)$	15653.5324	0.1649D+13	0.00000D+00
$1s^2 2p11g$ $^3F_4(^1G_4)$	15653.5838	0.2090D+13	0.00000D+00
$1s^2 2p11g$ $^3H_5$	15653.5912	0.1243D+13	0.00000D+00
$1s^2 2p11s$ $^3P_2$	15772.9685	0.6534D+12	0.18119D+13
$1s^2 2p11s$ $^1P_1$	15773.2301	0.4501D+12	0.15510D+14
$1s^2 2p11p$ $^3P_1(^1P_1)$	15775.6569	0.2925D+13	0.85709D+13
$1s^2 2p11p$ $^3D_2(^1D_2)$	15775.8272	0.3965D+13	0.23866D+13
$1s^2 2p11p$ $^3D_3$	15776.0679	0.2908D+13	0.99405D+12
$1s^2 2p11p$ $^3S_1$	15776.1702	0.1740D+13	0.14867D+14
$1s^2 2p11p$ $^3P_2$	15776.8700	0.1894D+13	0.12840D+13
$1s^2 2p11p$ $^1S_0$	15777.6913	0.3080D+12	0.28387D+14
$1s^2 2p11d$ $^1D_2$	15777.8405	0.1701D+13	0.21348D+12
$1s^2 2p11d$ $^3F_4$	15777.8970	0.1756D+13	0.24760D+12
$1s^2 2p11d$ $^3D_3$	15778.0985	0.2790D+13	0.37899D+12
$1s^2 2p11d$ $^3P_1$	15778.3508	0.1350D+13	0.55553D+13
$1s^2 2p11d$ $^3P_0$	15778.3814	0.4005D+12	0.79394D+13
$1s^2 2p11d$ $^3P_2(^3D_2)$	15778.3873	0.4387D+13	0.76363D+13
$1s^2 2p11f$ $^1F_3$	15778.8155	0.1466D+13	0.12016D+13
$1s^2 2p11f$ $^3F_4$	15778.9369	0.1866D+13	0.17260D+13
$1s^2 2p11d$ $^1F_3$	15779.0224	0.4084D+13	0.31448D+13
$1s^2 2p11f$ $^3F_2(^3D_2)$	15779.0512	0.2151D+13	0.39313D+12
$1s^2 2p11f$ $^3D_3(^3F_3)$	15779.0784	0.2873D+13	0.57809D+12
$1s^2 2p11f$ $^3G_5$	15779.1066	0.2365D+13	0.47561D+13
$1s^2 2p11g$ $^3G_4$	15779.1516	0.1070D+13	0.17090D+13

Table II (continued)

state	energy(1000cm <sup>-1</sup> )	$\sum gAr(1/s)$	$Aa(1/s)$
1s <sup>2</sup> 2p11d <sup>1</sup> P <sub>1</sub>	15779.1948	0.1341D+13	0.59719D+13
1s <sup>2</sup> 2p11g <sup>3</sup> G <sub>5</sub>	15779.2135	0.1304D+13	0.18329D+13
1s <sup>2</sup> 2p11g <sup>3</sup> G <sub>3</sub> ( <sup>1</sup> F <sub>3</sub> )	15779.2308	0.1649D+13	0.26390D+12
1s <sup>2</sup> 2p11f <sup>3</sup> D <sub>1</sub>	15779.2630	0.6335D+12	0.51873D+12
1s <sup>2</sup> 2p11g <sup>3</sup> F <sub>4</sub> ( <sup>1</sup> G <sub>4</sub> )	15779.2823	0.2090D+13	0.21623D+12
1s <sup>2</sup> 2p11f <sup>1</sup> G <sub>4</sub>	15779.2961	0.1732D+13	0.55143D+13
1s <sup>2</sup> 2p11g <sup>1</sup> H <sub>5</sub>	15779.3601	0.1292D+13	0.64105D+13
1s <sup>2</sup> 2p11g <sup>3</sup> H <sub>6</sub>	15779.3943	0.1534D+13	0.64023D+13
1s <sup>2</sup> 2p11g <sup>3</sup> F <sub>2</sub>	15779.4119	0.5926D+12	0.14780D+12
1s <sup>2</sup> 2p11g <sup>3</sup> F <sub>3</sub>	15779.4832	0.8279D+12	0.33883D+12
1s <sup>2</sup> 2p11f <sup>1</sup> D <sub>2</sub>	15779.4843	0.1196D+13	0.35418D+12
1s <sup>2</sup> 2p13s <sup>3</sup> P <sub>0</sub>	15786.1087	0.7407D+11	0.13917D+13
1s <sup>2</sup> 2p13s <sup>3</sup> P <sub>1</sub>	15786.1850	0.2416D+12	0.55504D+13
1s <sup>2</sup> 2p13p <sup>3</sup> D <sub>1</sub>	15787.5415	0.7032D+12	0.11522D+13
1s <sup>2</sup> 2p13p <sup>3</sup> P <sub>1</sub> ( <sup>1</sup> P <sub>1</sub> )	15788.2035	0.1430D+13	0.56517D+13
1s <sup>2</sup> 2p13p <sup>3</sup> P <sub>0</sub>	15788.2080	0.2808D+12	0.83592D+13
1s <sup>2</sup> 2p13p <sup>3</sup> D <sub>2</sub> ( <sup>1</sup> D <sub>2</sub> )	15788.2702	0.2494D+13	0.18626D+13
1s <sup>2</sup> 2p13d <sup>3</sup> F <sub>2</sub>	15788.9715	0.6975D+12	0.13683D+12
1s <sup>2</sup> 2p13d <sup>3</sup> P <sub>2</sub> ( <sup>3</sup> D <sub>2</sub> )	15789.3959	0.2637D+13	0.49550D+13
1s <sup>2</sup> 2p13d <sup>3</sup> F <sub>3</sub>	15789.4849	0.1745D+13	0.78846D+12
1s <sup>2</sup> 2p13d <sup>3</sup> D <sub>1</sub>	15789.5175	0.8844D+12	0.24440D+13
1s <sup>2</sup> 2p13f <sup>3</sup> G <sub>3</sub>	15789.7068	0.8432D+12	0.17220D+13
1s <sup>2</sup> 2p13f <sup>3</sup> F <sub>2</sub> ( <sup>3</sup> D <sub>2</sub> )	15789.8244	0.1181D+13	0.25255D+12
1s <sup>2</sup> 2p13f <sup>3</sup> D <sub>3</sub> ( <sup>3</sup> F <sub>3</sub> )	15789.8308	0.1628D+13	0.35823D+12
1s <sup>2</sup> 2p13f <sup>3</sup> G <sub>4</sub>	15789.8630	0.1014D+13	0.20997D+13
1s <sup>2</sup> 2p13g <sup>3</sup> H <sub>4</sub>	15789.9056	0.5940D+12	0.88698D+12
1s <sup>2</sup> 2p13g <sup>3</sup> G <sub>3</sub> ( <sup>1</sup> F <sub>3</sub> )	15789.9171	0.9514D+12	0.17643D+12
1s <sup>2</sup> 2p13g <sup>3</sup> F <sub>4</sub> ( <sup>1</sup> G <sub>4</sub> )	15789.9481	0.1208D+13	0.13299D+12
1s <sup>2</sup> 2p13g <sup>3</sup> H <sub>5</sub>	15789.9523	0.7202D+12	0.24732D+13
1s <sup>2</sup> 2p13s <sup>3</sup> P <sub>2</sub>	15911.8487	0.3915D+12	0.13071D+13
1s <sup>2</sup> 2p13s <sup>1</sup> P <sub>1</sub>	15912.0035	0.2683D+12	0.55504D+13
1s <sup>2</sup> 2p13p <sup>3</sup> P <sub>1</sub> ( <sup>1</sup> P <sub>1</sub> )	15913.4738	0.1430D+13	0.56517D+13
1s <sup>2</sup> 2p13p <sup>3</sup> D <sub>2</sub> ( <sup>1</sup> D <sub>2</sub> )	15913.5741	0.2494D+13	0.18626D+13
1s <sup>2</sup> 2p13p <sup>3</sup> D <sub>3</sub>	15913.7240	0.1748D+13	0.81151D+12
1s <sup>2</sup> 2p13p <sup>3</sup> S <sub>1</sub>	15913.7830	0.6568D+12	0.95589D+13
1s <sup>2</sup> 2p13p <sup>3</sup> P <sub>2</sub>	15914.2033	0.1262D+13	0.94014D+12
1s <sup>2</sup> 2p13p <sup>1</sup> S <sub>0</sub>	15914.6893	0.2943D+12	0.17662D+14
1s <sup>2</sup> 2p13d <sup>1</sup> D <sub>2</sub>	15914.7946	0.1022D+13	0.11043D+12
1s <sup>2</sup> 2p13d <sup>3</sup> F <sub>4</sub>	15914.8297	0.1059D+13	0.56002D+11

Table II (continued)

state	energy(1000cm <sup>-1</sup> )	$\sum gAr(1/s)$	$Aa(1/s)$
1s <sup>2</sup> 2p13d <sup>3</sup> D <sub>3</sub>	15914.9484	0.1666D+13	0.15768D+12
1s <sup>2</sup> 2p13d <sup>3</sup> P <sub>1</sub>	15915.0997	0.8111D+12	0.35416D+13
1s <sup>2</sup> 2p13d <sup>3</sup> P <sub>0</sub>	15915.1177	0.2404D+12	0.50670D+13
1s <sup>2</sup> 2p13d <sup>3</sup> P <sub>2</sub> ( <sup>3</sup> D <sub>2</sub> )	15915.1226	0.2637D+13	0.49550D+13
1s <sup>2</sup> 2p13f <sup>1</sup> F <sub>3</sub>	15915.3833	0.8240D+12	0.59448D+12
1s <sup>2</sup> 2p13f <sup>3</sup> F <sub>4</sub>	15915.4566	0.1048D+13	0.85875D+12
1s <sup>2</sup> 2p13d <sup>1</sup> F <sub>3</sub>	15915.5015	0.2445D+13	0.15395D+13
1s <sup>2</sup> 2p13f <sup>3</sup> F <sub>2</sub> ( <sup>3</sup> D <sub>2</sub> )	15915.5252	0.1181D+13	0.25255D+12
1s <sup>2</sup> 2p13f <sup>3</sup> D <sub>3</sub> ( <sup>3</sup> F <sub>3</sub> )	15915.5414	0.1628D+13	0.35823D+12
1s <sup>2</sup> 2p13f <sup>3</sup> G <sub>5</sub>	15915.5576	0.1333D+13	0.23641D+13
1s <sup>2</sup> 2p13g <sup>3</sup> G <sub>4</sub>	15915.5871	0.6189D+12	0.88698D+12
1s <sup>2</sup> 2p13d <sup>1</sup> P <sub>1</sub>	15915.6060	0.8068D+12	0.37856D+13
1s <sup>2</sup> 2p13g <sup>3</sup> G <sub>5</sub>	15915.6248	0.7540D+12	0.95406D+12
1s <sup>2</sup> 2p13g <sup>3</sup> G <sub>3</sub> ( <sup>1</sup> F <sub>3</sub> )	15915.6352	0.9514D+12	0.17643D+12
1s <sup>2</sup> 2p13f <sup>3</sup> D <sub>1</sub>	15915.6523	0.3521D+12	0.33654D+12
1s <sup>2</sup> 2p13g <sup>3</sup> F <sub>4</sub> ( <sup>1</sup> G <sub>4</sub> )	15915.6662	0.1208D+13	0.13299D+12
1s <sup>2</sup> 2p13f <sup>1</sup> G <sub>4</sub>	15915.6738	0.9642D+12	0.28097D+13
1s <sup>2</sup> 2p13g <sup>1</sup> H <sub>5</sub>	15915.7136	0.7425D+12	0.33463D+13
1s <sup>2</sup> 2p13g <sup>3</sup> H <sub>6</sub>	15915.7336	0.8862D+12	0.33385D+13
1s <sup>2</sup> 2p13g <sup>3</sup> F <sub>2</sub>	15915.7444	0.3397D+12	0.98431D+11
1s <sup>2</sup> 2p13f <sup>1</sup> D <sub>2</sub>	15915.7865	0.6072D+12	0.22343D+12
1s <sup>2</sup> 2p13g <sup>3</sup> F <sub>3</sub>	15915.7879	0.4738D+12	0.21772D+12
1s <sup>2</sup> 2p20s <sup>3</sup> P <sub>0</sub>	15987.3909	0.2052D+11	0.40611D+12
1s <sup>2</sup> 2p20s <sup>3</sup> P <sub>1</sub>	15987.4112	0.6882D+11	0.15220D+13
1s <sup>2</sup> 2p20p <sup>3</sup> D <sub>1</sub>	15987.7834	0.1966D+12	0.32755D+12
1s <sup>2</sup> 2p20p <sup>3</sup> P <sub>1</sub> ( <sup>1</sup> P <sub>1</sub> )	15987.9636	0.3951D+12	0.15774D+13
1s <sup>2</sup> 2p20p <sup>3</sup> P <sub>0</sub>	15987.9642	0.7051D+11	0.23574D+13
1s <sup>2</sup> 2p20p <sup>3</sup> D <sub>2</sub> ( <sup>1</sup> D <sub>2</sub> )	15987.9826	0.6722D+12	0.54246D+12
1s <sup>2</sup> 2p20d <sup>3</sup> F <sub>2</sub> ( <sup>3</sup> D <sub>2</sub> )	15988.1741	0.1935D+12	0.32269D+11
1s <sup>2</sup> 2p20d <sup>3</sup> P <sub>2</sub> ( <sup>3</sup> D <sub>2</sub> )	15988.2894	0.7330D+12	0.13416D+13
1s <sup>2</sup> 2p20d <sup>3</sup> F <sub>3</sub>	15988.3149	0.9519D+12	0.23012D+12
1s <sup>2</sup> 2p20d <sup>3</sup> D <sub>1</sub>	15988.3224	0.2450D+12	0.66108D+12
1s <sup>2</sup> 2p20f <sup>3</sup> G <sub>3</sub>	15988.3728	0.2200D+12	0.40939D+12
1s <sup>2</sup> 2p20f <sup>3</sup> F <sub>2</sub> ( <sup>3</sup> D <sub>2</sub> )	15988.4057	0.2980D+12	0.67991D+11
1s <sup>2</sup> 2p20f <sup>3</sup> D <sub>3</sub> ( <sup>3</sup> F <sub>3</sub> )	15988.4072	0.4225D+12	0.92246D+11
1s <sup>2</sup> 2p20f <sup>3</sup> G <sub>4</sub>	15988.4161	0.2661D+12	0.50942D+12
1s <sup>2</sup> 2p20g <sup>3</sup> H <sub>4</sub>	15988.4257	0.1566D+12	0.58517D+12
1s <sup>2</sup> 2p20g <sup>3</sup> G <sub>3</sub> ( <sup>1</sup> F <sub>3</sub> )	15988.4290	0.2646D+12	0.53209D+11
1s <sup>2</sup> 2p20g <sup>3</sup> F <sub>4</sub> ( <sup>1</sup> G <sub>4</sub> )	15988.4371	0.3350D+12	0.38366D+11
1s <sup>2</sup> 2p20g <sup>3</sup> H <sub>5</sub>	15988.4383	0.1903D+12	0.60201D+12



Table II (continued)

state	energy(1000cm <sup>-1</sup> )	$\sum gAr(1/s)$	$Aa(1/s)$
$1s^2 2p 20s$ $^3P_2$	16113.1385	0.1177D+12	0.37611D+12
$1s^2 2p 20s$ $^1P_1$	16113.1797	0.8354D+11	0.26808D+13
$1s^2 2p 20p$ $^3P_1(^1P_1)$	16113.5831	0.3951D+12	0.15774D+13
$1s^2 2p 20p$ $^3D_2$	16113.6098	0.6722D+12	0.54246D+12
$1s^2 2p 20p$ $^3D_3(^1D_2)$	16113.6522	0.4758D+12	0.23715D+12
$1s^2 2p 20p$ $^3S_1$	16113.6673	0.1900D+12	0.25927D+13
$1s^2 2p 20p$ $^3P_2$	16113.7814	0.3355D+12	0.26844D+12
$1s^2 2p 20p$ $^1S_0$	16113.9113	0.7079D+11	0.48141D+13
$1s^2 2p 20d$ $^1D_2$	16113.9442	0.2903D+12	0.27321D+11
$1s^2 2p 20d$ $^3F_4$	16113.9541	0.3106D+12	0.78877D+10
$1s^2 2p 20d$ $^3F_3(^3D_2)$	16113.9857	0.9519D+12	0.23012D+12
$1s^2 2p 20d$ $^3P_1$	16114.0266	0.2283D+12	0.23064D+11
$1s^2 2p 20d$ $^3P_0$	16114.0313	0.6796D+11	0.13640D+13
$1s^2 2p 20d$ $^3P_2(^3D_2)$	16114.0331	0.7330D+12	0.13416D+13
$1s^2 2p 20f$ $^1F_3$	16114.1035	0.2157D+12	0.14038D+12
$1s^2 2p 20f$ $^3F_4$	16114.1236	0.2745D+12	0.20494D+12
$1s^2 2p 20d$ $^1F_3$	16114.1354	0.6794D+12	0.37037D+12
$1s^2 2p 20f$ $^3F_2(^3D_2)$	16114.1422	0.2980D+12	0.67991D+11
$1s^2 2p 20f$ $^3D_3(^3F_3)$	16114.1467	0.4225D+12	0.92246D+11
$1s^2 2p 20f$ $^3G_5$	16114.1507	0.3474D+12	0.56269D+12
$1s^2 2p 20g$ $^3G_4$	16114.1574	0.1793D+12	0.21442D+12
$1s^2 2p 20d$ $^1P_1$	16114.1639	0.2331D+12	0.78202D+12
$1s^2 2p 20g$ $^3G_5$	16114.1675	0.2187D+12	0.23077D+12
$1s^2 2p 20g$ $^3G_3(^1F_3)$	16114.1707	0.2646D+12	0.53209D+11
$1s^2 2p 20f$ $^3D_1$	16114.1766	0.9028D+11	0.88970D+11
$1s^2 2p 20g$ $^3F_4(^1G_4)$	16114.1788	0.3350D+12	0.38366D+11
$1s^2 2p 20f$ $^1G_4$	16114.1831	0.2541D+12	0.68804D+12
$1s^2 2p 20g$ $^1H_5$	16114.1922	0.2163D+12	0.81329D+12
$1s^2 2p 20g$ $^3H_6$	16114.1971	0.2571D+12	0.80966D+12
$1s^2 2p 20g$ $^3F_2$	16114.2005	0.9887D+11	0.30223D+11
$1s^2 2p 20g$ $^3F_3$	16114.2122	0.1381D+12	0.64931D+11
$1s^2 2p 20f$ $^1D_2$	16114.2137	0.1460D+12	0.61361D+11

Table II (continued)

state	energy(1000cm <sup>-1</sup> )	$\sum gAr(1/s)$	Aa(1/s) to 2s	$\sum Aa(1/s)$
1s <sup>2</sup> 3s3p <sup>3</sup> P <sub>0</sub>	18175.2650	0.1009D+14	0.58203D+14	0.12316D+15
1s <sup>2</sup> 3s3p <sup>3</sup> P <sub>1</sub>	18184.5179	0.3052D+14	0.60554D+14	0.13123D+15
1s <sup>2</sup> 3s3p <sup>3</sup> P <sub>2</sub>	18210.4066	0.5138D+14	0.57618D+14	0.12532D+15
1s <sup>2</sup> 3s3d <sup>1</sup> D <sub>2</sub>	18263.5123	0.1041D+15	0.10646D+15	0.23008D+15
1s <sup>2</sup> 3s3d <sup>3</sup> D <sub>1</sub>	18282.4402	0.7365D+14	0.21147D+14	0.36788D+14
1s <sup>2</sup> 3s3d <sup>3</sup> D <sub>2</sub>	18287.8013	0.1201D+15	0.22802D+14	0.44108D+14
1s <sup>2</sup> 3s3d <sup>3</sup> D <sub>3</sub>	18294.0146	0.1671D+15	0.21145D+14	0.36780D+14
1s <sup>2</sup> 3s3p <sup>1</sup> P <sub>1</sub>	18296.2095	0.3703D+14	0.17226D+15	0.48082D+15
1s <sup>2</sup> 3p <sup>2</sup> <sup>3</sup> P <sub>0</sub>	18299.1022	0.1575D+14	0.57373D+13	0.27124D+15
1s <sup>2</sup> 3p <sup>2</sup> <sup>3</sup> P <sub>1</sub>	18318.6961	0.4742D+14	0.12302D+03	0.23374D+15
1s <sup>2</sup> 3p <sup>2</sup> <sup>3</sup> P <sub>2</sub>	18341.1206	0.8218D+14	0.12826D+14	0.23998D+15
1s <sup>2</sup> 3p <sup>2</sup> <sup>1</sup> S <sub>0</sub>	18417.2299	0.1962D+14	0.62735D+14	0.80821D+15
1s <sup>2</sup> 3p <sup>2</sup> <sup>1</sup> D <sub>2</sub>	18418.0893	0.1105D+15	0.16799D+15	0.42864D+15
1s <sup>2</sup> 3p3d <sup>3</sup> F <sub>2</sub>	18423.0127	0.1458D+15	0.11038D+12	0.45377D+13
1s <sup>2</sup> 3p3d <sup>3</sup> F <sub>3</sub>	18444.8773	0.2018D+15	0.66003D+12	0.80913D+13
1s <sup>2</sup> 3p3d <sup>1</sup> D <sub>2</sub>	18455.9009	0.1427D+15	0.32496D+12	0.21054D+14
1s <sup>2</sup> 3p3d <sup>3</sup> F <sub>4</sub>	18469.1360	0.2577D+15	0.11491D+12	0.22613D+12
1s <sup>2</sup> 3p3d <sup>3</sup> D <sub>1</sub>	18481.6784	0.8641D+14	0.51749D+12	0.15671D+15
1s <sup>2</sup> 3p3d <sup>3</sup> D <sub>2</sub>	18490.6019	0.1433D+15	0.64428D+12	0.14044D+15
1s <sup>2</sup> 3p3d <sup>3</sup> D <sub>3</sub>	18506.0530	0.2000D+15	0.16989D+13	0.16710D+15
1s <sup>2</sup> 3p3d <sup>3</sup> P <sub>1</sub>	18514.1083	0.8658D+14	0.28298D+13	0.98154D+14
1s <sup>2</sup> 3p3d <sup>3</sup> P <sub>0</sub>	18514.4327	0.2896D+14	0.30666D+13	0.90218D+14
1s <sup>2</sup> 3p3d <sup>3</sup> P <sub>2</sub>	18516.2171	0.1422D+15	0.26863D+13	0.10101D+15
1s <sup>2</sup> 3d <sup>2</sup> <sup>3</sup> F <sub>2</sub>	18521.0363	0.2159D+15	0.18437D+12	0.19016D+15
1s <sup>2</sup> 3d <sup>2</sup> <sup>3</sup> F <sub>3</sub>	18528.3758	0.2997D+15	0.58155D+05	0.18977D+15
1s <sup>2</sup> 3d <sup>2</sup> <sup>3</sup> F <sub>4</sub>	18537.0743	0.3810D+15	0.19288D+13	0.20086D+15
1s <sup>2</sup> 3d <sup>2</sup> <sup>3</sup> P <sub>0</sub>	18576.9833	0.4225D+14	0.18231D+12	0.22110D+14
1s <sup>2</sup> 3d <sup>2</sup> <sup>3</sup> P <sub>1</sub>	18580.6515	0.1258D+15	0.49098D+03	0.20102D+14
1s <sup>2</sup> 3d <sup>2</sup> <sup>3</sup> P <sub>2</sub>	18582.6793	0.2069D+15	0.11760D+12	0.69696D+14
1s <sup>2</sup> 3d <sup>2</sup> <sup>1</sup> G <sub>4</sub>	18583.1269	0.3762D+15	0.19099D+15	0.12878D+16
1s <sup>2</sup> 3p3d <sup>1</sup> F <sub>3</sub>	18583.4460	0.2007D+15	0.94980D+14	0.63168D+15
1s <sup>2</sup> 3d <sup>2</sup> <sup>1</sup> D <sub>2</sub>	18600.7984	0.1952D+15	0.29137D+12	0.23575D+15
1s <sup>2</sup> 3p3d <sup>1</sup> P <sub>1</sub>	18629.0642	0.7985D+14	0.16217D+12	0.28288D+15
1s <sup>2</sup> 3d <sup>2</sup> <sup>1</sup> S <sub>0</sub>	18736.9354	0.3759D+14	0.31480D+14	0.93879D+14
1s <sup>2</sup> 3s4s <sup>3</sup> S <sub>1</sub>	21130.9173	0.1180D+14	0.16984D+13	0.17315D+13
1s <sup>2</sup> 3s4s <sup>1</sup> S <sub>0</sub>	21157.2833	0.4138D+13	0.86975D+14	0.87152D+14
1s <sup>2</sup> 3s4d <sup>3</sup> D <sub>1</sub>	21235.5553	0.3045D+14	0.92551D+13	0.11943D+14
1s <sup>2</sup> 3s4p <sup>1</sup> P <sub>1</sub>	21238.8416	0.1915D+14	0.21838D+14	0.53449D+14
1s <sup>2</sup> 3s4d <sup>3</sup> D <sub>2</sub>	21239.4271	0.5021D+14	0.11531D+14	0.14441D+14

Table II (continued)

state	energy(1000cm <sup>-1</sup> )	$\sum gAr(1/s)$	$Aa(1/s)$ to $2s$	$\sum Aa(1/s)$
$1s^23s4p$ $^3P_0$	21243.4627	0.6135D+13	0.32167D+14	0.53449D+14
$1s^23s4d$ $^3D_3$	21244.7279	0.6876D+14	0.10855D+14	0.13838D+14
$1s^23s4d$ $^1D_2$	21248.9432	0.4897D+14	0.51520D+14	0.67014D+14
$1s^23s4p$ $^3P_1$	21250.6607	0.1893D+14	0.28441D+14	0.47771D+14
$1s^23s4p$ $^3P_2$	21257.8516	0.3103D+14	0.30355D+14	0.50325D+14
$1s^23p4p$ $^1P_1$	21304.6989	0.7269D+14	0.49739D+13	0.32525D+14
$1s^23p4s$ $^3P_0$	21308.1417	0.9360D+13	0.50734D+13	0.25582D+14
$1s^23p4s$ $^3P_1$	21315.8405	0.2811D+14	0.15348D+14	0.51053D+14
$1s^23p4p$ $^3D_2$	21322.7477	0.7118D+14	0.47540D+13	0.38854D+14
$1s^23p4p$ $^1P_1$	21325.0249	0.7269D+14	0.49739D+13	0.32525D+14
$1s^23s4f$ $^3F_2$	21328.5149	0.1575D+14	0.55761D+13	0.61170D+13
$1s^23s4f$ $^3F_3$	21329.9125	0.2127D+14	0.57989D+13	0.63586D+13
$1s^23s4f$ $^3F_4$	21331.7360	0.2639D+14	0.60249D+13	0.65968D+13
$1s^23p4p$ $^3P_0$	21333.3032	0.1135D+14	0.25788D+13	0.16209D+15
$1s^23d4s$ $^1D_2$	21339.7850	0.8224D+14	0.11317D+14	0.77278D+14
$1s^23p4s$ $^3P_2$	21341.3325	0.4947D+14	0.65432D+13	0.29552D+14
$1s^23s4f$ $^1F_3$	21343.2365	0.3389D+14	0.43156D+13	0.42992D+14
$1s^23p4p$ $^3D_3$	21345.3639	0.1071D+15	0.43625D+13	0.11197D+14
$1s^23p4p$ $^3P_1$	21348.0979	0.3819D+14	0.51969D+12	0.89659D+14
$1s^23p4p$ $^3S_1$	21358.7781	0.3916D+14	0.18888D+12	0.37079D+14
$1s^23p4p$ $^3P_2$	21363.5775	0.7403D+14	0.73321D+13	0.10859D+15
$1s^23p4s$ $^1P_1$	21364.3007	0.2950D+14	0.84617D+14	0.25181D+15
$1s^23d4s$ $^3D_1$	21366.2473	0.6226D+14	0.92526D+12	0.48200D+13
$1s^23d4s$ $^3D_2$	21372.6235	0.9240D+14	0.64419D+12	0.34280D+14
$1s^23d4s$ $^3D_3$	21381.2451	0.1267D+15	0.43760D+12	0.15468D+13
$1s^23p4f$ $^3G_3$	21386.9228	0.5955D+14	0.16611D+13	0.85762D+13
$1s^23p4f$ $^3G_4$	21398.2329	0.7889D+14	0.64416D+13	0.17038D+14
$1s^23p4f$ $^3F_3$	21400.8592	0.1316D+15	0.63164D+12	0.70849D+13
$1s^23p4d$ $^3D_1$	21401.4394	0.5166D+14	0.25062D+12	0.14373D+14
$1s^23p4f$ $^3F_2$	21403.3251	0.4552D+14	0.10315D+10	0.14712D+13
$1s^23p4d$ $^3D_2$	21408.1231	0.8170D+14	0.73392D+12	0.14182D+14
$1s^23p4d$ $^3D_3$	21410.6932	0.2359D+15	0.13643D+13	0.15349D+14
$1s^23p4d$ $^1D_2$	21412.9685	0.1610D+15	0.22493D+13	0.14395D+14
$1s^23p4p$ $^1S_0$	21415.7967	0.1365D+14	0.26456D+14	0.46842D+15
$1s^23p4p$ $^1D_2$	21420.4174	0.8730D+14	0.10361D+15	0.23128D+15
$1s^23p4f$ $^3F_3$	21421.3630	0.1316D+15	0.63164D+12	0.70849D+13
$1s^23p4f$ $^3F_4$	21425.5760	0.8827D+14	0.20891D+13	0.31705D+13
$1s^23p4f$ $^3G_5$	21425.5899	0.1086D+15	0.20891D+13	0.79550D+13
$1s^23p4d$ $^3D_3$	21432.9045	0.2359D+15	0.13643D+13	0.15349D+14

Table II (continued)

state	energy(1000cm <sup>-1</sup> )	$\sum gAr(1/s)$	$Aa(1/s)$ to $2s$	$\sum Aa(1/s)$
$1s^2 3p4d$ $^1D_2$	21435.9849	0.1610D+15	0.22493D+13	0.14395D+14
$1s^2 3p4d$ $^3F_3$	21440.7867	0.1180D+15	0.23006D+13	0.30016D+13
$1s^2 3p4d$ $^1P_1$	21442.7602	0.5547D+14	0.19875D+11	0.10121D+14
$1s^2 3d4d$ $^3D_1$	21443.4922	0.6858D+14	0.53189D+11	0.72395D+12
$1s^2 3p4f$ $^3D_3$	21444.2548	0.8963D+14	0.28565D+10	0.35128D+13
$1s^2 3p4f$ $^3D_2$	21444.9671	0.8299D+14	0.13444D+12	0.21205D+13
$1s^2 3d4d$ $^3G_3$	21452.1622	0.1906D+15	0.87109D+12	0.13242D+14
$1s^2 3p4d$ $^3F_4$	21454.2805	0.1427D+15	0.26878D+13	0.30284D+13
$1s^2 3p4f$ $^1G_4$	21454.8069	0.1367D+15	0.78331D+14	0.23401D+15
$1s^2 3d4p$ $^1D_2$	21457.4129	0.1178D+15	0.69900D+12	0.61168D+13
$1s^2 3d4d$ $^3G_4$	21457.4337	0.2338D+15	0.84641D+13	0.43265D+14
$1s^2 3d4d$ $^1P_1$	21460.6963	0.6363D+14	0.49827D+11	0.36318D+13
$1s^2 3p4d$ $^3P_2$	21461.5009	0.7873D+14	0.48385D+11	0.22530D+14
$1s^2 3d4d$ $^3D_2$	21463.3507	0.9859D+14	0.66838D+12	0.12775D+14
$1s^2 3p4d$ $^3P_1$	21465.2142	0.4947D+14	0.10770D+12	0.27119D+14
$1s^2 3p4d$ $^3P_0$	21466.2223	0.1613D+14	0.33235D+10	0.17315D+14
$1s^2 3d4d$ $^3G_5$	21466.4644	0.2865D+15	0.12352D+13	0.15283D+14
$1s^2 3p4f$ $^3D_1$	21468.9240	0.5888D+14	0.17998D+12	0.65454D+13
$1s^2 3d4p$ $^3F_2$	21468.9648	0.1238D+15	0.45413D+12	0.39912D+13
$1s^2 3d4d$ $^3D_3$	21469.1768	0.1599D+15	0.21521D+12	0.83595D+13
$1s^2 3d4p$ $^3F_3$	21473.0062	0.1691D+15	0.95999D+12	0.50594D+13
$1s^2 3d4p$ $^3D_1$	21477.2908	0.6619D+14	0.84150D+11	0.74804D+14
$1s^2 3d4d$ $^3F_2$	21479.5310	0.1341D+15	0.27725D+12	0.97856D+14
$1s^2 3p4f$ $^1D_2$	21481.3849	0.6273D+14	0.70413D+13	0.26114D+14
$1s^2 3d4p$ $^3D_2$	21484.2893	0.1101D+15	0.15694D+12	0.77264D+14
$1s^2 3d4d$ $^3F_3$	21485.2719	0.1875D+15	0.18093D+11	0.96861D+14
$1s^2 3d4d$ $^3F_4$	21490.7478	0.2434D+15	0.15499D+12	0.10253D+15
$1s^2 3d4p$ $^3D_3$	21491.9659	0.1541D+15	0.12375D+13	0.89885D+14
$1s^2 3d4d$ $^3S_1$	21493.3534	0.7778D+14	0.69759D+12	0.36107D+13
$1s^2 3d4p$ $^3P_0$	21493.8888	0.2368D+14	0.14962D+13	0.43935D+14
$1s^2 3d4p$ $^3P_1$	21494.8835	0.7063D+14	0.14716D+13	0.46180D+14
$1s^2 3d4p$ $^3P_2$	21499.5008	0.1167D+15	0.15468D+13	0.39332D+14
$1s^2 3d4d$ $^3P_0$	21506.3436	0.2862D+14	0.32427D+12	0.17444D+14
$1s^2 3d4f$ $^3H_4$	21508.4313	0.1996D+15	0.90092D+13	0.43856D+14
$1s^2 3d4d$ $^3P_1$	21509.9802	0.8460D+14	0.24939D+11	0.14365D+14
$1s^2 3d4d$ $^3P_2$	21512.1569	0.1366D+15	0.35075D+12	0.22016D+14
$1s^2 3d4f$ $^3H_5$	21515.5351	0.2406D+15	0.11504D+14	0.52684D+14
$1s^2 3d4f$ $^1G_4$	21515.8876	0.1962D+15	0.24385D+13	0.21003D+14
$1s^2 3d4f$ $^3H_6$	21522.6505	0.2795D+15	0.11472D+14	0.52434D+14
$1s^2 3d4f$ $^3F_2$	21527.8767	0.1095D+15	0.63006D+10	0.96534D+13

Table II (continued)

state	energy(1000cm <sup>-1</sup> )	$\sum gAr(1/s)$	$Aa(1/s)$ to $2s$	$\sum Aa(1/s)$
$1s^2 3d4d \ ^1D_2$	21527.9896	0.1138D+15	0.23903D+13	0.12240D+15
$1s^2 3d4f \ ^3F_3$	21532.0578	0.1517D+15	0.14065D+13	0.18477D+14
$1s^2 3d4d \ ^1G_4$	21534.9278	0.1967D+15	0.15386D+14	0.40835D+15
$1s^2 3d4f \ ^3F_4$	21537.8367	0.1917D+15	0.54898D+11	0.10282D+14
$1s^2 3d4p \ ^1F_3$	21539.9946	0.1456D+15	0.60973D+14	0.34458D+15
$1s^2 3d4f \ ^1D_2$	21547.4547	0.1082D+15	0.31022D+11	0.13246D+13
$1s^2 3d4p \ ^1P_1$	21557.7249	0.5934D+14	0.23949D+13	0.17273D+15
$1s^2 3d4f \ ^3G_3$	21559.2303	0.1566D+15	0.28389D+09	0.28636D+14
$1s^2 3d4f \ ^3G_4$	21564.2750	0.1991D+15	0.21928D+11	0.28322D+14
$1s^2 3d4f \ ^3G_5$	21568.8199	0.2417D+15	0.17320D+12	0.29209D+14
$1s^2 3d4f \ ^3D_1$	21571.1528	0.6632D+14	0.41147D+11	0.92013D+12
$1s^2 3d4f \ ^3D_2$	21572.1256	0.1102D+15	0.31022D+11	0.40427D+12
$1s^2 3d4f \ ^3D_3$	21573.8669	0.1533D+15	0.22486D+12	0.24294D+13
$1s^2 3d4f \ ^3P_2$	21585.5318	0.1070D+15	0.23166D+12	0.13685D+13
$1s^2 3d4d \ ^1S_0$	21586.3556	0.2582D+14	0.25030D+14	0.82886D+14
$1s^2 3d4f \ ^3P_1$	21587.3875	0.6386D+14	0.22889D+12	0.13063D+13
$1s^2 3d4f \ ^3P_0$	21588.8011	0.2118D+14	0.22589D+12	0.12606D+13
$1s^2 3d4f \ ^1F_3$	21592.7565	0.1463D+15	0.49951D+13	0.42173D+14
$1s^2 3d4f \ ^1H_5$	21617.9305	0.2465D+15	0.45812D+14	0.22648D+15
$1s^2 3d4f \ ^1P_1$	21626.1895	0.6523D+14	0.42235D+12	0.22996D+13

**Table III** Comparison between our  $Ar$  (in units of  $1/s$ ) values from the  $1s^2 2l_1 n l_2$  ( $n = 2 \sim 5$ ) to the  $1s^2 2l_1 2l_2$  and those by others[17]. The first to fourth columns give the lower states, upper states, our results, and other results.

lower level	upper level	$Ar$ ( $1/s$ )	
		present work	Shirai et.al.[17]
$1s^2 2s^2 \ ^1S_0$	$1s^2 2s 2p \ ^3P_1$	0.62750D+08	0.48D+08
$1s^2 2s^2 \ ^1S_0$	$1s^2 2s 2p \ ^1P_1$	0.18009D+11	0.195D+11
$1s^2 2p^2 \ ^3P_1$	$1s^2 2s 3p \ ^3P_0$	0.39410D+11	
$1s^2 2s^2 \ ^1S_0$	$1s^2 2s 3p \ ^3P_1$	0.45013D+13	
$1s^2 2p^2 \ ^3P_0$	$1s^2 2s 3p \ ^3P_1$	0.85073D+10	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2s 3p \ ^3P_1$	0.75420D+10	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2s 3p \ ^3P_1$	0.81450D+11	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2s 3p \ ^3P_1$	0.38010D+11	
$1s^2 2p^2 \ ^1S_0$	$1s^2 2s 3p \ ^3P_1$	0.11835D+11	
$1s^2 2s^2 \ ^1S_0$	$1s^2 2s 3p \ ^1P_1$	0.88040D+13	0.114D+14
$1s^2 2p^2 \ ^3P_0$	$1s^2 2s 3p \ ^1P_1$	0.97747D+10	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2s 3p \ ^1P_1$	0.13012D+10	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2s 3p \ ^1P_1$	0.54940D+11	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2s 3p \ ^1P_1$	0.18626D+12	0.17D+12
$1s^2 2p^2 \ ^1S_0$	$1s^2 2s 3p \ ^1P_1$	0.67130D+10	0.14D+12
$1s^2 2p^2 \ ^3P_1$	$1s^2 2s 3p \ ^3P_2$	0.26000D+11	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2s 3p \ ^3P_2$	0.35698D+11	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2s 3p \ ^3P_2$	0.87492D+10	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2p 3s \ ^3P_0$	0.24014D+13	
$1s^2 2s^2 \ ^1S_0$	$1s^2 2p 3s \ ^3P_1$	0.25259D+12	
$1s^2 2p^2 \ ^3P_0$	$1s^2 2p 3s \ ^3P_1$	0.93063D+12	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2p 3s \ ^3P_1$	0.48010D+12	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p 3s \ ^3P_1$	0.12272D+13	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p 3s \ ^3P_1$	0.24562D+10	
$1s^2 2p^2 \ ^1S_0$	$1s^2 2p 3s \ ^3P_1$	0.35093D+11	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2p 3s \ ^3P_2$	0.99638D+12	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p 3s \ ^3P_2$	0.16671D+13	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p 3s \ ^3P_2$	0.48466D+12	
$1s^2 2s^2 \ ^1S_0$	$1s^2 2p 3s \ ^1P_1$	0.42457D+12	
$1s^2 2p^2 \ ^3P_0$	$1s^2 2p 3s \ ^1P_1$	0.30703D+11	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2p 3s \ ^1P_1$	0.12133D+12	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p 3s \ ^1P_1$	0.15005D+12	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p 3s \ ^1P_1$	0.21441D+13	
$1s^2 2p^2 \ ^1S_0$	$1s^2 2p 3s \ ^1P_1$	0.90287D+12	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2p 3d \ ^3F_2$	0.23496D+12	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p 3d \ ^3F_2$	0.17403D+13	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p 3d \ ^3F_2$	0.71622D+12	

Table III (continued)

lower level	upper level	Ar (1/s)	
		present work	Shirai et.al.[17]
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p3d \ ^3F_3$	0.74247D+13	0.77D+13
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p3d \ ^3F_3$	0.17793D+12	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2p3d \ ^3P_2$	0.13544D+14	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p3d \ ^3P_2$	0.17606D+13	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p3d \ ^3P_2$	0.15470D+13	
$1s^2 2s^2 \ ^1S_0$	$1s^2 2p3d \ ^3D_1$	0.10922D+12	
$1s^2 2p^2 \ ^3P_0$	$1s^2 2p3d \ ^3D_1$	0.22123D+14	0.24D+14
$1s^2 2p^2 \ ^3P_1$	$1s^2 2p3d \ ^3D_1$	0.43640D+13	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p3d \ ^3D_1$	0.23120D+12	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p3d \ ^3D_1$	0.32181D+11	
$1s^2 2p^2 \ ^1S_0$	$1s^2 2p3d \ ^3D_1$	0.34397D+12	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2p3d \ ^1D_2$	0.10499D+14	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p3d \ ^1D_2$	0.25624D+13	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p3d \ ^1D_2$	0.24662D+13	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p3d \ ^3D_3$	0.22769D+14	0.23D+14
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p3d \ ^3D_3$	0.18420D+13	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2p3d \ ^3P_2$	0.21318D+13	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p3d \ ^3P_2$	0.10886D+14	0.12D+14
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p3d \ ^3P_2$	0.81146D+13	
$1s^2 2s^2 \ ^1S_0$	$1s^2 2p3d \ ^3P_1$	0.44530D+10	
$1s^2 2p^2 \ ^3P_0$	$1s^2 2p3d \ ^3P_1$	0.28345D+11	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2p3d \ ^3P_1$	0.12427D+14	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p3d \ ^3P_1$	0.56513D+13	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p3d \ ^3P_1$	0.17278D+13	
$1s^2 2p^2 \ ^1S_0$	$1s^2 2p3d \ ^3P_1$	0.17935D+12	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2p3d \ ^3P_0$	0.16895D+14	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p3d \ ^1F_3$	0.26200D+13	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p3d \ ^1F_3$	0.34233D+14	0.35D+14
$1s^2 2s^2 \ ^1S_0$	$1s^2 2p3d \ ^1P_1$	0.32918D+12	
$1s^2 2p^2 \ ^3P_0$	$1s^2 2p3d \ ^1P_1$	0.38900D+12	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2p3d \ ^1P_1$	0.27450D+12	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p3d \ ^1P_1$	0.25140D+11	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p3d \ ^1P_1$	0.10741D+13	
$1s^2 2p^2 \ ^1S_0$	$1s^2 2p3d \ ^1P_1$	0.20658D+14	0.21D+14
$1s^2 2p^2 \ ^3P_1$	$1s^2 2s4p \ ^3P_0$	0.65185D+08	
$1s^2 2s^2 \ ^1S_0$	$1s^2 2s4p \ ^3P_1$	0.11206D+13	0.12D+13
$1s^2 2p^2 \ ^3P_0$	$1s^2 2s4p \ ^3P_1$	0.16136D+10	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2s4p \ ^3P_1$	0.41483D+06	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2s4p \ ^3P_1$	0.23862D+10	

Table III (continued)

lower level	upper level	Ar (1/s)	
		present work	Shirai et.al.[17]
$1s^2 2p^2 \ ^1D_2$	$1s^2 2s4p \ ^3P_1$	0.32016D+10	
$1s^2 2p^2 \ ^1S_0$	$1s^2 2s4p \ ^3P_1$	0.22287D+11	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2s4p \ ^3P_2$	0.74396D+09	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2s4p \ ^3P_2$	0.40136D+09	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2s4p \ ^3P_2$	0.18119D+08	
$1s^2 2s^2 \ ^1S_0$	$1s^2 2s4p \ ^1P_1$	0.48173D+13	0.497D+13
$1s^2 2p^2 \ ^3P_0$	$1s^2 2s4p \ ^1P_1$	0.14513D+11	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2s4p \ ^1P_1$	0.13887D+07	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2s4p \ ^1P_1$	0.84153D+10	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2s4p \ ^1P_1$	0.18430D+11	
$1s^2 2p^2 \ ^1S_0$	$1s^2 2s4p \ ^1P_1$	0.94127D+11	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2s4f \ ^3F_2$	0.65684D+07	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2s4f \ ^3F_2$	0.61082D+07	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2s4f \ ^3F_2$	0.11053D+07	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2s4f \ ^3F_3$	0.15627D+08	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2s4f \ ^3F_3$	0.18321D+09	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2s4f \ ^1F_3$	0.24299D+10	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2s4f \ ^1F_3$	0.66563D+10	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2p4s \ ^3P_0$	0.76285D+12	
$1s^2 2s^2 \ ^1S_0$	$1s^2 2p4s \ ^3P_1$	0.76753D+10	
$1s^2 2p^2 \ ^3P_0$	$1s^2 2p4s \ ^3P_1$	0.30755D+12	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2p4s \ ^3P_1$	0.12232D+12	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p4s \ ^3P_1$	0.49933D+12	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p4s \ ^3P_1$	0.50020D+11	
$1s^2 2p^2 \ ^1S_0$	$1s^2 2p4s \ ^3P_1$	0.42790D+11	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2p4d \ ^3F_2$	0.64170D+11	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p4d \ ^3F_2$	0.81412D+12	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p4d \ ^3F_2$	0.16255D+12	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2p4d \ ^3P_2$	0.53218D+13	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p4d \ ^3P_2$	0.90126D+12	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p4d \ ^3P_2$	0.18036D+11	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p4d \ ^3F_3$	0.45807D+13	0.46D+13
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p4d \ ^3F_3$	0.43109D+12	
$1s^2 2s^2 \ ^1S_0$	$1s^2 2p4d \ ^3D_1$	0.75490D+11	
$1s^2 2p^2 \ ^3P_0$	$1s^2 2p4d \ ^3D_1$	0.68883D+13	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2p4d \ ^3D_1$	0.11499D+13	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p4d \ ^3D_1$	0.13255D+12	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p4d \ ^3D_1$	0.38057D+11	



Table III (continued)

lower level	upper level	Ar (1/s)	
		present work	Shirai et.al.[17]
$1s^2 2p^2 \ ^1S_0$	$1s^2 2p4d \ ^3D_1$	0.49417D+12	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2p4s \ ^3P_2$	0.57448D+12	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p4s \ ^3P_2$	0.45474D+11	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p4s \ ^3P_2$	0.37270D+12	
$1s^2 2s^2 \ ^1S_0$	$1s^2 2p4s \ ^1P_1$	0.10472D+11	
$1s^2 2p^2 \ ^3P_0$	$1s^2 2p4s \ ^1P_1$	0.92397D+11	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2p4s \ ^1P_1$	0.39477D+11	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p4s \ ^1P_1$	0.31138D+11	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p4s \ ^1P_1$	0.86187D+12	
$1s^2 2p^2 \ ^1S_0$	$1s^2 2p4s \ ^1P_1$	0.30390D+12	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2p4d \ ^1D_2$	0.22182D+13	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p4d \ ^1D_2$	0.10463D+13	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p4d \ ^1D_2$	0.91442D+12	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p4d \ ^3D_3$	0.61799D+13	0.61D+13
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p4d \ ^3D_3$	0.32147D+12	
$1s^2 2s^2 \ ^1S_0$	$1s^2 2p4d \ ^3P_1$	0.21566D+10	
$1s^2 2p^2 \ ^3P_0$	$1s^2 2p4d \ ^3P_1$	0.16605D+10	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2p4d \ ^3P_1$	0.43667D+13	0.45D+13
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p4d \ ^3P_1$	0.18514D+13	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p4d \ ^3P_1$	0.55380D+12	
$1s^2 2p^2 \ ^1S_0$	$1s^2 2p4d \ ^3P_1$	0.91253D+11	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2p4d \ ^3P_2$	0.88178D+12	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p4d \ ^3P_2$	0.33218D+13	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p4d \ ^3P_2$	0.28580D+13	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2p4d \ ^3P_0$	0.56197D+13	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p4d \ ^1F_3$	0.22427D+12	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p4d \ ^1F_3$	0.11641D+14	0.12D+14
$1s^2 2s^2 \ ^1S_0$	$1s^2 2p4d \ ^1P_1$	0.14797D+12	
$1s^2 2p^2 \ ^3P_0$	$1s^2 2p4d \ ^1P_1$	0.46210D+12	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2p4d \ ^1P_1$	0.14091D+12	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p4d \ ^1P_1$	0.39353D+09	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p4d \ ^1P_1$	0.37863D+12	
$1s^2 2p^2 \ ^1S_0$	$1s^2 2p4d \ ^1P_1$	0.65910D+13	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2s5p \ ^3P_0$	0.12153D+10	
$1s^2 2s^2 \ ^1S_0$	$1s^2 2s5p \ ^3P_1$	0.49883D+12	
$1s^2 2p^2 \ ^3P_0$	$1s^2 2s5p \ ^3P_1$	0.36590D+10	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2s5p \ ^3P_1$	0.34203D+09	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2s5p \ ^3P_1$	0.17102D+09	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2s5p \ ^3P_1$	0.28798D+10	

Table III (continued)

		<i>Ar</i> (1/s)	
lower level	upper level	present work	Shirai et.al.[17]
$1s^2 2p^2 \ ^1S_0$	$1s^2 2s5p \ ^3P_1$	0.17190D+11	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2s5p \ ^3P_2$	0.20522D+08	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2s5p \ ^3P_2$	0.34108D+09	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2s5p \ ^3P_2$	0.25492D+09	
$1s^2 2s^2 \ ^1S_0$	$1s^2 2s5p \ ^1P_1$	0.25394D+13	0.25D+13
$1s^2 2p^2 \ ^3P_0$	$1s^2 2s5p \ ^1P_1$	0.11266D+11	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2s5p \ ^1P_1$	0.39893D+08	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2s5p \ ^1P_1$	0.56360D+10	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2s5p \ ^1P_1$	0.93230D+10	
$1s^2 2p^2 \ ^1S_0$	$1s^2 2s5p \ ^1P_1$	0.97000D+11	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2s5f \ ^3F_2$	0.24724D+04	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2s5f \ ^3F_2$	0.77086D+03	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2s5f \ ^3F_2$	0.12893D+04	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2s5f \ ^3F_3$	0.90606D+08	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2s5f \ ^3F_3$	0.25924D+09	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2s5f \ ^1F_3$	0.37033D+10	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2s5f \ ^1F_3$	0.10433D+11	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2p5s \ ^3P_0$	0.33996D+12	
$1s^2 2s^2 \ ^1S_0$	$1s^2 2p5s \ ^3P_1$	0.15987D+08	
$1s^2 2p^2 \ ^3P_0$	$1s^2 2p5s \ ^3P_1$	0.14325D+12	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2p5s \ ^3P_1$	0.49947D+11	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p5s \ ^3P_1$	0.24376D+12	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p5s \ ^3P_1$	0.37397D+11	
$1s^2 2p^2 \ ^1S_0$	$1s^2 2p5s \ ^3P_1$	0.23001D+11	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2p5d \ ^3F_2$	0.44344D+11	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p5d \ ^3F_2$	0.36790D+12	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p5d \ ^3F_2$	0.68190D+11	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2p5d \ ^3P_2$	0.27490D+13	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p5d \ ^3P_2$	0.27398D+12	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p5d \ ^3P_2$	0.47450D+11	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p5d \ ^3F_3$	0.24327D+13	0.23D+13
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p5d \ ^3F_3$	0.33859D+12	
$1s^2 2s^2 \ ^1S_0$	$1s^2 2p5d \ ^3D_1$	0.52653D+11	
$1s^2 2p^2 \ ^3P_0$	$1s^2 2p5d \ ^3D_1$	0.31276D+13	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2p5d \ ^3D_1$	0.47490D+12	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p5d \ ^3D_1$	0.70980D+11	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p5d \ ^3D_1$	0.17725D+11	
$1s^2 2p^2 \ ^1S_0$	$1s^2 2p5d \ ^3D_1$	0.30813D+12	

Table III (continued)

lower level	upper level	Ar (1/s)	
		present work	Shirai et.al.[17]
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p5g \ ^3G_3$	0.25691D+10	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p5g \ ^3G_3$	0.44254D+09	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2p5s \ ^3P_2$	0.43028D+11	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p5s \ ^3P_2$	0.18674D+12	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p5s \ ^3P_2$	0.10122D+12	
$1s^2 2s^2 \ ^1S_0$	$1s^2 2p5s \ ^1P_1$	0.28633D+09	
$1s^2 2p^2 \ ^3P_0$	$1s^2 2p5s \ ^1P_1$	0.21441D+11	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2p5s \ ^1P_1$	0.28649D+11	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p5s \ ^1P_1$	0.18668D+11	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p5s \ ^1P_1$	0.44113D+12	
$1s^2 2p^2 \ ^1S_0$	$1s^2 2p5s \ ^1P_1$	0.15931D+12	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2p5d \ ^1D_2$	0.92998D+12	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p5d \ ^1D_2$	0.50306D+12	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p5d \ ^1D_2$	0.40896D+12	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p5d \ ^3D_3$	0.26561D+13	0.25D+13
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p5d \ ^3D_3$	0.11569D+12	
$1s^2 2s^2 \ ^1S_0$	$1s^2 2p5d \ ^3P_1$	0.13889D+10	
$1s^2 2p^2 \ ^3P_0$	$1s^2 2p5d \ ^3P_1$	0.28861D+10	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2p5d \ ^3P_1$	0.20584D+13	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p5d \ ^3P_1$	0.85987D+12	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p5d \ ^3P_1$	0.25361D+12	
$1s^2 2p^2 \ ^1S_0$	$1s^2 2p5d \ ^3P_1$	0.43467D+11	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2p5d \ ^3P_2$	0.42936D+12	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p5d \ ^3P_2$	0.14963D+13	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p5d \ ^3P_2$	0.13470D+13	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2p5d \ ^3P_0$	0.26176D+13	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p5d \ ^1F_3$	0.39657D+11	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p5d \ ^1F_3$	0.53754D+13	0.49D+13
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p5g \ ^3G_3$	0.13432D+10	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p5g \ ^3G_3$	0.29051D+11	
$1s^2 2s^2 \ ^1S_0$	$1s^2 2p5d \ ^1P_1$	0.91713D+11	
$1s^2 2p^2 \ ^3P_0$	$1s^2 2p5d \ ^1P_1$	0.29566D+12	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2p5d \ ^1P_1$	0.79137D+11	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p5d \ ^1P_1$	0.16381D+09	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p5d \ ^1P_1$	0.19361D+12	
$1s^2 2p^2 \ ^1S_0$	$1s^2 2p5d \ ^1P_1$	0.29984D+13	
$1s^2 2p^2 \ ^3P_1$	$1s^2 2p5g \ ^3F_2$	0.28108D+09	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p5g \ ^3F_2$	0.25262D+09	
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p5g \ ^3F_3$	0.15181D+10	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p5g \ ^3F_3$	0.66453D+10	
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p5g \ ^3F_2$	0.13651D+08	

Table III (continued)

lower level	upper level	Ar (1/s)	
		present work	Shirai et.al.[17]
$1s^2 2s 2p \ ^3P_0$	$1s^2 2p^2 \ ^3P_1$	0.65150D+10	0.662D+10
$1s^2 2s 2p \ ^3P_0$	$1s^2 2s 3s \ ^3S_1$	0.43203D+12	
$1s^2 2s 2p \ ^3P_0$	$1s^2 2s 3d \ ^3D_1$	0.13321D+14	0.13D+14
$1s^2 2s 2p \ ^3P_0$	$1s^2 2p 3p \ ^3D_1$	0.11430D+13	
$1s^2 2s 2p \ ^3P_0$	$1s^2 2p 3p \ ^3S_1$	0.35763D+13	
$1s^2 2s 2p \ ^3P_0$	$1s^2 2p 3p \ ^3P_1$	0.39743D+12	
$1s^2 2s 2p \ ^3P_0$	$1s^2 2p 3p \ ^3S_1$	0.31173D+11	
$1s^2 2s 2p \ ^3P_0$	$1s^2 2s 4s \ ^3S_1$	0.14467D+12	
$1s^2 2s 2p \ ^3P_0$	$1s^2 2s 4d \ ^3D_1$	0.40397D+13	0.40D+13
$1s^2 2s 2p \ ^3P_0$	$1s^2 2p 4p \ ^3D_1$	0.81303D+12	
$1s^2 2s 2p \ ^3P_0$	$1s^2 2p 4p \ ^3P_1$	0.18903D+13	
$1s^2 2s 2p \ ^3P_0$	$1s^2 2p 4p \ ^1P_1$	0.13789D+11	
$1s^2 2s 2p \ ^3P_0$	$1s^2 2p 4p \ ^3S_1$	0.10305D+05	
$1s^2 2s 2p \ ^3P_0$	$1s^2 2p 4f \ ^3D_1$	0.55650D+08	
$1s^2 2s 2p \ ^3P_0$	$1s^2 2s 5s \ ^3S_1$	0.61277D+11	
$1s^2 2s 2p \ ^3P_0$	$1s^2 2s 5d \ ^3D_1$	0.18019D+13	0.18D+13
$1s^2 2s 2p \ ^3P_0$	$1s^2 2p 5p \ ^3D_1$	0.45323D+12	
$1s^2 2s 2p \ ^3P_0$	$1s^2 2p 5p \ ^3P_1$	0.10033D+13	
$1s^2 2s 2p \ ^3P_0$	$1s^2 2p 5p \ ^1P_1$	0.82843D+09	
$1s^2 2s 2p \ ^3P_0$	$1s^2 2p 5p \ ^3S_1$	0.47687D+09	
$1s^2 2s 2p \ ^3P_0$	$1s^2 2p 5f \ ^3D_1$	0.57383D+09	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2p^2 \ ^3P_0$	0.12315D+11	0.124D+11
$1s^2 2s 2p \ ^3P_1$	$1s^2 2p^2 \ ^3P_1$	0.41270D+10	0.419D+10
$1s^2 2s 2p \ ^3P_1$	$1s^2 2p^2 \ ^3P_2$	0.53354D+10	0.546D+10
$1s^2 2s 2p \ ^3P_1$	$1s^2 2p^2 \ ^1D_2$	0.49374D+09	0.44D+9
$1s^2 2s 2p \ ^3P_1$	$1s^2 2p^2 \ ^1S_0$	0.18458D+09	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2s 3s \ ^3S_1$	0.12434D+13	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2s 3s \ ^1S_0$	0.26550D+11	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2s 3d \ ^3D_1$	0.95917D+13	0.93D+13
$1s^2 2s 2p \ ^3P_1$	$1s^2 2s 3d \ ^3D_2$	0.17434D+14	0.17D+14
$1s^2 2s 2p \ ^3P_1$	$1s^2 2s 3d \ ^1D_2$	0.14662D+12	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2p 3p \ ^3D_2$	0.46192D+13	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2p 3p \ ^3S_1$	0.11360D+13	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2p 3p \ ^3P_0$	0.65580D+13	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2p 3p \ ^3P_1$	0.19002D+13	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2p 3p \ ^3S_1$	0.12437D+12	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2p 3p \ ^3P_2$	0.25912D+12	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2p 3p \ ^1D_2$	0.25248D+12	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2p 3p \ ^3D_1$	0.22618D+13	

Table III (continued)

lower level	upper level	Ar (1/s)	
		present work	Shirai et.al.[17]
$1s^2 2s 2p \ ^3P_1$	$1s^2 2p 3p \ ^1S_0$	0.24966D+12	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2s 4s \ ^3S_1$	0.41493D+12	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2s 4s \ ^1S_0$	0.18496D+11	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2s 4d \ ^3D_1$	0.29131D+13	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2s 4d \ ^3D_2$	0.53824D+13	0.53D+13
$1s^2 2s 2p \ ^3P_1$	$1s^2 2s 4d \ ^1D_2$	0.36208D+11	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2p 4p \ ^3D_1$	0.13835D+13	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2p 4p \ ^3P_1$	0.76180D+12	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2p 4p \ ^3D_2$	0.24070D+13	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2p 4p \ ^3P_0$	0.26936D+13	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2p 4f \ ^3F_2$	0.19260D+10	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2p 4p \ ^1P_1$	0.45630D+12	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2p 4p \ ^3P_2$	0.19800D+11	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2p 4p \ ^3S_1$	0.18424D+12	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2p 4p \ ^1D_2$	0.26512D+12	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2p 4p \ ^1S_0$	0.31006D+12	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2p 4f \ ^3F_2$	0.21966D+08	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2p 4f \ ^3D_1$	0.75213D+07	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2p 4f \ ^1D_2$	0.87096D+08	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2s 5s \ ^3S_1$	0.17552D+12	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2s 5s \ ^1S_0$	0.12118D+11	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2s 5d \ ^3D_1$	0.13017D+13	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2s 5d \ ^3D_2$	0.24172D+13	0.25D+13
$1s^2 2s 2p \ ^3P_1$	$1s^2 2s 5d \ ^1D_2$	0.14242D+11	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2p 5p \ ^3D_1$	0.78273D+12	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2p 5p \ ^3P_1$	0.43123D+12	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2p 5p \ ^3P_0$	0.12761D+13	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2p 5p \ ^3D_2$	0.12330D+13	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2p 5f \ ^3F_2$	0.19173D+10	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2p 5p \ ^1P_1$	0.17815D+12	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2p 5p \ ^3D_2$	0.41970D+11	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2p 5p \ ^3S_1$	0.12604D+12	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2p 5p \ ^1D_2$	0.16362D+12	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2p 5p \ ^1S_0$	0.22921D+12	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2p 5f \ ^3F_2$	0.22424D+08	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2p 5f \ ^3D_1$	0.25570D+09	
$1s^2 2s 2p \ ^3P_1$	$1s^2 2p 5f \ ^1D_2$	0.61180D+09	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p^2 \ ^3P_1$	0.44810D+10	0.451D+10
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p^2 \ ^3P_2$	0.73242D+10	0.756D+10

Table III (continued)

lower level	upper level	Ar (1/s)	
		present work	Shirai et.al.[17]
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p^2 \ ^1D_2$	0.51046D+10	0.494D+10
$1s^2 2s 2p \ ^3P_2$	$1s^2 2s 3s \ ^3S_1$	0.19960D+13	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2s 3d \ ^3D_1$	0.62090D+12	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2s 3d \ ^3D_2$	0.55562D+13	0.54D+13
$1s^2 2s 2p \ ^3P_2$	$1s^2 2s 3d \ ^3D_3$	0.22283D+14	0.22D+14
$1s^2 2s 2p \ ^3P_2$	$1s^2 2s 3d \ ^1D_2$	0.26904D+11	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 3p \ ^3D_1$	0.17301D+09	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 3p \ ^3D_2$	0.10883D+12	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 3p \ ^3S_1$	0.63077D+11	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 3p \ ^3P_1$	0.70627D+12	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 3p \ ^3D_3$	0.49026D+13	0.52D+13
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 3p \ ^3S_1$	0.56037D+13	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 3p \ ^3P_2$	0.47776D+13	0.49D+13
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 3p \ ^1D_2$	0.14174D+13	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2s 4s \ ^3S_1$	0.67060D+12	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2s 4d \ ^3D_1$	0.19322D+12	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2s 4d \ ^3D_2$	0.17210D+13	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2s 4d \ ^3D_3$	0.69559D+13	0.71D+13
$1s^2 2s 2p \ ^3P_2$	$1s^2 2s 4d \ ^1D_2$	0.17841D+11	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 4p \ ^3D_1$	0.92477D+06	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 4p \ ^3P_1$	0.66080D+09	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 4p \ ^3D_2$	0.17987D+10	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 4f \ ^3G_3$	0.16843D+08	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 4f \ ^3F_2$	0.43734D+09	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 4f \ ^3D_3$	0.11601D+10	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 4p \ ^1P_1$	0.85360D+12	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 4p \ ^3P_2$	0.21182D+13	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 4p \ ^3D_3$	0.26241D+13	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 4p \ ^3S_1$	0.21190D+13	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 4p \ ^1D_2$	0.78892D+12	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 4f \ ^1F_3$	0.21033D+06	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 4f \ ^3F_2$	0.35228D+08	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 4f \ ^3D_3$	0.14248D+09	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 4f \ ^3D_1$	0.16450D+07	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 4f \ ^1D_2$	0.21796D+09	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2s 5s \ ^3S_1$	0.28416D+12	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2s 5d \ ^3D_1$	0.86960D+11	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2s 5d \ ^3D_2$	0.77226D+12	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2s 5d \ ^3D_3$	0.31311D+13	0.30D+13
$1s^2 2s 2p \ ^3P_2$	$1s^2 2s 5d \ ^1D_2$	0.10426D+11	

Table III (continued)

lower level	upper level	$Ar$ (1/s)	
		present work	Shirai et.al.[17]
$1s^2 2s 2p \ ^3P_2$	$1s^2 2s 5g \ ^3G_3$	0.17193D+00	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 5p \ ^3D_1$	0.66447D+08	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 5p \ ^3P_1$	0.10119D+10	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 5p \ ^3D_2$	0.85984D+08	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 5f \ ^3G_3$	0.15224D+08	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 5f \ ^3F_2$	0.12218D+09	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 5f \ ^3D_3$	0.91756D+09	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 5p \ ^1P_1$	0.54833D+12	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 5p \ ^3D_2$	0.10660D+13	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 5p \ ^3D_3$	0.14294D+13	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 5p \ ^3S_1$	0.10581D+13	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 5p \ ^1D_2$	0.42284D+12	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 5f \ ^1F_3$	0.40004D+04	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 5f \ ^3F_2$	0.66254D+08	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 5f \ ^3D_3$	0.11437D+06	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 5f \ ^3D_1$	0.23106D+07	
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 5f \ ^1D_2$	0.48682D+08	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2p^2 \ ^3P_0$	0.33120D+08	0.21D8
$1s^2 2s 2p \ ^1P_1$	$1s^2 2p^2 \ ^3P_1$	0.12097D+08	0.70D+7
$1s^2 2s 2p \ ^1P_1$	$1s^2 2p^2 \ ^3P_2$	0.51678D+09	0.35D+9
$1s^2 2s 2p \ ^1P_1$	$1s^2 2p^2 \ ^1D_2$	0.50452D+10	0.454D+10
$1s^2 2s 2p \ ^1P_1$	$1s^2 2p^2 \ ^1S_0$	0.29805D+11	0.328D+11
$1s^2 2s 2p \ ^1P_1$	$1s^2 2s 3 \ ^3S_1$	0.25503D+11	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2s 3s \ ^1S_0$	0.13143D+13	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2s 3d \ ^3D_1$	0.20772D+12	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2s 3d \ ^3D_2$	0.12016D+12	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2s 3d \ ^1D_2$	0.16663D+14	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2p 3p \ ^3D_1$	0.13310D+13	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2p 3p \ ^3D_2$	0.43088D+12	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2p 3p \ ^3S_1$	0.14316D+13	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2p 3p \ ^3P_0$	0.26355D+12	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2p 3p \ ^3P_1$	0.36493D+13	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2p 3p \ ^3S_1$	0.66923D+12	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2p 3p \ ^3P_2$	0.21890D+13	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2p 3p \ ^1D_2$	0.68004D+13	0.68D+13
$1s^2 2s 2p \ ^1P_1$	$1s^2 2p 3p \ ^1S_0$	0.65685D+13	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2s 4s \ ^3S_1$	0.92620D+10	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2s 4s \ ^1S_0$	0.76997D+12	0.16D+13
$1s^2 2s 2p \ ^1P_1$	$1s^2 2s 4d \ ^3D_1$	0.74093D+11	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2s 4d \ ^3D_2$	0.13811D+11	

Table III (continued)

lower level	upper level	Ar (1/s)	
		present work	Shirai et.al.[17]
$1s^2 2s 2p \ ^1P_1$	$1s^2 2s 4d \ ^1D_2$	0.60714D+13	0.62D+13
$1s^2 2s 2p \ ^1P_1$	$1s^2 2p 4p \ ^3D_1$	0.46317D+12	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2p 4p \ ^3P_1$	0.29244D+12	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2p 4p \ ^3D_2$	0.45850D+12	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2p 4p \ ^3P_0$	0.28077D+12	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2p 4f \ ^3F_2$	0.47602D+09	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2p 4p \ ^1P_1$	0.16788D+13	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2p 4p \ ^3F_2$	0.82858D+12	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2p 4p \ ^3S_1$	0.70557D+12	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2p 4p \ ^1D_2$	0.21972D+13	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2p 4p \ ^1S_0$	0.25333D+13	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2p 4f \ ^3F_2$	0.82368D+10	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2p 4f \ ^3D_1$	0.43453D+08	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2p 4f \ ^1D_2$	0.15022D+11	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2s 5s \ ^3S_1$	0.39717D+10	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2s 5s \ ^1S_0$	0.48107D+12	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2s 5d \ ^3D_1$	0.34270D+11	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2s 5d \ ^3D_2$	0.26088D+10	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2s 5d \ ^1D_2$	0.29410D+13	0.28D+13
$1s^2 2s 2p \ ^1P_1$	$1s^2 2p 5p \ ^3D_1$	0.24119D+12	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2p 5p \ ^3P_1$	0.13066D+12	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2p 5p \ ^3P_0$	0.16437D+12	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2p 5p \ ^3D_2$	0.28312D+12	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2p 5f \ ^3F_2$	0.59628D+09	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2p 5p \ ^1P_1$	0.86923D+12	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2p 5p \ ^3D_2$	0.40378D+12	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2p 5p \ ^3S_1$	0.49663D+12	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2p 5p \ ^1D_2$	0.10079D+13	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2p 5p \ ^1S_0$	0.10487D+13	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2p 5f \ ^3F_2$	0.49854D+10	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2p 5f \ ^3D_1$	0.39707D+08	
$1s^2 2s 2p \ ^1P_1$	$1s^2 2p 5f \ ^1D_2$	0.11817D+11	



**Table IV**  $\lambda(\text{\AA})$ ,  $gAr(1/s)$ ,  $\sum Ar$ ,  $Aa(1/s)$ , and  $Qd$  from the upper level to the lower level.

lower level	upper level	$\lambda(\text{\AA})$	$gAr(1/s)$	$\sum Ar$	$Aa(1/s)$	$Qd$
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 11p \ ^3D_2$	6.53321	0.53620D+12	0.27382D+12	0.23866D+13	0.48101D+12
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 11p \ ^3D_3$	6.53311	0.10379D+13	0.27445D+12	0.99405D+12	0.81334D+12
$1s^2 2s 2p \ ^1P_1$	$1s^2 2p 11p \ ^3P_2$	6.64319	0.39400D+12	0.27210D+12	0.12840D+13	0.32511D+12
$1s^2 2p 3p \ ^3D_3$	$1s^2 2p 11d \ ^3F_4$	16.23643	0.82234D+12	0.19861D+12	0.24760D+12	0.45631D+12
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p 11d \ ^3D_3$	6.79835	0.13944D+13	0.39938D+12	0.37899D+12	0.67893D+12
$1s^2 2p^2 \ ^3P_1$	$1s^2 2p 11d \ ^3P_1$	6.77753	0.50007D+12	0.44984D+12	0.55553D+13	0.46261D+12
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p 11d \ ^3P_2$	6.79822	0.57323D+12	0.44936D+12	0.76363D+13	0.54137D+12
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p 11d \ ^3P_2$	6.85731	0.55557D+12	0.44936D+12	0.76363D+13	0.52469D+12
$1s^2 2p 3p \ ^1D_2$	$1s^2 2p 11d \ ^1F_3$	16.43888	0.45401D+12	0.57831D+12	0.31448D+13	0.38349D+12
$1s^2 2p^2 \ ^1S_0$	$1s^2 2p 11d \ ^1P_1$	6.94917	0.60853D+12	0.44017D+12	0.59719D+13	0.56676D+12
$1s^2 2p 3d \ ^1D_2$	$1s^2 2p 11f \ ^1F_3$	16.51326	0.43091D+12	0.21409D+12	0.12016D+13	0.36575D+12
$1s^2 2p 3d \ ^3D_3$	$1s^2 2p 11f \ ^3F_4$	16.57587	0.54286D+12	0.21122D+12	0.17260D+13	0.48367D+12
$1s^2 2p 4d \ ^3D_3$	$1s^2 2p 11f \ ^3F_4$	31.47646	0.33161D+12	0.21122D+12	0.17260D+13	0.29545D+12
$1s^2 2p 3d \ ^3P_2$	$1s^2 2p 11f \ ^3D_3$	16.62439	0.34845D+12	0.20852D+12	0.57809D+12	0.25608D+12
$1s^2 2p 3d \ ^3F_4$	$1s^2 2p 11f \ ^3G_5$	16.49139	0.97964D+12	0.22258D+12	0.47561D+13	0.93584D+12
$1s^2 2p 4d \ ^3F_4$	$1s^2 2p 11f \ ^3G_5$	31.40501	0.54667D+12	0.22258D+12	0.47561D+13	0.52223D+12
$1s^2 2p 5d \ ^3F_4$	$1s^2 2p 11f \ ^3G_5$	53.76820	0.31581D+12	0.22258D+12	0.47561D+13	0.30169D+12
$1s^2 2p 3d \ ^1F_3$	$1s^2 2p 11f \ ^1G_4$	16.77695	0.49172D+12	0.19075D+12	0.55143D+13	0.47528D+12
$1s^2 2p 4d \ ^1F_3$	$1s^2 2p 11f \ ^1G_4$	31.71539	0.34934D+12	0.19075D+12	0.55143D+13	0.33766D+12
$1s^2 2p 4f \ ^3F_4$	$1s^2 2p 11g \ ^3G_5$	31.64046	0.30657D+12	0.12286D+12	0.18329D+13	0.28731D+12
$1s^2 2p 5f \ ^3F_4$	$1s^2 2p 11g \ ^3G_5$	54.10618	0.25716D+12	0.12286D+12	0.18329D+13	0.24101D+12
$1s^2 2p 4f \ ^1G_4$	$1s^2 2p 11g \ ^1H_5$	31.70848	0.32203D+12	0.12106D+12	0.64105D+13	0.31606D+12
$1s^2 2p 5f \ ^1G_4$	$1s^2 2p 11g \ ^1H_5$	54.20768	0.27613D+12	0.12106D+12	0.64105D+13	0.27101D+12
$1s^2 2p 4f \ ^3G_5$	$1s^2 2p 11g \ ^3H_6$	31.68935	0.40649D+12	0.12232D+12	0.64023D+13	0.39887D+12
$1s^2 2p 5f \ ^3G_5$	$1s^2 2p 11g \ ^3H_6$	54.16296	0.35144D+12	0.12232D+12	0.64023D+13	0.34485D+12
$1s^2 2p 6f \ ^3G_5$	$1s^2 2p 11g \ ^3H_6$	88.10639	0.25807D+12	0.12232D+12	0.64023D+13	0.25323D+12
$1s^2 2p 6g \ ^3H_6$	$1s^2 2p 11h \ ^3I_7$	88.21890	0.25987D+12	0.66703D+11	0.60226D+13	0.25702D+12
$1s^2 2s 2p \ ^3P_1$	$1s^2 2p 13p \ ^3D_2$	6.48926	0.36985D+12	0.32783D+12	0.18626D+13	0.31450D+12
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 13p \ ^3D_2$	6.47494	0.32148D+12	0.32783D+12	0.18626D+13	0.27336D+12
$1s^2 2s 2p \ ^3P_2$	$1s^2 2p 13p \ ^3D_3$	6.47488	0.62823D+12	0.16714D+12	0.81151D+12	0.52094D+12
$1s^2 2p^2 \ ^3P_1$	$1s^2 2p 13d \ ^3P_2$	6.77246	0.65980D+12	0.53030D+12	0.49550D+13	0.59601D+12
$1s^2 2p 3p \ ^3S_1$	$1s^2 2p 13d \ ^3P_2$	15.95324	0.21184D+12	0.53030D+12	0.49550D+13	0.19136D+12
$1s^2 2p 3p \ ^3D_2$	$1s^2 2p 13d \ ^3F_3$	15.95072	0.31147D+12	0.25009D+12	0.78846D+12	0.23647D+12
$1s^2 2p^2 \ ^3P_0$	$1s^2 2p 13d \ ^3D_1$	6.74008	0.42727D+12	0.29409D+12	0.24440D+13	0.38138D+12
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p 13d \ ^3D_3$	6.73569	0.82744D+12	0.23975D+12	0.15768D+12	0.32828D+12
$1s^2 2p^2 \ ^3P_1$	$1s^2 2p 13d \ ^3P_1$	6.71529	0.30079D+12	0.27169D+12	0.35416D+13	0.27936D+12
$1s^2 2p^2 \ ^3P_2$	$1s^2 2p 13d \ ^3P_2$	6.73561	0.34233D+12	0.53030D+12	0.49550D+13	0.30923D+12
$1s^2 2p^2 \ ^1D_2$	$1s^2 2p 13d \ ^3P_2$	6.79361	0.33442D+12	0.53030D+12	0.49550D+13	0.30209D+12
$1s^2 2p 3p \ ^1D_2$	$1s^2 2p 13d \ ^1F_3$	16.07816	0.27616D+12	0.34803D+12	0.15395D+13	0.22524D+12
$1s^2 2p^2 \ ^1S_0$	$1s^2 2p 13d \ ^1P_1$	6.88391	0.36528D+12	0.26638D+12	0.37856D+13	0.34127D+12
$1s^2 2p 3d \ ^3F_2$	$1s^2 2p 13f \ ^3G_3$	16.10918	0.33400D+12	0.12880D+12	0.17220D+13	0.31076D+12

Table IV (continue)

lower level	upper level	$\lambda(\text{\AA})$	$gAr(1/s)$	$\sum Ar$	$Aa(1/s)$	$Qd$
$1s^2 2p 3s \ ^1P_1$	$1s^2 3s^2 \ ^1S_0$	11.64241	0.27587D+13	0.58250D+13	0.13576D+15	0.26397D+13
$1s^2 2s 3s \ ^3S_1$	$1s^2 3s 3p \ ^3P_0$	10.78845	0.76268D+13	0.10088D+14	0.58203D+14	0.33314D+13
$1s^2 2s 3s \ ^3S_1$	$1s^2 3s 3p \ ^3P_1$	10.77769	0.22533D+14	0.10173D+14	0.60554D+14	0.96497D+13
$1s^2 2s 3s \ ^3S_1$	$1s^2 3s 3p \ ^3P_2$	10.74770	0.38748D+14	0.10277D+14	0.57618D+14	0.16464D+14
$1s^2 2p 3p \ ^3D_3$	$1s^2 3s 3p \ ^3P_2$	11.63944	0.62247D+13	0.10277D+14	0.57618D+14	0.26449D+13
$1s^2 2s 3s \ ^1S_0$	$1s^2 3s 3p \ ^1P_1$	10.72000	0.20011D+14	0.12343D+14	0.17226D+15	0.69897D+13
$1s^2 2p 3p \ ^3S_1$	$1s^2 3s 3p \ ^1P_1$	11.54408	0.94524D+12	0.12343D+14	0.17226D+15	0.33016D+12
$1s^2 2s 3p \ ^3P_1$	$1s^2 3s 3d \ ^1D_2$	10.82257	0.20443D+14	0.20819D+14	0.10646D+15	0.86739D+13
$1s^2 2s 3p \ ^1P_1$	$1s^2 3s 3d \ ^1D_2$	10.86036	0.33101D+14	0.20819D+14	0.10646D+15	0.14045D+14
$1s^2 2p 3s \ ^3P_1$	$1s^2 3s 3d \ ^1D_2$	11.17119	0.15455D+14	0.20819D+14	0.10646D+15	0.65575D+13
$1s^2 2p 3s \ ^1P_1$	$1s^2 3s 3d \ ^1D_2$	11.35902	0.26705D+14	0.20819D+14	0.10646D+15	0.11331D+14
$1s^2 2p 3s \ ^3P_0$	$1s^2 3s 3d \ ^3D_1$	11.12775	0.37358D+14	0.24550D+14	0.21147D+14	0.12880D+14
$1s^2 2p 3s \ ^3P_1$	$1s^2 3s 3d \ ^3D_1$	11.14762	0.22777D+14	0.24550D+14	0.21147D+14	0.78529D+13
$1s^2 2p 3s \ ^1P_1$	$1s^2 3s 3d \ ^3D_1$	11.33465	0.46923D+13	0.24550D+14	0.21147D+14	0.16178D+13
$1s^2 2p 3s \ ^3P_1$	$1s^2 3s 3d \ ^3D_2$	11.14096	0.60706D+14	0.24018D+14	0.22802D+14	0.20319D+14
$1s^2 2p 3s \ ^3P_2$	$1s^2 3s 3d \ ^3D_2$	11.27583	0.25929D+14	0.24018D+14	0.22802D+14	0.86786D+13
$1s^2 2p 3s \ ^1P_1$	$1s^2 3s 3d \ ^3D_2$	11.32777	0.21313D+14	0.24018D+14	0.22802D+14	0.71336D+13
$1s^2 2p 3s \ ^3P_2$	$1s^2 3s 3d \ ^3D_3$	11.26793	0.15277D+15	0.23875D+14	0.21145D+14	0.53256D+14
$1s^2 2p 3d \ ^3F_4$	$1s^2 3s 3d \ ^3D_3$	11.59170	0.76039D+13	0.23875D+14	0.21145D+14	0.26507D+13
$1s^2 2s 3p \ ^3P_2$	$1s^2 3p^2 \ ^3P_2$	10.77166	0.52951D+14	0.16436D+14	0.12826D+14	0.26486D+13
$1s^2 2s 3p \ ^3P_1$	$1s^2 3p^2 \ ^1D_2$	10.64450	0.48592D+13	0.22096D+14	0.16799D+15	0.18111D+13
$1s^2 2s 3p \ ^1P_1$	$1s^2 3p^2 \ ^1D_2$	10.68105	0.16534D+14	0.22096D+14	0.16799D+15	0.61623D+13
$1s^2 2p 3s \ ^3P_1$	$1s^2 3p^2 \ ^1D_2$	10.98156	0.95042D+13	0.22096D+14	0.16799D+15	0.35423D+13
$1s^2 2p 3s \ ^1P_1$	$1s^2 3p^2 \ ^1D_2$	11.16301	0.40016D+14	0.22096D+14	0.16799D+15	0.14914D+14
$1s^2 2p 3d \ ^3F_2$	$1s^2 3p^2 \ ^1D_2$	11.25592	0.74537D+13	0.22096D+14	0.16799D+15	0.27780D+13
$1s^2 2p 3d \ ^3P_2$	$1s^2 3p^2 \ ^1D_2$	11.32129	0.86539D+13	0.22096D+14	0.16799D+15	0.32254D+13
$1s^2 2p 3d \ ^1D_2$	$1s^2 3p^2 \ ^1D_2$	11.43741	0.14248D+14	0.22096D+14	0.16799D+15	0.53103D+13
$1s^2 2p 3d \ ^1P_1$	$1s^2 3p^2 \ ^1D_2$	11.57627	0.43870D+13	0.22096D+14	0.16799D+15	0.16351D+13
$1s^2 2p 3p \ ^3D_2$	$1s^2 3p 3d \ ^3F_3$	11.20489	0.86020D+14	0.28834D+14	0.66003D+12	0.15376D+13
$1s^2 2s 3d \ ^1D_2$	$1s^2 3p 3d \ ^1F_3$	10.73085	0.30605D+14	0.28676D+14	0.94980D+14	0.44019D+13
$1s^2 2p 3p \ ^3D_2$	$1s^2 3p 3d \ ^1F_3$	11.03358	0.19829D+14	0.28676D+14	0.94980D+14	0.28520D+13
$1s^2 2p 3p \ ^3P_2$	$1s^2 3p 3d \ ^1F_3$	11.17672	0.23984D+14	0.28676D+14	0.94980D+14	0.34496D+13
$1s^2 2p 3p \ ^1D_2$	$1s^2 3p 3d \ ^1F_3$	11.25173	0.12425D+15	0.28676D+14	0.94980D+14	0.17871D+14
$1s^2 2p 3d \ ^1P_1$	$1s^2 3d^2 \ ^1S_0$	11.16419	0.31609D+14	0.37593D+14	0.31480D+14	0.75685D+13
$1s^2 2p 3d \ ^3D_3$	$1s^2 3d^2 \ ^3F_4$	11.31322	0.23631D+15	0.42332D+14	0.19288D+13	0.18741D+13
$1s^2 2p 3d \ ^3F_3$	$1s^2 3d^2 \ ^1G_4$	11.10128	0.29946D+14	0.41801D+14	0.19099D+15	0.43014D+13
$1s^2 2p 3d \ ^1F_3$	$1s^2 3d^2 \ ^1G_4$	11.34755	0.33778D+15	0.41801D+14	0.19099D+15	0.48518D+14
$1s^2 2s 4s \ ^3S_1$	$1s^2 3p 4s \ ^3P_1$	10.68504	0.18854D+14	0.93702D+13	0.15348D+14	0.47891D+13
$1s^2 2s 4s \ ^3S_1$	$1s^2 3p 4s \ ^3P_2$	10.65602	0.35183D+14	0.98948D+13	0.65432D+13	0.58360D+13
$1s^2 2s 4s \ ^1S_0$	$1s^2 3p 4s \ ^1P_1$	10.65425	0.16148D+14	0.98335D+13	0.84617D+14	0.52224D+13
$1s^2 2s 4p \ ^3P_1$	$1s^2 3d 4s \ ^1D_2$	10.67874	0.17982D+14	0.16448D+14	0.11317D+14	0.21713D+13
$1s^2 2p 4s \ ^1P_1$	$1s^2 3d 4s \ ^1D_2$	11.25082	0.39870D+14	0.16448D+14	0.11317D+14	0.48142D+13
$1s^2 2p 4p \ ^3D_2$	$1s^2 3d 4p \ ^3F_3$	11.07282	0.79744D+14	0.24159D+14	0.95999D+12	0.26200D+13
$1s^2 2p 3p \ ^1D_2$	$1s^2 3d 4p \ ^1F_3$	8.44304	0.13638D+14	0.20795D+14	0.60973D+14	0.22759D+13

Table IV (continue)

lower level	upper level	$\lambda(\text{\AA})$	$gAr(1/s)$	$\sum Ar$	$Aa(1/s)$	$Qd$
$1s^2 2s 4d \ ^1D_2$	$1s^2 3d 4p \ ^1F_3$	10.58853	0.14773D+14	0.20795D+14	0.60973D+14	0.24653D+13
$1s^2 2p 4p \ ^3D_2$	$1s^2 3d 4p \ ^1F_3$	10.99129	0.16451D+14	0.20795D+14	0.60973D+14	0.27453D+13
$1s^2 2p 4p \ ^3P_2$	$1s^2 3d 4p \ ^1F_3$	11.12894	0.14669D+14	0.20795D+14	0.60973D+14	0.24479D+13
$1s^2 2p 4p \ ^3D_3$	$1s^2 3d 4p \ ^3F_4$	11.19866	0.16395D+15	0.23821D+14	0.90436D+12	0.51172D+13
$1s^2 2p 4f \ ^3G_3$	$1s^2 3s 4f \ ^3F_2$	11.31927	0.27839D+13	0.31499D+13	0.55761D+13	0.16751D+13
$1s^2 2p 4f \ ^3F_2$	$1s^2 3s 4f \ ^3F_2$	11.48581	0.34387D+13	0.31499D+13	0.55761D+13	0.20691D+13
$1s^2 2p 4f \ ^3F_4$	$1s^2 3s 4f \ ^3F_3$	11.47949	0.42123D+13	0.30387D+13	0.57989D+13	0.25993D+13
$1s^2 2p 4f \ ^3D_3$	$1s^2 3s 4f \ ^3F_3$	11.48494	0.28362D+13	0.30387D+13	0.57989D+13	0.17502D+13
$1s^2 2p 4f \ ^3D_3$	$1s^2 3s 4f \ ^3F_4$	11.32041	0.35584D+13	0.29326D+13	0.60249D+13	0.22498D+13
$1s^2 2p 4f \ ^3G_4$	$1s^2 3s 4f \ ^3F_4$	11.32051	0.43377D+13	0.29326D+13	0.60249D+13	0.27425D+13
$1s^2 2p 4f \ ^3F_4$	$1s^2 3s 4f \ ^3F_4$	11.47708	0.44519D+13	0.29326D+13	0.60249D+13	0.28147D+13
$1s^2 2p 4f \ ^3G_5$	$1s^2 3s 4f \ ^3F_4$	11.48379	0.81945D+13	0.29326D+13	0.60249D+13	0.51809D+13
$1s^2 2s 4f \ ^3F_2$	$1s^2 3p 4f \ ^3G_3$	10.71147	0.25960D+14	0.85067D+13	0.16611D+13	0.25243D+13
$1s^2 2s 4f \ ^1F_3$	$1s^2 3p 4f \ ^1D_2$	10.61138	0.27012D+14	0.12546D+14	0.70413D+13	0.49198D+13
$1s^2 2p 4d \ ^1D_2$	$1s^2 3p 4f \ ^1D_2$	11.19218	0.93989D+13	0.12546D+14	0.70413D+13	0.17119D+13
$1s^2 2s 4f \ ^3F_3$	$1s^2 3p 4f \ ^3G_4$	10.69935	0.20518D+14	0.87660D+13	0.64416D+13	0.51221D+13
$1s^2 2s 4f \ ^3F_4$	$1s^2 3p 4f \ ^3G_4$	10.70101	0.35221D+14	0.87660D+13	0.64416D+13	0.87926D+13
$1s^2 2s 4f \ ^1F_3$	$1s^2 3p 4f \ ^3G_4$	10.70584	0.86883D+13	0.87660D+13	0.64416D+13	0.21689D+13
$1s^2 2p 4d \ ^3F_3$	$1s^2 3p 4f \ ^3G_4$	11.15826	0.82661D+13	0.87660D+13	0.64416D+13	0.20636D+13
$1s^2 2s 4f \ ^3F_3$	$1s^2 3p 4f \ ^3F_4$	10.66814	0.35642D+14	0.98078D+13	0.86367D+12	0.23719D+13
$1s^2 2s 4f \ ^3F_4$	$1s^2 3p 4f \ ^3F_4$	10.66979	0.25912D+14	0.98078D+13	0.86367D+12	0.17244D+13
$1s^2 2s 4f \ ^3F_4$	$1s^2 3p 4f \ ^3G_5$	10.66978	0.75116D+14	0.98726D+13	0.20891D+13	0.88022D+13
$1s^2 2p 4d \ ^3F_4$	$1s^2 3p 4f \ ^3G_5$	11.26266	0.24641D+14	0.98726D+13	0.20891D+13	0.28875D+13
$1s^2 2p 3d \ ^3F_3$	$1s^2 3p 4f \ ^1G_4$	8.41776	0.94891D+13	0.15189D+14	0.78331D+14	0.29828D+13
$1s^2 2p 3d \ ^1F_3$	$1s^2 3p 4f \ ^1G_4$	8.55860	0.13720D+14	0.15189D+14	0.78331D+14	0.43127D+13
$1s^2 2s 4f \ ^1F_3$	$1s^2 3p 4f \ ^1G_4$	10.64139	0.39841D+14	0.15189D+14	0.78331D+14	0.12524D+14
$1s^2 2p 4d \ ^3D_3$	$1s^2 3p 4f \ ^1G_4$	11.23461	0.10884D+14	0.15189D+14	0.78331D+14	0.34213D+13
$1s^2 2p 4d \ ^1F_3$	$1s^2 3p 4f \ ^1G_4$	11.26545	0.54224D+14	0.15189D+14	0.78331D+14	0.17045D+14
$1s^2 2p 4f \ ^1D_2$	$1s^2 3d 4f \ ^1F_3$	11.15992	0.28398D+14	0.20906D+14	0.49951D+13	0.22488D+13
$1s^2 2p 4f \ ^3G_3$	$1s^2 3d 4f \ ^3H_4$	11.09335	0.16279D+15	0.22175D+14	0.90092D+13	0.22211D+14
$1s^2 2p 4f \ ^1G_4$	$1s^2 3d 4f \ ^3H_4$	11.25778	0.20816D+14	0.22175D+14	0.90092D+13	0.28401D+13
$1s^2 2p 4f \ ^3G_4$	$1s^2 3d 4f \ ^3H_5$	11.08976	0.11057D+15	0.21873D+14	0.11504D+14	0.17060D+14
$1s^2 2p 4f \ ^3F_4$	$1s^2 3d 4f \ ^3H_5$	11.23998	0.66203D+14	0.21873D+14	0.11504D+14	0.10215D+14
$1s^2 2p 4f \ ^3G_5$	$1s^2 3d 4f \ ^3H_5$	11.24641	0.12322D+14	0.21873D+14	0.11504D+14	0.19012D+13
$1s^2 2p 4f \ ^1G_4$	$1s^2 3d 4f \ ^3H_5$	11.24878	0.51449D+14	0.21873D+14	0.11504D+14	0.79383D+13
$1s^2 2p 4f \ ^3D_3$	$1s^2 3d 4f \ ^1G_4$	11.08923	0.33060D+14	0.21798D+14	0.24385D+13	0.18835D+13
$1s^2 2p 4f \ ^3G_4$	$1s^2 3d 4f \ ^1G_4$	11.08933	0.39707D+14	0.21798D+14	0.24385D+13	0.22622D+13
$1s^2 2p 4f \ ^1F_3$	$1s^2 3d 4f \ ^1G_4$	11.23641	0.79681D+14	0.21798D+14	0.24385D+13	0.45397D+13
$1s^2 2p 4f \ ^3G_5$	$1s^2 3d 4f \ ^3H_6$	11.23742	0.27946D+15	0.21503D+14	0.11472D+14	0.43360D+14
$1s^2 2p 4f \ ^3D_3$	$1s^2 3d 4f \ ^3F_3$	11.06938	0.47979D+14	0.21678D+14	0.14065D+13	0.16805D+13
$1s^2 2p 4f \ ^1F_3$	$1s^2 3d 4f \ ^1F_3$	11.14019	0.29545D+14	0.20906D+14	0.49951D+13	0.23396D+13
$1s^2 2p 4f \ ^3D_3$	$1s^2 3d 4f \ ^1F_3$	11.14839	0.36437D+14	0.20906D+14	0.49951D+13	0.28854D+13
$1s^2 2p 4f \ ^3G_4$	$1s^2 3d 4f \ ^1H_5$	10.96525	0.64997D+14	0.22407D+14	0.45812D+14	0.11964D+14
$1s^2 2p 4f \ ^1G_4$	$1s^2 3d 4f \ ^1H_5$	11.12069	0.17681D+15	0.22407D+14	0.45812D+14	0.32545D+14

**Table V**  $\alpha$  values (in units of  $\text{cm}^3/\text{s}$ ) to the final bound states as a function of  $T_e(\text{eV})$ .

$T_e$	$1s^2 2s^2 \ ^1S_0$	$1s^2 2s 2p \ ^3P_0$	$1s^2 2s 2p \ ^1P_1$	$1s^2 2s 2p \ ^3P_1$	$1s^2 2s 2p \ ^3P_2$	$1s^2 2p^2 \ ^1S_0$
1.00	0.3560D-14	0.5563D-14	0.2345D-12	0.7293D-13	0.5199D-12	0.1247D-12
2.00	0.2912D-13	0.7130D-13	0.1505D-11	0.5510D-12	0.3216D-11	0.9205D-12
3.00	0.4762D-13	0.1395D-12	0.2217D-11	0.8932D-12	0.4663D-11	0.1425D-11
4.00	0.5598D-13	0.1862D-12	0.2452D-11	0.1062D-11	0.5104D-11	0.1615D-11
5.00	0.5916D-13	0.2200D-12	0.2488D-11	0.1149D-11	0.5138D-11	0.1661D-11
6.00	0.5989D-13	0.2465D-12	0.2448D-11	0.1199D-11	0.5021D-11	0.1646D-11
7.00	0.5945D-13	0.2684D-12	0.2381D-11	0.1231D-11	0.4852D-11	0.1608D-11
8.00	0.5844D-13	0.2867D-12	0.2305D-11	0.1252D-11	0.4672D-11	0.1559D-11
9.00	0.5715D-13	0.3018D-12	0.2230D-11	0.1265D-11	0.4495D-11	0.1509D-11
10.00	0.5572D-13	0.3139D-12	0.2156D-11	0.1272D-11	0.4327D-11	0.1458D-11
20.00	0.4186D-13	0.3353D-12	0.1594D-11	0.1166D-11	0.3109D-11	0.1059D-11
30.00	0.3189D-13	0.2908D-12	0.1227D-11	0.9670D-12	0.2364D-11	0.8048D-12
40.00	0.2505D-13	0.2435D-12	0.9724D-12	0.7939D-12	0.1863D-11	0.6334D-12
50.00	0.2024D-13	0.2044D-12	0.7916D-12	0.6590D-12	0.1511D-11	0.5132D-12
60.00	0.1676D-13	0.1734D-12	0.6591D-12	0.5554D-12	0.1255D-11	0.4258D-12
70.00	0.1416D-13	0.1490D-12	0.5590D-12	0.4750D-12	0.1063D-11	0.3603D-12
80.00	0.1216D-13	0.1295D-12	0.4815D-12	0.4116D-12	0.9144D-12	0.3098D-12
90.00	0.1058D-13	0.1138D-12	0.4201D-12	0.3608D-12	0.7971D-12	0.2699D-12
100.00	0.9319D-14	0.1010D-12	0.3706D-12	0.3194D-12	0.7027D-12	0.2378D-12
200.00	0.3812D-14	0.4257D-13	0.1527D-12	0.1335D-12	0.2883D-12	0.9743D-13
300.00	0.2189D-14	0.2460D-13	0.8772D-13	0.7692D-13	0.1652D-12	0.5581D-13
400.00	0.1463D-14	0.1647D-13	0.5857D-13	0.5143D-13	0.1102D-12	0.3720D-13
500.00	0.1066D-14	0.1200D-13	0.4262D-13	0.3745D-13	0.8008D-13	0.2704D-13
600.00	0.8216D-15	0.9241D-14	0.3280D-13	0.2882D-13	0.6157D-13	0.2079D-13
700.00	0.6582D-15	0.7397D-14	0.2624D-13	0.2306D-13	0.4923D-13	0.1663D-13
800.00	0.5427D-15	0.6094D-14	0.2161D-13	0.1900D-13	0.4053D-13	0.1369D-13
900.00	0.4575D-15	0.5134D-14	0.1820D-13	0.1600D-13	0.3411D-13	0.1152D-13
1000.00	0.3925D-15	0.4401D-14	0.1560D-13	0.1371D-13	0.2923D-13	0.9874D-14
2000.00	0.1420D-15	0.1585D-14	0.5618D-14	0.4935D-14	0.1050D-13	0.3549D-14
3000.00	0.7791D-16	0.8680D-15	0.3077D-14	0.2702D-14	0.5747D-14	0.1943D-14
4000.00	0.5082D-16	0.5655D-15	0.2005D-14	0.1760D-14	0.3743D-14	0.1265D-14
5000.00	0.3646D-16	0.4054D-15	0.1437D-14	0.1262D-14	0.2683D-14	0.9069D-15
6000.00	0.2778D-16	0.3088D-15	0.1095D-14	0.9610D-15	0.2043D-14	0.6906D-15
7000.00	0.2208D-16	0.2453D-15	0.8695D-15	0.7633D-15	0.1623D-14	0.5485D-15
8000.00	0.1808D-16	0.2009D-15	0.7122D-15	0.6251D-15	0.1329D-14	0.4492D-15
9000.00	0.1517D-16	0.1684D-15	0.5971D-15	0.5242D-15	0.1114D-14	0.3766D-15
10000.00	0.1296D-16	0.1439D-15	0.5101D-15	0.4477D-15	0.9515D-15	0.3217D-15

Table V (continued)

$T_e$	$1s^2 2p^2 \ ^3P_0$	$1s^2 2p^2 \ ^3P_1$	$1s^2 2p^2 \ ^3P_2$	$1s^2 2p^2 \ ^1D_2$	$1s^2 2s 3s \ ^1S_0$	$1s^2 2s 3s \ ^3S_1$
1.00	0.4066D-13	0.2258D-12	0.3966D-12	0.6625D-12	0.6905D-15	0.1392D-14
2.00	0.4513D-12	0.1609D-11	0.3042D-11	0.4796D-11	0.5458D-14	0.9939D-14
3.00	0.8430D-12	0.2468D-11	0.4882D-11	0.7360D-11	0.8821D-14	0.1526D-13
4.00	0.1073D-11	0.2801D-11	0.5690D-11	0.8295D-11	0.1030D-13	0.1728D-13
5.00	0.1195D-11	0.2909D-11	0.5981D-11	0.8498D-11	0.1084D-13	0.1784D-13
6.00	0.1257D-11	0.2929D-11	0.6031D-11	0.8400D-11	0.1094D-13	0.1782D-13
7.00	0.1283D-11	0.2913D-11	0.5966D-11	0.8182D-11	0.1083D-13	0.1755D-13
8.00	0.1289D-11	0.2882D-11	0.5844D-11	0.7921D-11	0.1062D-13	0.1719D-13
9.00	0.1281D-11	0.2845D-11	0.5694D-11	0.7649D-11	0.1036D-13	0.1680D-13
10.00	0.1266D-11	0.2804D-11	0.5530D-11	0.7382D-11	0.1009D-13	0.1639D-13
20.00	0.9917D-12	0.2319D-11	0.3996D-11	0.5318D-11	0.7514D-14	0.1270D-13
30.00	0.7544D-12	0.1861D-11	0.2961D-11	0.4025D-11	0.5797D-14	0.1046D-13
40.00	0.5890D-12	0.1506D-11	0.2285D-11	0.3161D-11	0.5311D-14	0.1252D-13
50.00	0.4734D-12	0.1240D-11	0.1825D-11	0.2558D-11	0.6539D-14	0.2171D-13
60.00	0.3901D-12	0.1040D-11	0.1498D-11	0.2121D-11	0.9411D-14	0.3759D-13
70.00	0.3282D-12	0.8860D-12	0.1257D-11	0.1793D-11	0.1340D-13	0.5760D-13
80.00	0.2808D-12	0.7659D-12	0.1073D-11	0.1541D-11	0.1792D-13	0.7907D-13
90.00	0.2437D-12	0.6700D-12	0.9301D-12	0.1342D-11	0.2251D-13	0.1001D-12
100.00	0.2140D-12	0.5923D-12	0.8161D-12	0.1182D-11	0.2688D-13	0.1195D-12
200.00	0.8618D-13	0.2460D-12	0.3272D-12	0.4836D-12	0.4940D-13	0.2078D-12
300.00	0.4904D-13	0.1415D-12	0.1860D-12	0.2768D-12	0.5095D-13	0.2061D-12
400.00	0.3258D-13	0.9450D-13	0.1235D-12	0.1845D-12	0.4707D-13	0.1857D-12
500.00	0.2363D-13	0.6877D-13	0.8952D-13	0.1341D-12	0.4218D-13	0.1638D-12
600.00	0.1814D-13	0.5291D-13	0.6871D-13	0.1030D-12	0.3752D-13	0.1440D-12
700.00	0.1449D-13	0.4232D-13	0.5487D-13	0.8238D-13	0.3341D-13	0.1272D-12
800.00	0.1192D-13	0.3485D-13	0.4513D-13	0.6780D-13	0.2987D-13	0.1130D-12
900.00	0.1002D-13	0.2935D-13	0.3796D-13	0.5707D-13	0.2684D-13	0.1010D-12
1000.00	0.8585D-14	0.2515D-13	0.3251D-13	0.4890D-13	0.2425D-13	0.9093D-13
2000.00	0.3077D-14	0.9045D-14	0.1165D-13	0.1756D-13	0.1116D-13	0.4108D-13
3000.00	0.1683D-14	0.4952D-14	0.6370D-14	0.9612D-14	0.6652D-14	0.2434D-13
4000.00	0.1096D-14	0.3226D-14	0.4147D-14	0.6260D-14	0.4525D-14	0.1651D-13
5000.00	0.7850D-15	0.2312D-14	0.2971D-14	0.4486D-14	0.3329D-14	0.1212D-13
6000.00	0.5977D-15	0.1761D-14	0.2262D-14	0.3417D-14	0.2580D-14	0.9384D-14
7000.00	0.4746D-15	0.1398D-14	0.1797D-14	0.2713D-14	0.2075D-14	0.7541D-14
8000.00	0.3887D-15	0.1145D-14	0.1471D-14	0.2222D-14	0.1715D-14	0.6230D-14
9000.00	0.3259D-15	0.9602D-15	0.1233D-14	0.1863D-14	0.1449D-14	0.5260D-14
10000.00	0.2783D-15	0.8202D-15	0.1053D-14	0.1591D-14	0.1245D-14	0.4517D-14

Table V (continued)

$T_e$	$1s^2 2s 3p^3 P_0$	$1s^2 2s 3p^1 P_1$	$1s^2 2s 3p^3 P_1$	$1s^2 2s 3p^3 P_2$	$1s^2 2s 3d^3 D_1$	$1s^2 2s 3d^1 D_2$
1.00	0.6749D-16	0.3816D-14	0.2446D-14	0.4027D-15	0.5752D-15	0.2487D-14
2.00	0.5352D-15	0.2605D-13	0.1584D-13	0.2687D-14	0.5522D-14	0.1968D-13
3.00	0.8551D-15	0.3985D-13	0.2347D-13	0.4014D-14	0.9729D-14	0.3179D-13
4.00	0.9931D-15	0.4538D-13	0.2607D-13	0.4485D-14	0.1198D-13	0.3713D-13
5.00	0.1046D-14	0.4727D-13	0.2656D-13	0.4607D-14	0.1305D-13	0.3906D-13
6.00	0.1063D-14	0.4768D-13	0.2625D-13	0.4598D-14	0.1351D-13	0.3942D-13
7.00	0.1063D-14	0.4745D-13	0.2563D-13	0.4544D-14	0.1363D-13	0.3902D-13
8.00	0.1055D-14	0.4694D-13	0.2492D-13	0.4473D-14	0.1356D-13	0.3827D-13
9.00	0.1044D-14	0.4630D-13	0.2420D-13	0.4398D-14	0.1337D-13	0.3735D-13
10.00	0.1030D-14	0.4559D-13	0.2349D-13	0.4322D-14	0.1312D-13	0.3636D-13
20.00	0.8535D-15	0.3740D-13	0.1778D-13	0.3557D-14	0.9948D-14	0.2701D-13
30.00	0.6846D-15	0.3015D-13	0.1396D-13	0.2896D-14	0.7499D-14	0.2048D-13
40.00	0.5563D-15	0.2651D-13	0.1239D-13	0.2638D-14	0.5901D-14	0.1637D-13
50.00	0.4668D-15	0.2830D-13	0.1389D-13	0.3080D-14	0.4955D-14	0.1434D-13
60.00	0.4086D-15	0.3541D-13	0.1832D-13	0.4266D-14	0.4490D-14	0.1406D-13
70.00	0.3759D-15	0.4629D-13	0.2478D-13	0.6010D-14	0.4355D-14	0.1509D-13
80.00	0.3656D-15	0.5908D-13	0.3223D-13	0.8065D-14	0.4425D-14	0.1693D-13
90.00	0.3764D-15	0.7228D-13	0.3985D-13	0.1023D-13	0.4611D-14	0.1918D-13
100.00	0.4074D-15	0.8491D-13	0.4710D-13	0.1235D-13	0.4850D-14	0.2157D-13
200.00	0.1326D-14	0.1440D-12	0.8121D-13	0.2510D-13	0.6449D-14	0.3571D-13
300.00	0.2090D-14	0.1389D-12	0.7907D-13	0.2721D-13	0.6333D-14	0.3621D-13
400.00	0.2375D-14	0.1219D-12	0.6988D-13	0.2577D-13	0.5726D-14	0.3299D-13
500.00	0.2392D-14	0.1051D-12	0.6059D-13	0.2340D-13	0.5064D-14	0.2925D-13
600.00	0.2290D-14	0.9087D-13	0.5258D-13	0.2098D-13	0.4464D-14	0.2580D-13
700.00	0.2143D-14	0.7914D-13	0.4593D-13	0.1877D-13	0.3949D-14	0.2282D-13
800.00	0.1987D-14	0.6952D-13	0.4044D-13	0.1684D-13	0.3512D-14	0.2030D-13
900.00	0.1835D-14	0.6160D-13	0.3590D-13	0.1517D-13	0.3143D-14	0.1817D-13
1000.00	0.1693D-14	0.5501D-13	0.3211D-13	0.1372D-13	0.2831D-14	0.1636D-13
2000.00	0.8508D-15	0.2392D-13	0.1407D-13	0.6349D-14	0.1283D-14	0.7399D-14
3000.00	0.5210D-15	0.1398D-13	0.8245D-14	0.3788D-14	0.7606D-15	0.4385D-14
4000.00	0.3590D-15	0.9412D-14	0.5559D-14	0.2578D-14	0.5159D-15	0.2973D-14
5000.00	0.2661D-15	0.6882D-14	0.4068D-14	0.1897D-14	0.3790D-15	0.2183D-14
6000.00	0.2073D-15	0.5312D-14	0.3141D-14	0.1470D-14	0.2934D-15	0.1690D-14
7000.00	0.1673D-15	0.4259D-14	0.2520D-14	0.1182D-14	0.2358D-15	0.1358D-14
8000.00	0.1387D-15	0.3513D-14	0.2079D-14	0.9775D-15	0.1948D-15	0.1122D-14
9000.00	0.1174D-15	0.2962D-14	0.1754D-14	0.8257D-15	0.1645D-15	0.9472D-15
10000.00	0.1010D-15	0.2542D-14	0.1505D-14	0.7094D-15	0.1413D-15	0.8134D-15

Table V (continued)

$T_e$	$1s^2 2s 3d \ ^3D_2$	$1s^2 2s 3d \ ^3D_3$	$1s^2 2p 3s \ ^3P_0$	$1s^2 2p 3s \ ^1P_1$	$1s^2 2p 3s \ ^3P_1$	$1s^2 2p 3s \ ^3P_2$
1.00	0.1367D-14	0.2069D-14	0.1733D-14	0.8329D-13	0.9075D-14	0.1571D-12
2.00	0.1074D-13	0.1401D-13	0.2201D-13	0.5277D-12	0.8831D-13	0.9720D-12
3.00	0.1739D-13	0.2098D-13	0.4293D-13	0.7718D-12	0.1593D-12	0.1410D-11
4.00	0.2037D-13	0.2333D-13	0.5721D-13	0.8483D-12	0.2029D-12	0.1543D-11
5.00	0.2148D-13	0.2370D-13	0.6748D-13	0.8560D-12	0.2313D-12	0.1554D-11
6.00	0.2173D-13	0.2328D-13	0.7554D-13	0.8378D-12	0.2520D-12	0.1518D-11
7.00	0.2155D-13	0.2257D-13	0.8218D-13	0.8106D-12	0.2680D-12	0.1468D-11
8.00	0.2117D-13	0.2176D-13	0.8771D-13	0.7811D-12	0.2807D-12	0.1413D-11
9.00	0.2068D-13	0.2094D-13	0.9225D-13	0.7519D-12	0.2909D-12	0.1360D-11
10.00	0.2014D-13	0.2013D-13	0.9593D-13	0.7241D-12	0.2987D-12	0.1309D-11
20.00	0.1493D-13	0.1396D-13	0.1022D-12	0.5211D-12	0.3008D-12	0.9413D-12
30.00	0.1126D-13	0.1029D-13	0.8874D-13	0.3969D-12	0.2566D-12	0.7165D-12
40.00	0.8940D-14	0.8077D-14	0.7550D-13	0.3164D-12	0.2167D-12	0.5706D-12
50.00	0.7699D-14	0.6917D-14	0.6702D-13	0.2672D-12	0.1919D-12	0.4806D-12
60.00	0.7307D-14	0.6549D-14	0.6346D-13	0.2409D-12	0.1820D-12	0.4313D-12
70.00	0.7511D-14	0.6729D-14	0.6360D-13	0.2302D-12	0.1832D-12	0.4095D-12
80.00	0.8079D-14	0.7247D-14	0.6603D-13	0.2293D-12	0.1911D-12	0.4049D-12
90.00	0.8835D-14	0.7946D-14	0.6961D-13	0.2339D-12	0.2025D-12	0.4100D-12
100.00	0.9658D-14	0.8722D-14	0.7358D-13	0.2411D-12	0.2151D-12	0.4198D-12
200.00	0.1478D-13	0.1440D-13	0.9329D-13	0.2898D-12	0.2797D-12	0.4814D-12
300.00	0.1509D-13	0.1570D-13	0.8720D-13	0.2759D-12	0.2645D-12	0.4450D-12
400.00	0.1393D-13	0.1512D-13	0.7631D-13	0.2460D-12	0.2332D-12	0.3889D-12
500.00	0.1250D-13	0.1393D-13	0.6602D-13	0.2159D-12	0.2027D-12	0.3367D-12
600.00	0.1113D-13	0.1264D-13	0.5730D-13	0.1895D-12	0.1766D-12	0.2925D-12
700.00	0.9914D-14	0.1141D-13	0.5010D-13	0.1671D-12	0.1548D-12	0.2559D-12
800.00	0.8868D-14	0.1031D-13	0.4417D-13	0.1483D-12	0.1368D-12	0.2258D-12
900.00	0.7974D-14	0.9342D-14	0.3926D-13	0.1326D-12	0.1218D-12	0.2008D-12
1000.00	0.7207D-14	0.8497D-14	0.3516D-13	0.1193D-12	0.1092D-12	0.1800D-12
2000.00	0.3322D-14	0.4024D-14	0.1552D-13	0.5380D-13	0.4850D-13	0.7970D-13
3000.00	0.1982D-14	0.2421D-14	0.9126D-14	0.3187D-13	0.2857D-13	0.4691D-13
4000.00	0.1348D-14	0.1654D-14	0.6163D-14	0.2160D-13	0.1932D-13	0.3170D-13
5000.00	0.9924D-15	0.1221D-14	0.4516D-14	0.1586D-13	0.1416D-13	0.2323D-13
6000.00	0.7693D-15	0.9477D-15	0.3490D-14	0.1228D-13	0.1095D-13	0.1796D-13
7000.00	0.6187D-15	0.7632D-15	0.2801D-14	0.9867D-14	0.8791D-14	0.1442D-13
8000.00	0.5116D-15	0.6315D-15	0.2312D-14	0.8152D-14	0.7259D-14	0.1190D-13
9000.00	0.4321D-15	0.5338D-15	0.1951D-14	0.6882D-14	0.6125D-14	0.1004D-13
10000.00	0.3713D-15	0.4589D-15	0.1675D-14	0.5910D-14	0.5258D-14	0.8621D-14

Table V (continued)

$T_e$	$1s^2 2p 3p \ ^1S_0$	$1s^2 2p 3p \ ^3P_0$	$1s^2 2p 3p \ ^3S_1$	$1s^2 2p 3p \ ^3P_1$	$1s^2 2p 3p \ ^3D_1$	$1s^2 2p 3p \ ^3P_2$
1.00	0.2114D-13	0.5884D-14	0.9870D-13	0.5826D-13	0.1159D-13	0.9713D-13
2.00	0.1518D-12	0.6928D-13	0.6998D-12	0.3982D-12	0.1600D-12	0.6653D-12
3.00	0.2325D-12	0.1321D-12	0.1071D-11	0.5989D-12	0.3175D-12	0.1005D-11
4.00	0.2617D-12	0.1698D-12	0.1212D-11	0.6679D-12	0.4159D-12	0.1125D-11
5.00	0.2680D-12	0.1905D-12	0.1254D-11	0.6800D-12	0.4710D-12	0.1148D-11
6.00	0.2649D-12	0.2012D-12	0.1256D-11	0.6696D-12	0.5000D-12	0.1131D-11
7.00	0.2580D-12	0.2061D-12	0.1243D-11	0.6506D-12	0.5133D-12	0.1098D-11
8.00	0.2498D-12	0.2075D-12	0.1223D-11	0.6289D-12	0.5170D-12	0.1059D-11
9.00	0.2414D-12	0.2067D-12	0.1201D-11	0.6068D-12	0.5146D-12	0.1019D-11
10.00	0.2331D-12	0.2045D-12	0.1177D-11	0.5853D-12	0.5083D-12	0.9788D-12
20.00	0.1690D-12	0.1609D-12	0.9461D-12	0.4212D-12	0.3910D-12	0.6703D-12
30.00	0.1285D-12	0.1224D-12	0.7515D-12	0.3186D-12	0.2926D-12	0.4885D-12
40.00	0.1013D-12	0.9550D-13	0.6055D-12	0.2503D-12	0.2263D-12	0.3745D-12
50.00	0.8220D-13	0.7678D-13	0.4990D-12	0.2032D-12	0.1813D-12	0.2992D-12
60.00	0.6833D-13	0.6337D-13	0.4207D-12	0.1698D-12	0.1501D-12	0.2475D-12
70.00	0.5795D-13	0.5348D-13	0.3623D-12	0.1455D-12	0.1278D-12	0.2108D-12
80.00	0.4997D-13	0.4598D-13	0.3177D-12	0.1274D-12	0.1113D-12	0.1842D-12
90.00	0.4369D-13	0.4016D-13	0.2831D-12	0.1135D-12	0.9888D-13	0.1643D-12
100.00	0.3865D-13	0.3553D-13	0.2555D-12	0.1026D-12	0.8924D-13	0.1491D-12
200.00	0.1664D-13	0.1595D-13	0.1337D-12	0.5585D-13	0.4981D-13	0.8848D-13
300.00	0.1003D-13	0.1028D-13	0.9147D-13	0.3957D-13	0.3719D-13	0.6738D-13
400.00	0.6982D-14	0.7615D-14	0.6880D-13	0.3049D-13	0.3008D-13	0.5443D-13
500.00	0.5262D-14	0.6041D-14	0.5448D-13	0.2455D-13	0.2518D-13	0.4526D-13
600.00	0.4168D-14	0.4984D-14	0.4462D-13	0.2035D-13	0.2152D-13	0.3840D-13
700.00	0.3415D-14	0.4221D-14	0.3745D-13	0.1724D-13	0.1868D-13	0.3309D-13
800.00	0.2870D-14	0.3643D-14	0.3204D-13	0.1486D-13	0.1641D-13	0.2889D-13
900.00	0.2458D-14	0.3189D-14	0.2782D-13	0.1297D-13	0.1456D-13	0.2550D-13
1000.00	0.2137D-14	0.2825D-14	0.2446D-13	0.1146D-13	0.1303D-13	0.2271D-13
2000.00	0.8300D-15	0.1202D-14	0.9955D-14	0.4766D-14	0.5774D-14	0.9832D-14
3000.00	0.4684D-15	0.7007D-15	0.5700D-14	0.2749D-14	0.3408D-14	0.5750D-14
4000.00	0.3100D-15	0.4717D-15	0.3800D-14	0.1839D-14	0.2307D-14	0.3874D-14
5000.00	0.2245D-15	0.3451D-15	0.2762D-14	0.1340D-14	0.1693D-14	0.2835D-14
6000.00	0.1721D-15	0.2664D-15	0.2124D-14	0.1032D-14	0.1310D-14	0.2189D-14
7000.00	0.1374D-15	0.2137D-15	0.1698D-14	0.8262D-15	0.1052D-14	0.1756D-14
8000.00	0.1129D-15	0.1763D-15	0.1398D-14	0.6807D-15	0.8692D-15	0.1449D-14
9000.00	0.9497D-16	0.1487D-15	0.1177D-14	0.5734D-15	0.7337D-15	0.1222D-14
10000.00	0.8130D-16	0.1276D-15	0.1008D-14	0.4916D-15	0.6300D-15	0.1048D-14



Table V (continued)

$T_e$	$1s^2 2p3p \ ^1D_2$	$1s^2 2p3p \ ^3D_2$	$1s^2 2p3p \ ^3D_3$	$1s^2 2p3d \ ^3P_0$	$1s^2 2p3d \ ^1P_1$	$1s^2 2p3d \ ^3P_1$
1.00	0.1152D-12	0.3322D-13	0.1632D-12	0.1573D-13	0.3401D-13	0.3998D-13
2.00	0.8146D-12	0.3483D-12	0.1089D-11	0.1177D-12	0.2574D-12	0.2970D-12
3.00	0.1239D-11	0.6409D-12	0.1620D-11	0.1818D-12	0.3992D-12	0.4578D-12
4.00	0.1389D-11	0.8101D-12	0.1794D-11	0.2052D-12	0.4516D-12	0.5161D-12
5.00	0.1417D-11	0.8996D-12	0.1816D-11	0.2102D-12	0.4634D-12	0.5285D-12
6.00	0.1398D-11	0.9441D-12	0.1777D-11	0.2077D-12	0.4585D-12	0.5219D-12
7.00	0.1358D-11	0.9631D-12	0.1716D-11	0.2021D-12	0.4469D-12	0.5078D-12
8.00	0.1313D-11	0.9666D-12	0.1647D-11	0.1955D-12	0.4329D-12	0.4912D-12
9.00	0.1266D-11	0.9607D-12	0.1578D-11	0.1886D-12	0.4183D-12	0.4739D-12
10.00	0.1221D-11	0.9487D-12	0.1511D-11	0.1819D-12	0.4039D-12	0.4570D-12
20.00	0.8775D-12	0.7415D-12	0.1012D-11	0.1305D-12	0.2920D-12	0.3278D-12
30.00	0.6646D-12	0.5634D-12	0.7288D-12	0.9859D-13	0.2213D-12	0.2476D-12
40.00	0.5233D-12	0.4400D-12	0.5549D-12	0.7736D-13	0.1742D-12	0.1942D-12
50.00	0.4269D-12	0.3549D-12	0.4408D-12	0.6259D-13	0.1423D-12	0.1572D-12
60.00	0.3605D-12	0.2953D-12	0.3621D-12	0.5193D-13	0.1207D-12	0.1305D-12
70.00	0.3146D-12	0.2525D-12	0.3058D-12	0.4397D-13	0.1063D-12	0.1107D-12
80.00	0.2825D-12	0.2210D-12	0.2641D-12	0.3787D-13	0.9677D-13	0.9557D-13
90.00	0.2598D-12	0.1974D-12	0.2325D-12	0.3308D-13	0.9055D-13	0.8373D-13
100.00	0.2433D-12	0.1791D-12	0.2079D-12	0.2924D-13	0.8651D-13	0.7429D-13
200.00	0.1810D-12	0.1060D-12	0.1108D-12	0.1257D-13	0.7575D-13	0.3355D-13
300.00	0.1483D-12	0.8122D-13	0.8454D-13	0.7597D-14	0.6664D-13	0.2131D-13
400.00	0.1224D-12	0.6608D-13	0.7052D-13	0.5307D-14	0.5702D-13	0.1550D-13
500.00	0.1023D-12	0.5527D-13	0.6056D-13	0.4010D-14	0.4872D-13	0.1208D-13
600.00	0.8677D-13	0.4709D-13	0.5278D-13	0.3182D-14	0.4193D-13	0.9822D-14
700.00	0.7462D-13	0.4071D-13	0.4649D-13	0.2612D-14	0.3643D-13	0.8217D-14
800.00	0.6497D-13	0.3563D-13	0.4131D-13	0.2197D-14	0.3197D-13	0.7020D-14
900.00	0.5717D-13	0.3151D-13	0.3698D-13	0.1884D-14	0.2831D-13	0.6095D-14
1000.00	0.5079D-13	0.2812D-13	0.3334D-13	0.1639D-14	0.2527D-13	0.5360D-14
2000.00	0.2161D-13	0.1226D-13	0.1529D-13	0.6389D-15	0.1099D-13	0.2199D-14
3000.00	0.1254D-13	0.7184D-14	0.9122D-14	0.3609D-15	0.6427D-14	0.1265D-14
4000.00	0.8417D-14	0.4846D-14	0.6210D-14	0.2390D-15	0.4328D-14	0.8462D-15
5000.00	0.6142D-14	0.3548D-14	0.4572D-14	0.1731D-15	0.3166D-14	0.6164D-15
6000.00	0.4734D-14	0.2741D-14	0.3545D-14	0.1327D-15	0.2444D-14	0.4746D-15
7000.00	0.3793D-14	0.2199D-14	0.2852D-14	0.1060D-15	0.1960D-14	0.3799D-15
8000.00	0.3126D-14	0.1815D-14	0.2359D-14	0.8712D-16	0.1617D-14	0.3130D-15
9000.00	0.2635D-14	0.1531D-14	0.1993D-14	0.7327D-16	0.1363D-14	0.2636D-15
10000.00	0.2260D-14	0.1314D-14	0.1712D-14	0.6273D-16	0.1170D-14	0.2260D-15

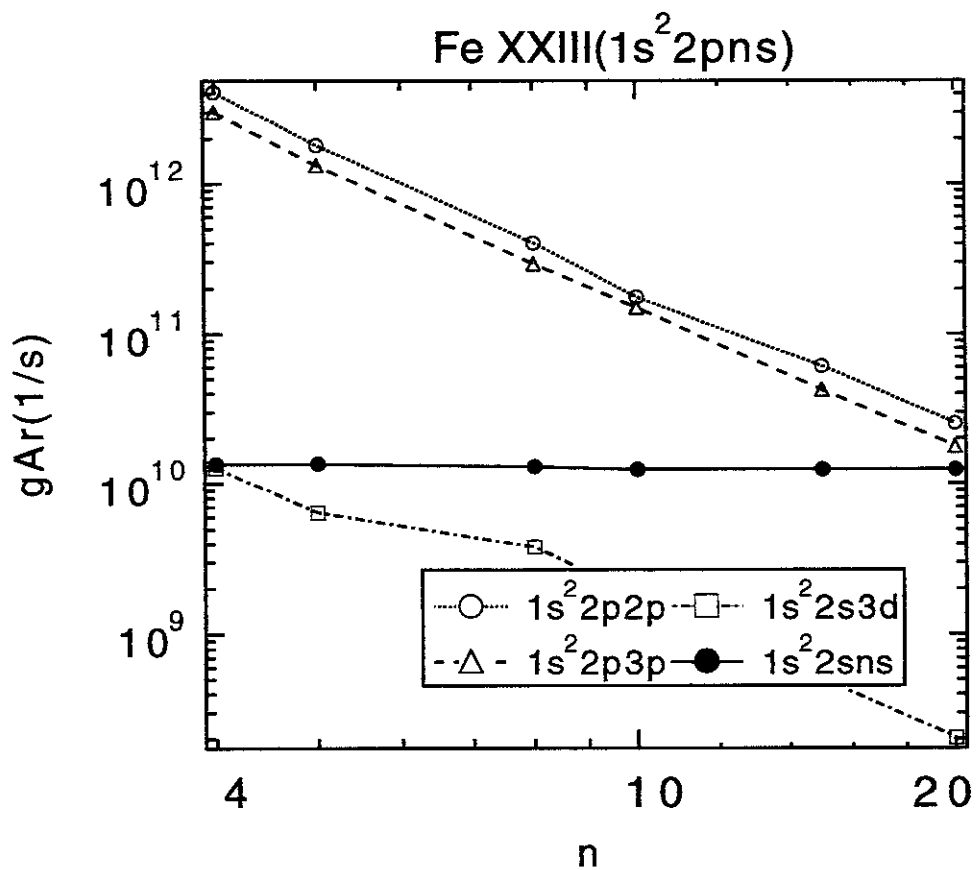


Fig.1(a)

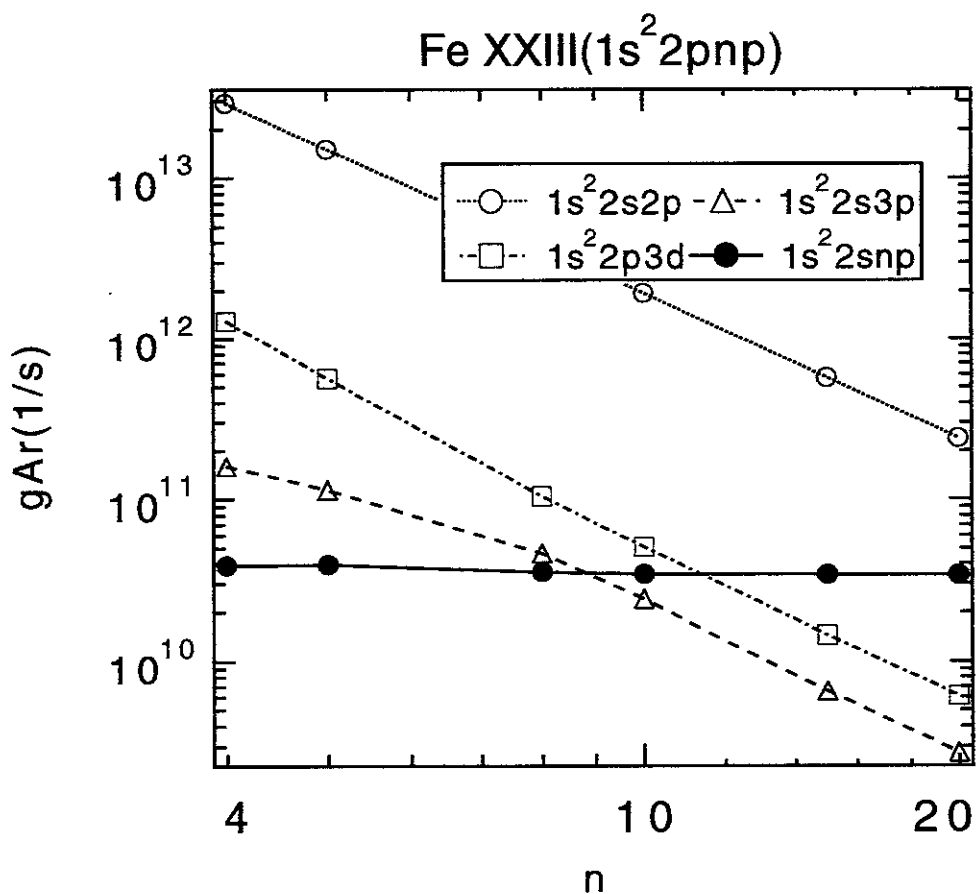


Fig.1(b)

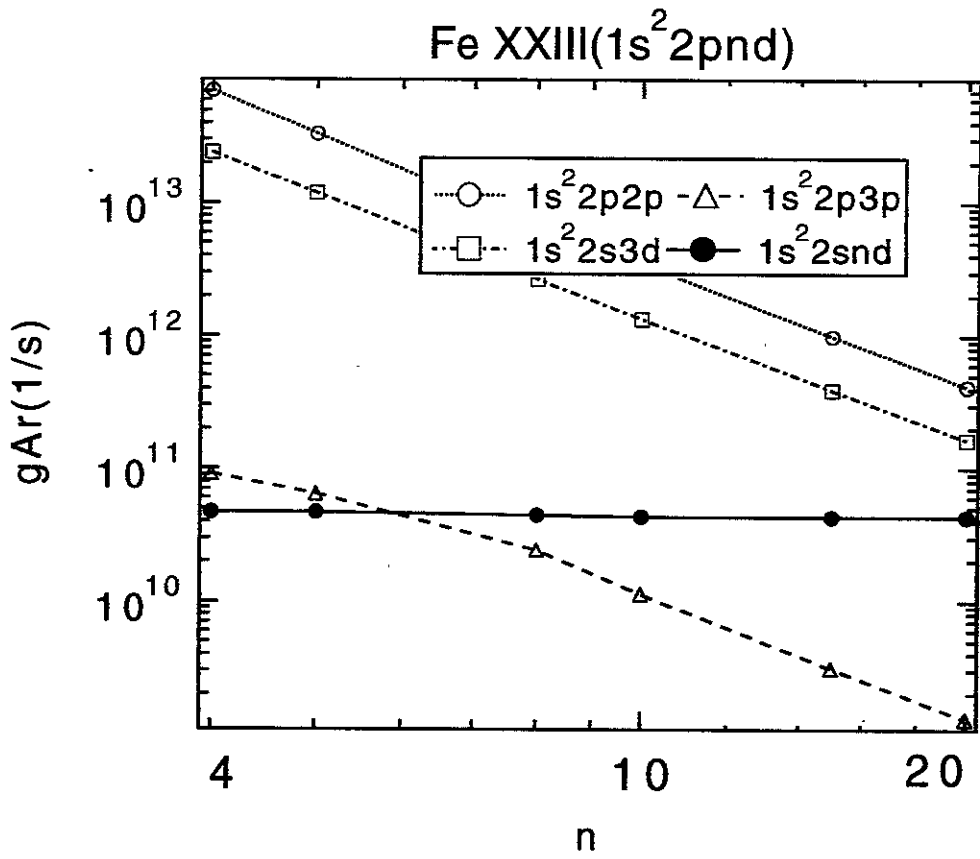


Fig.1(c)

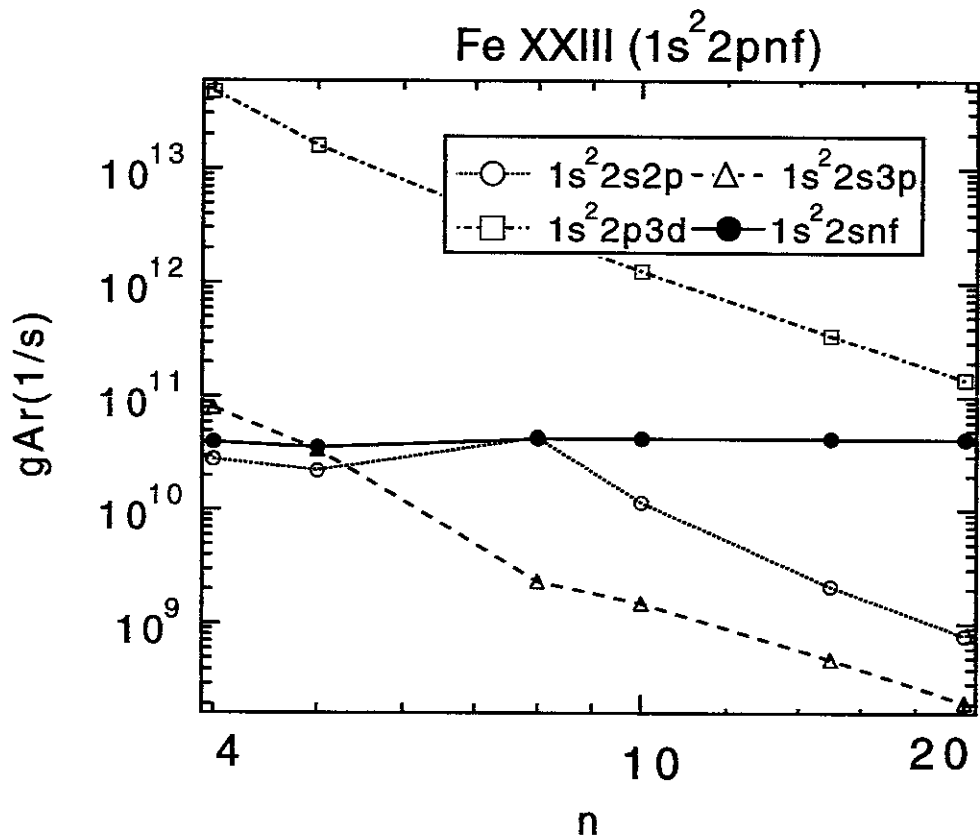


Fig.1(d)

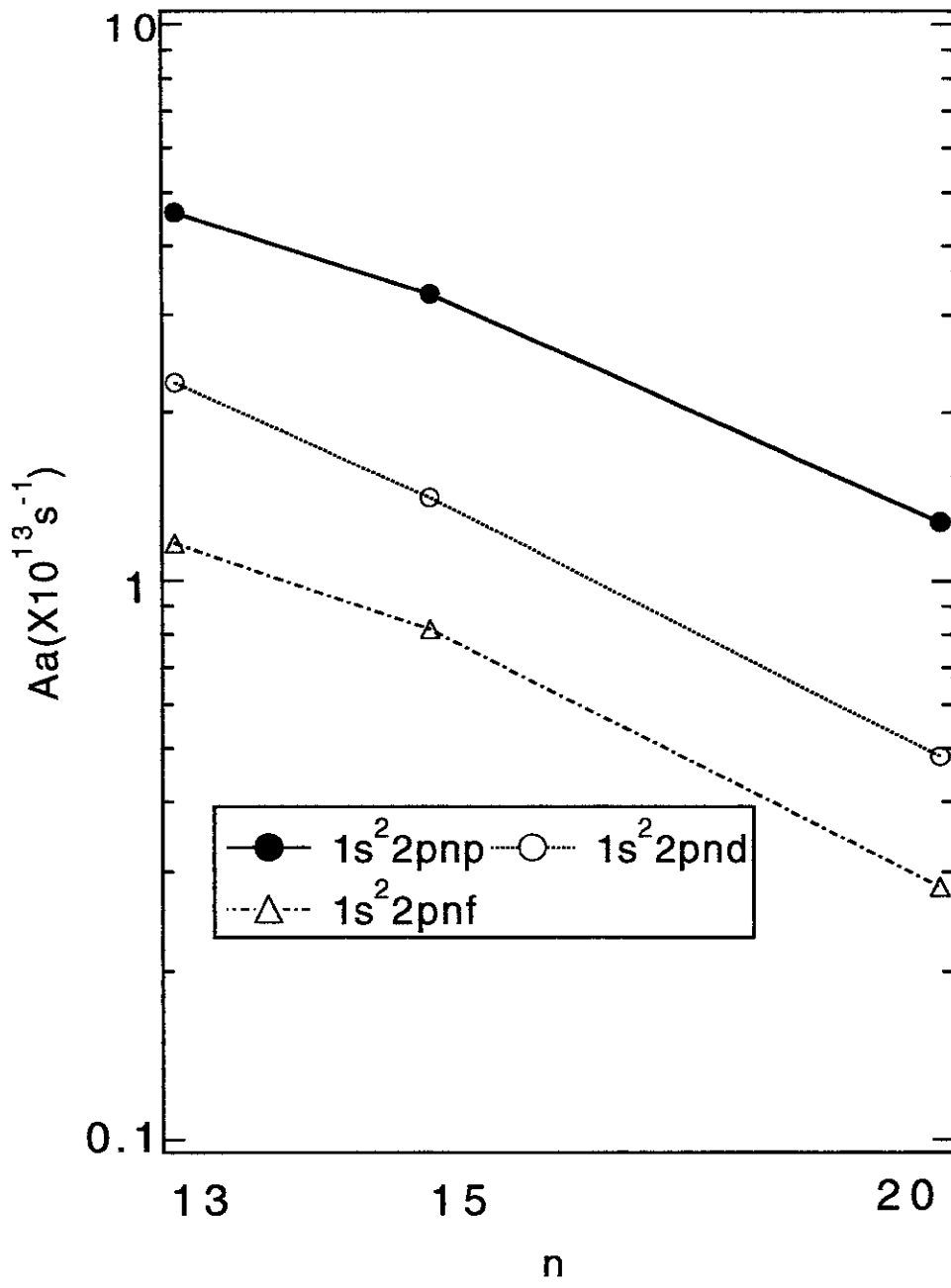


Fig.2

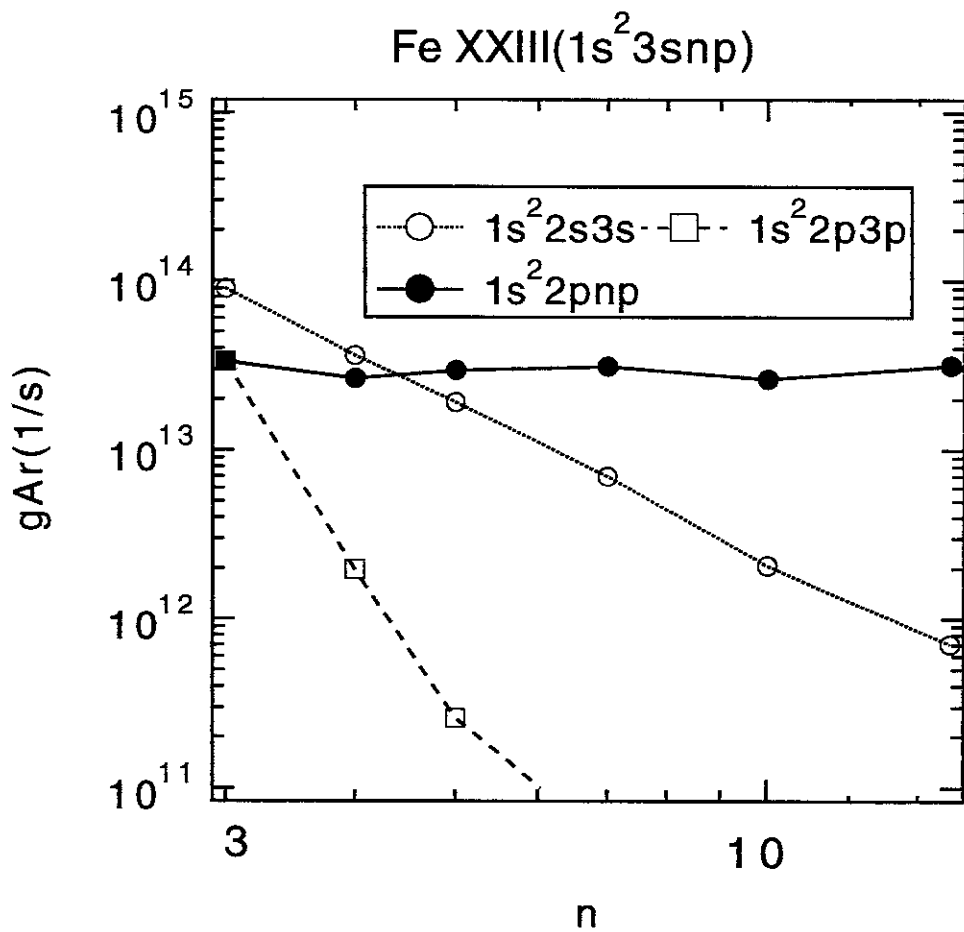


Fig.3(a)

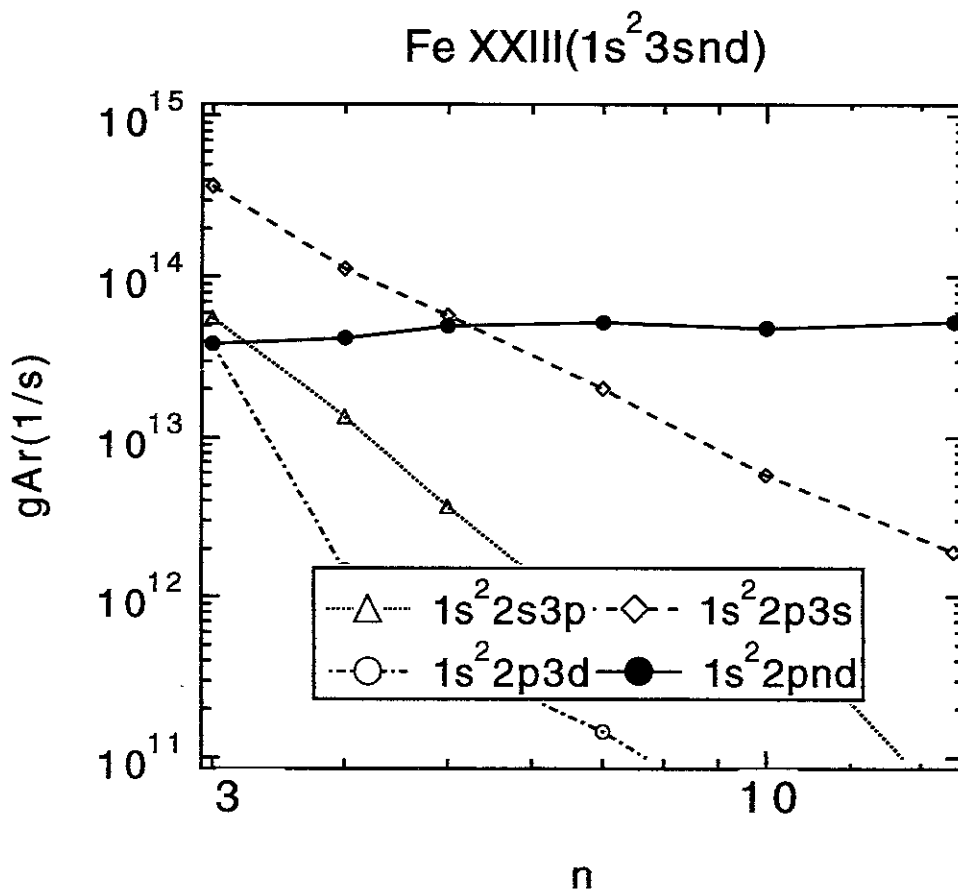


Fig.3(b)

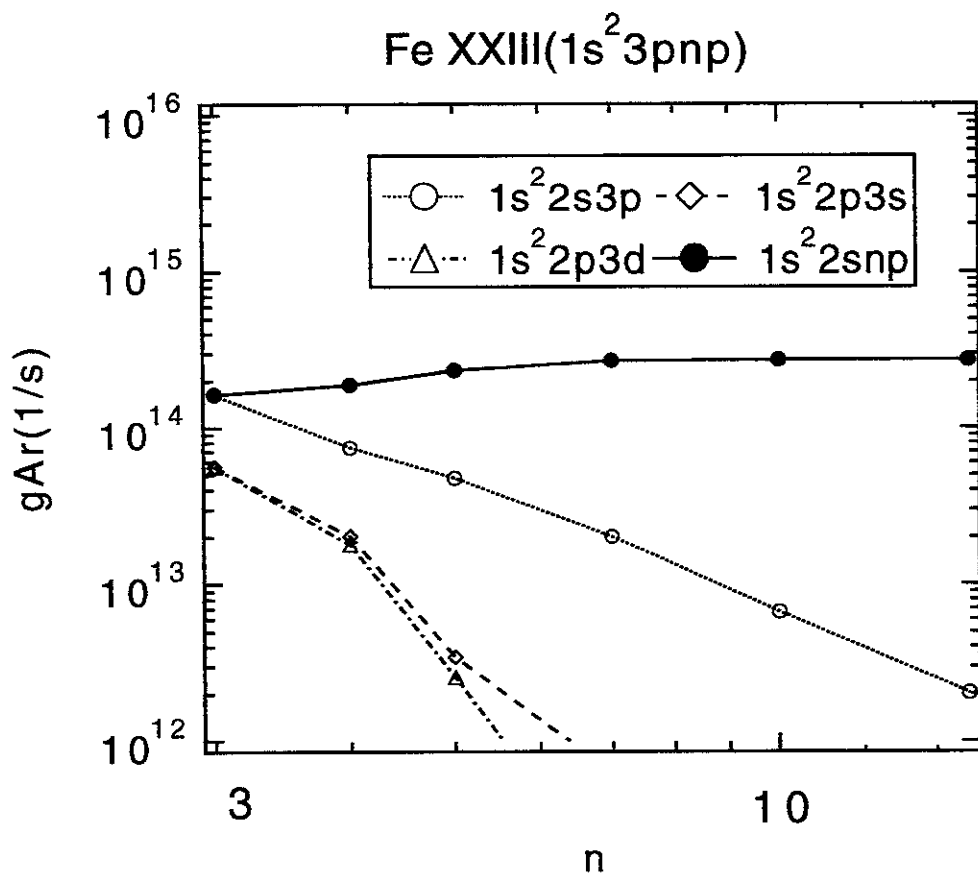


Fig.3(c)

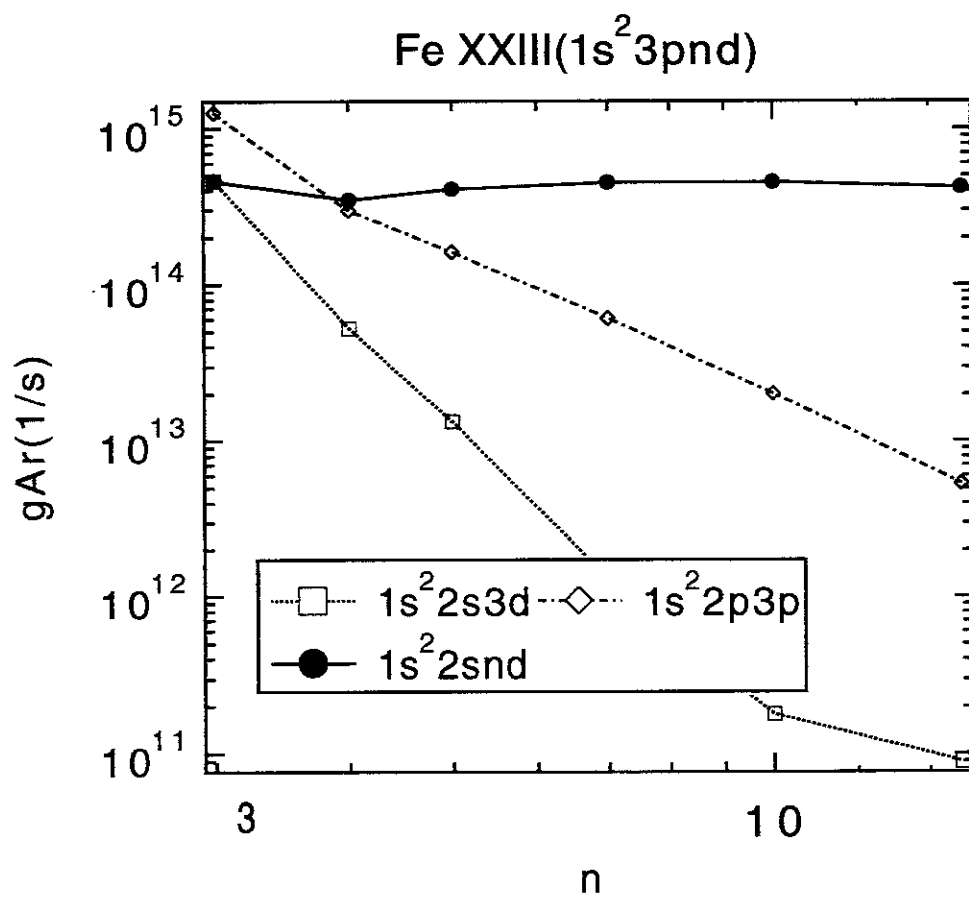


Fig.3(d)

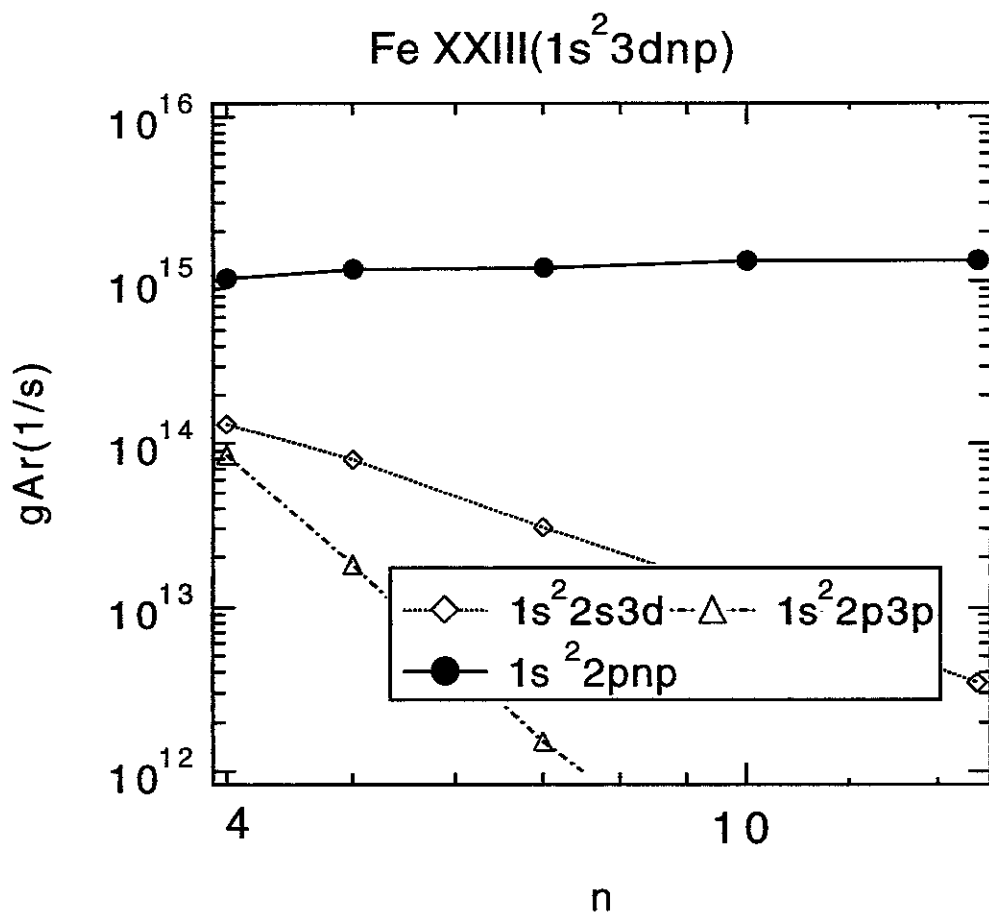


Fig.3(e)

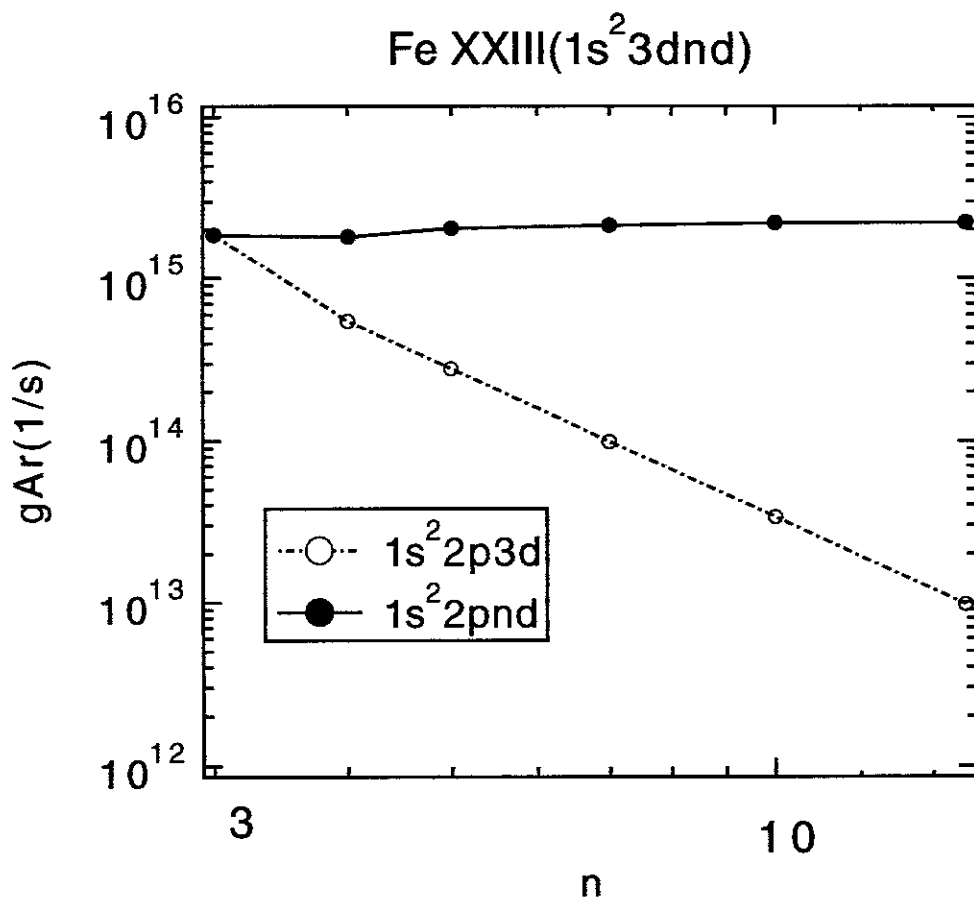


Fig.3(f)

Fe XXIII(1s<sup>2</sup>3snp)

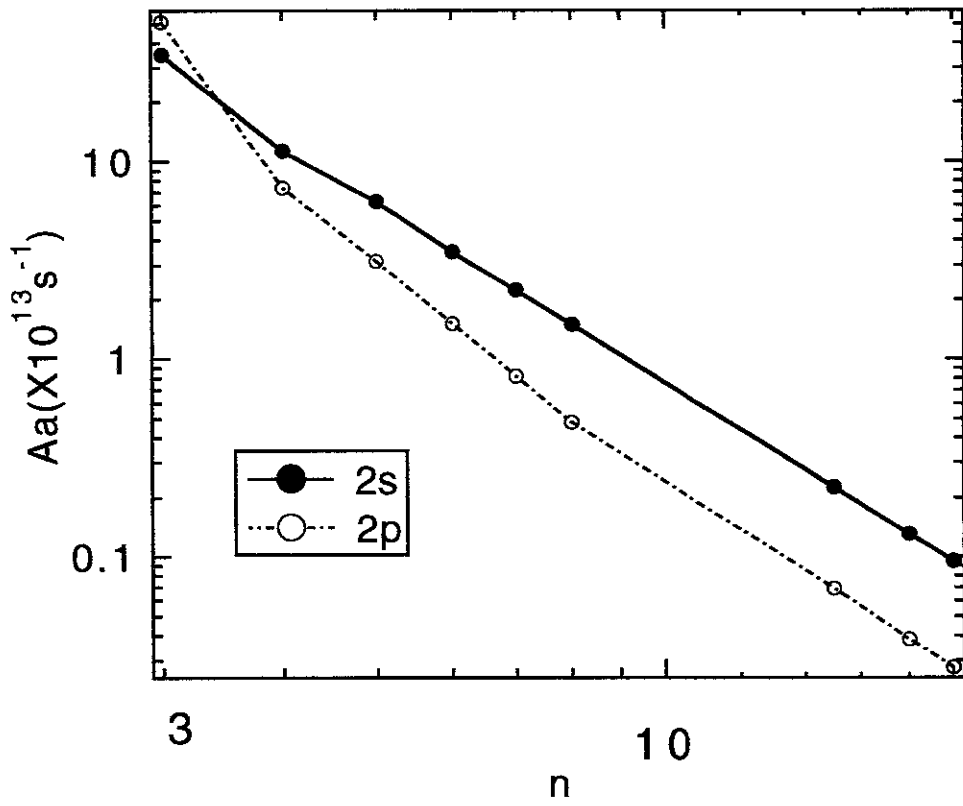


Fig.4(a)

Fe XXIII(1s<sup>2</sup>3snd)

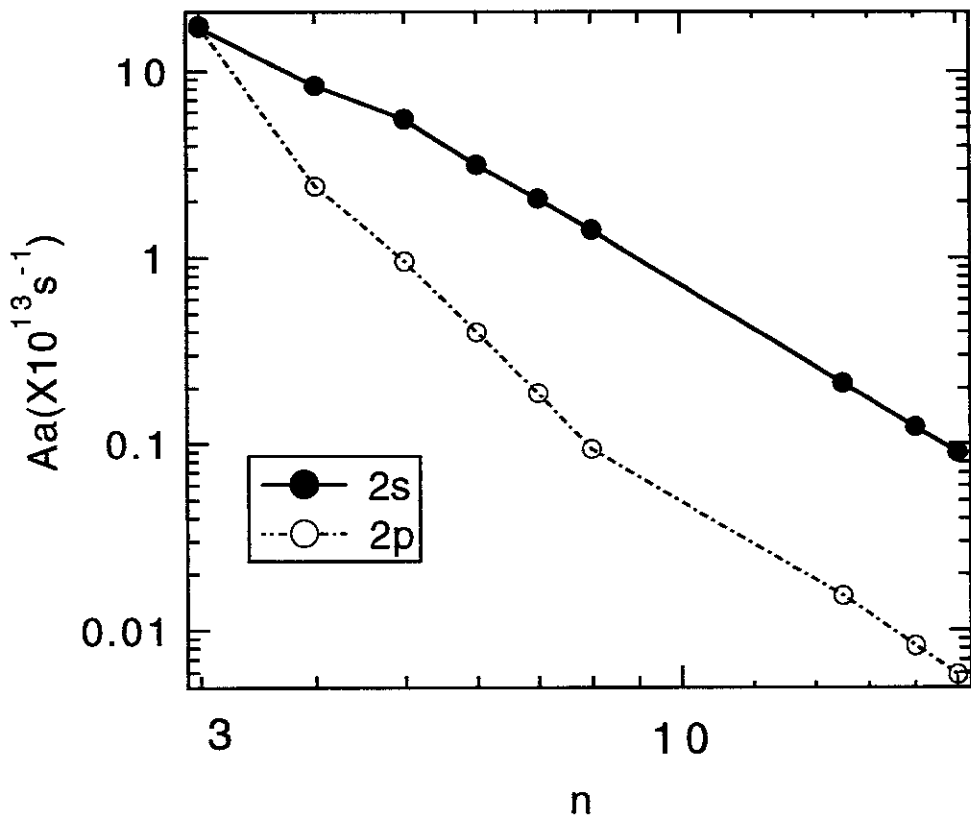


Fig.4(b)



### Fe XXIII(1s<sup>2</sup>3pnp)

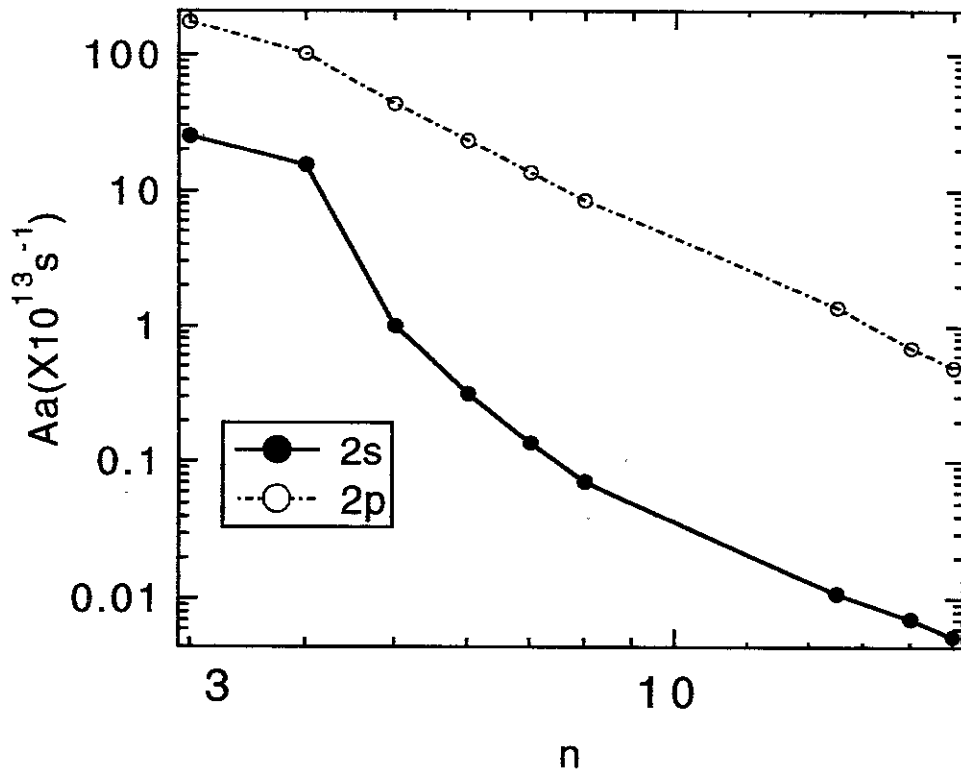


Fig.4(c)

### Fe XXIII(1s<sup>2</sup>3pnd)

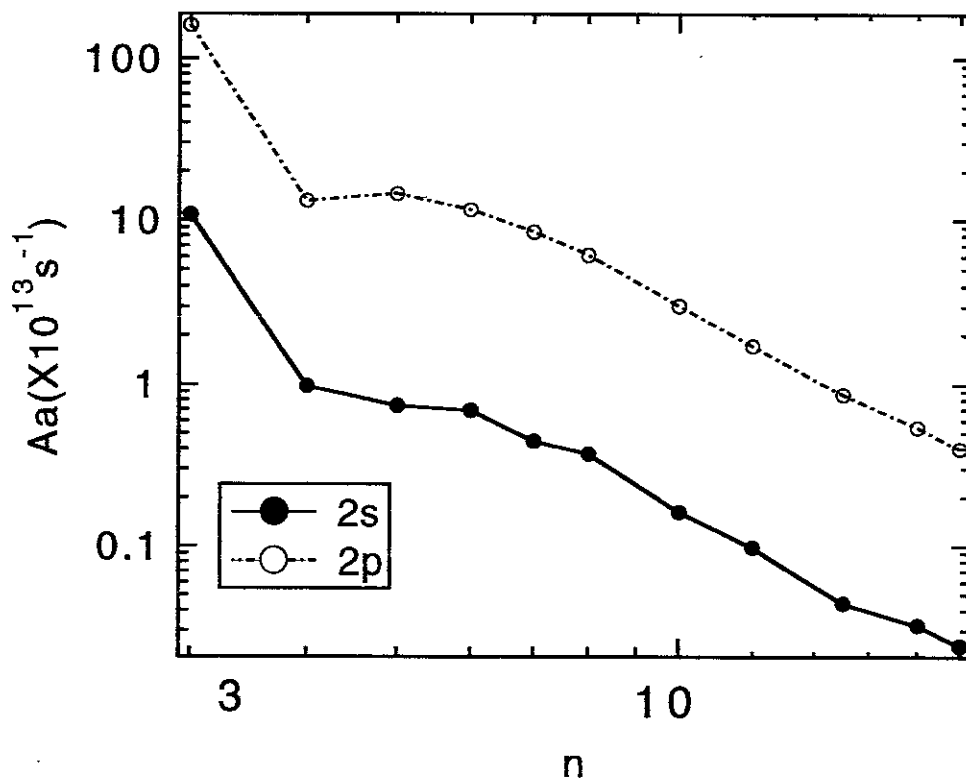


Fig.4(d)

### Fe XXIII(1s<sup>2</sup>3dnp)

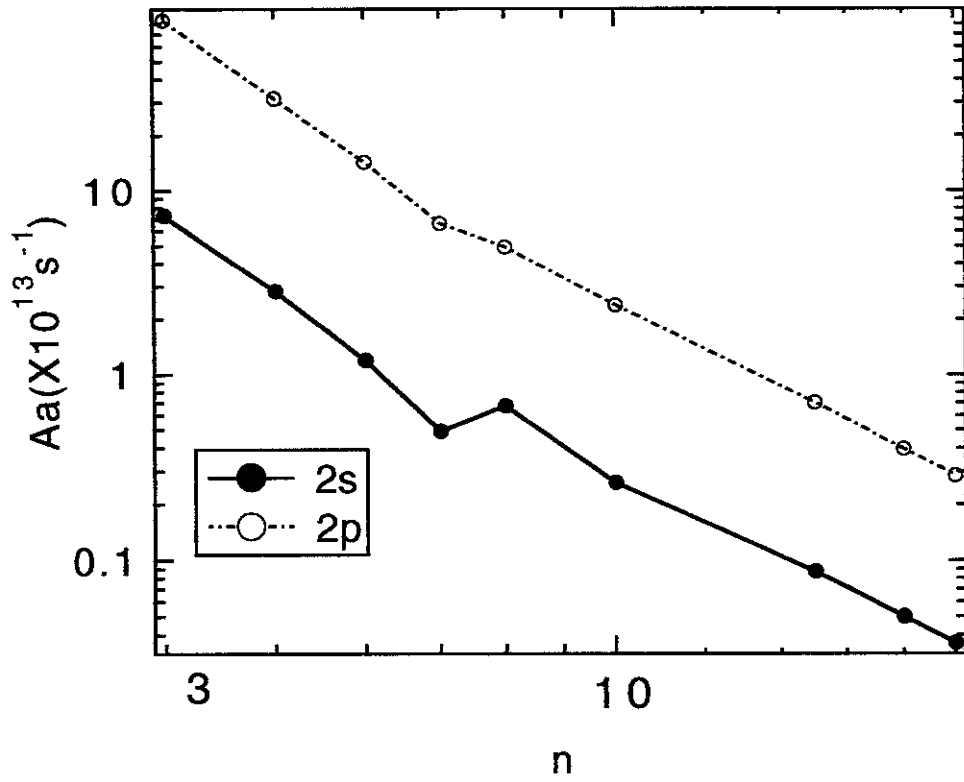


Fig.4(e)

### Fe XXIII(1s<sup>2</sup>3dnd)

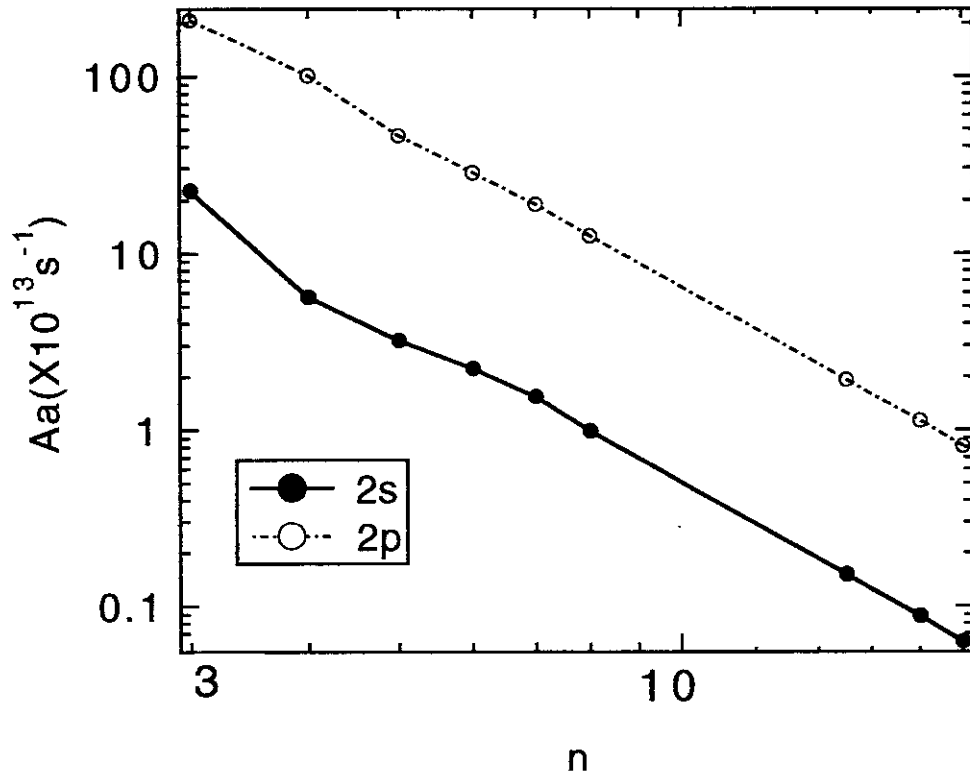


Fig.4(f)

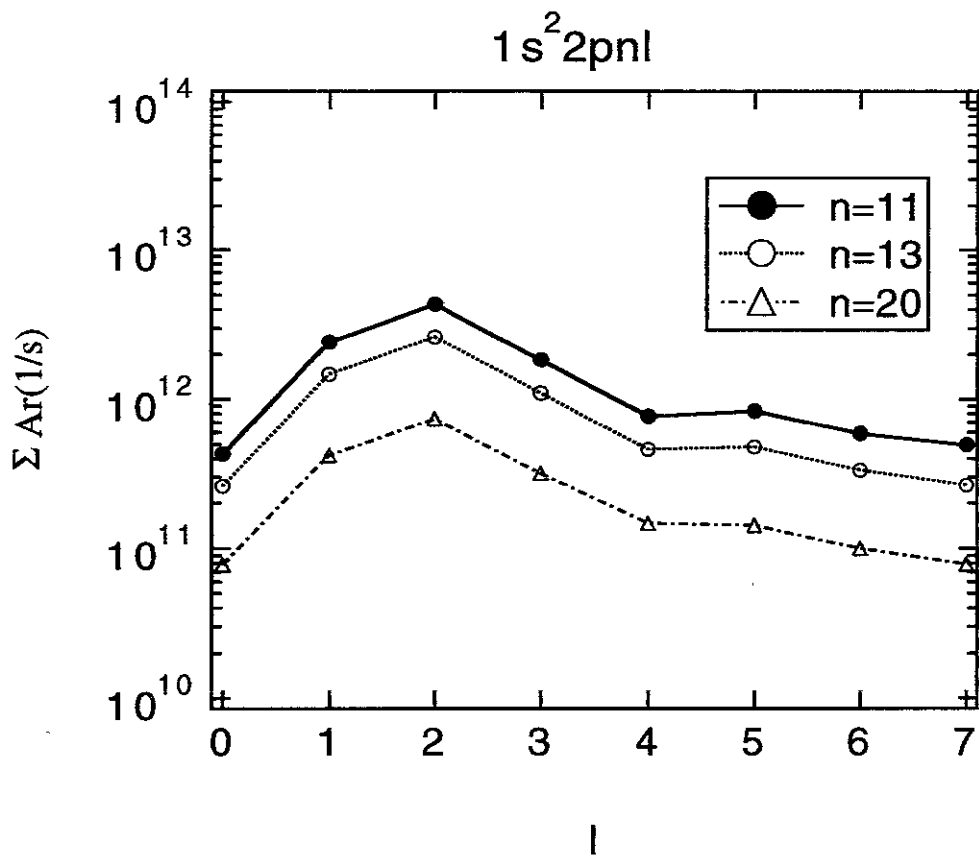


Fig.5

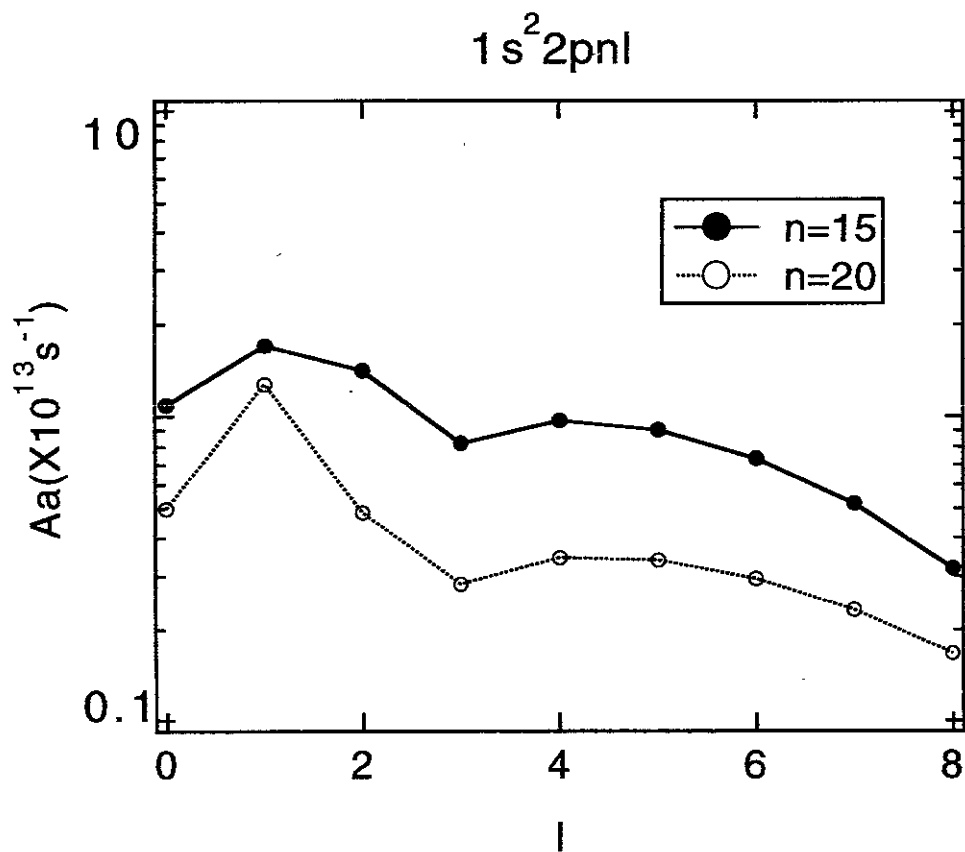


Fig.6

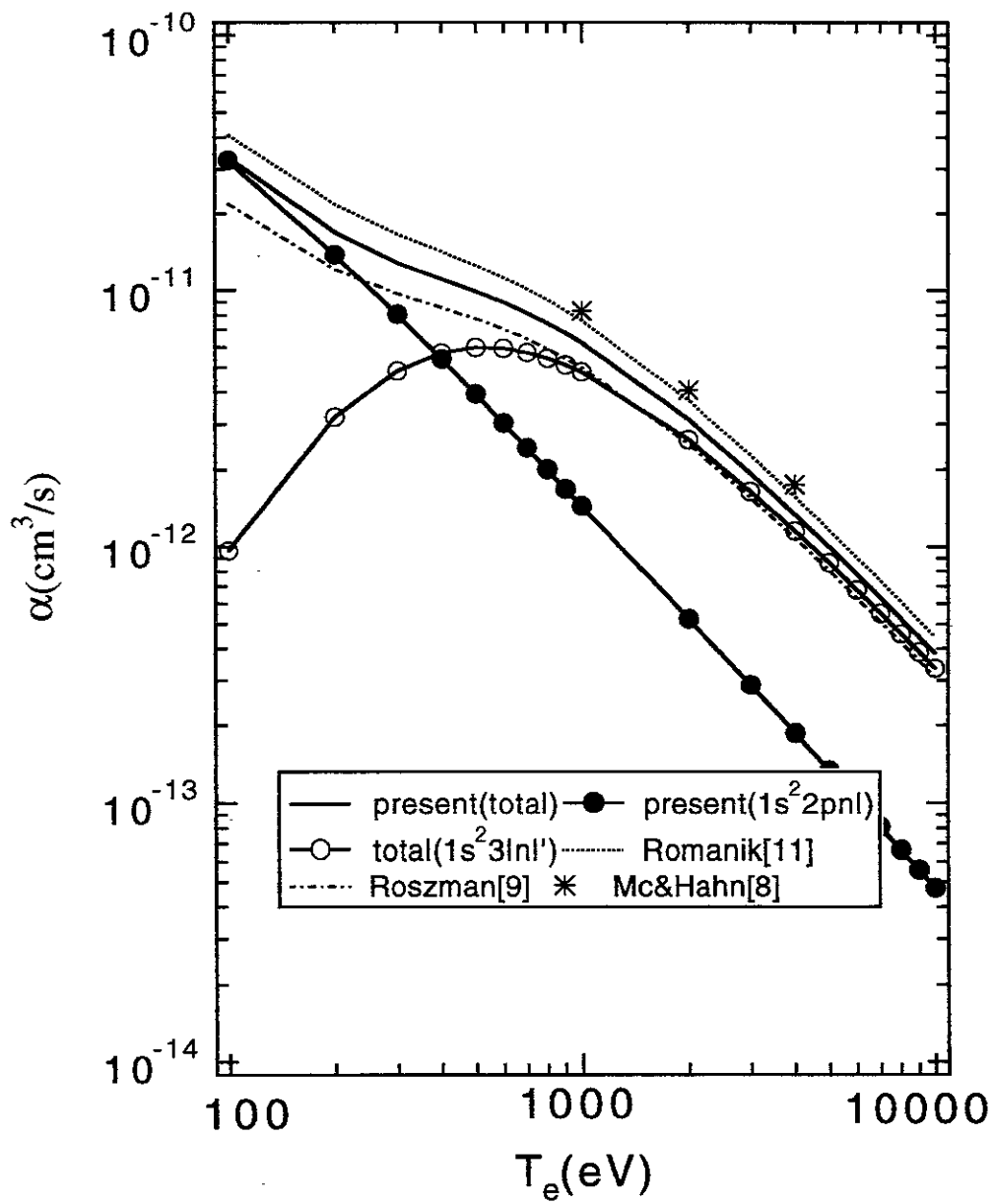


Fig.7

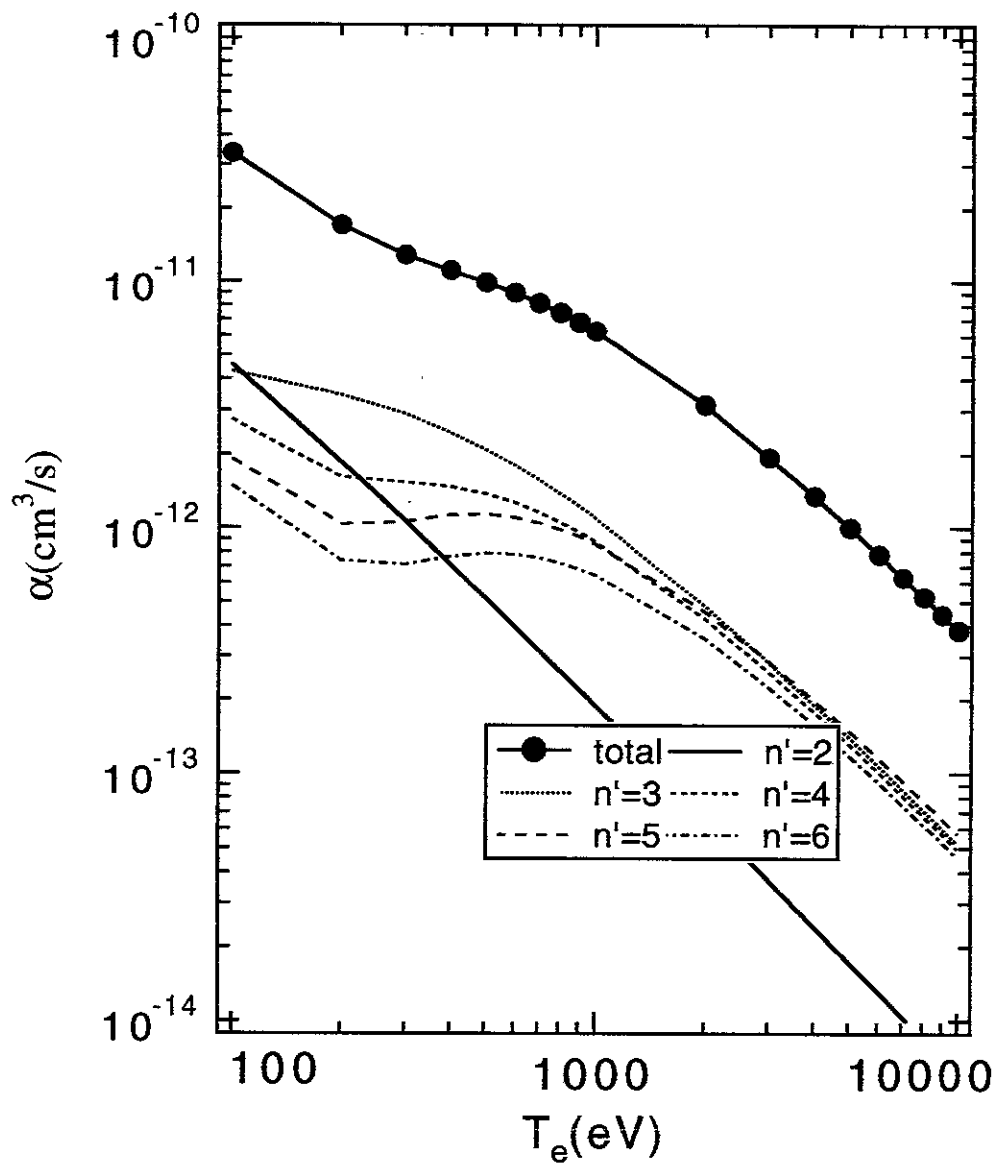


Fig.8

1 s<sup>2</sup> 2pnl

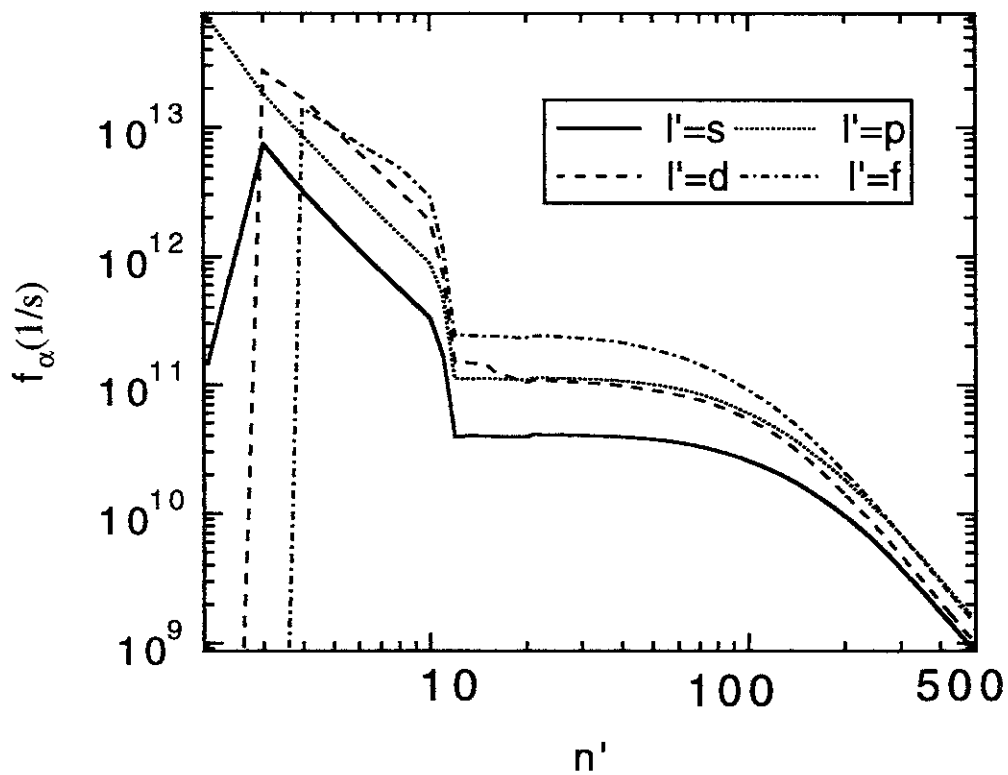


Fig.9(a)

1 s<sup>2</sup> 3pnl

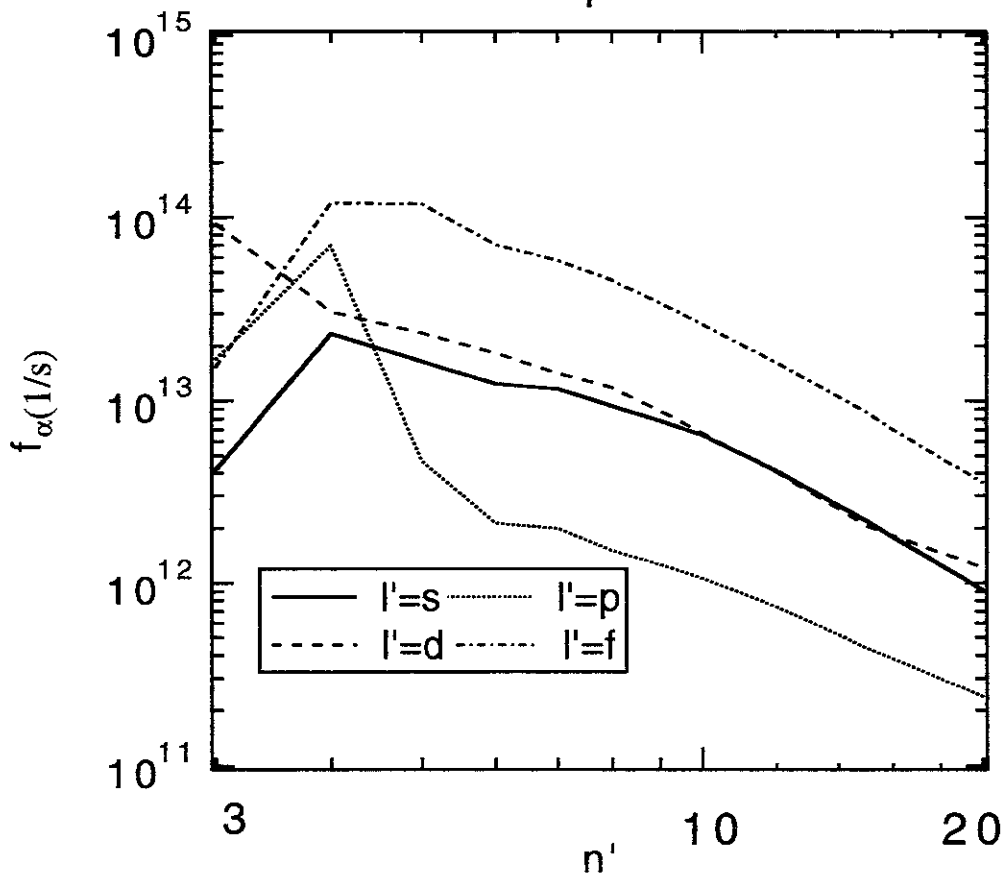


Fig.9(b)

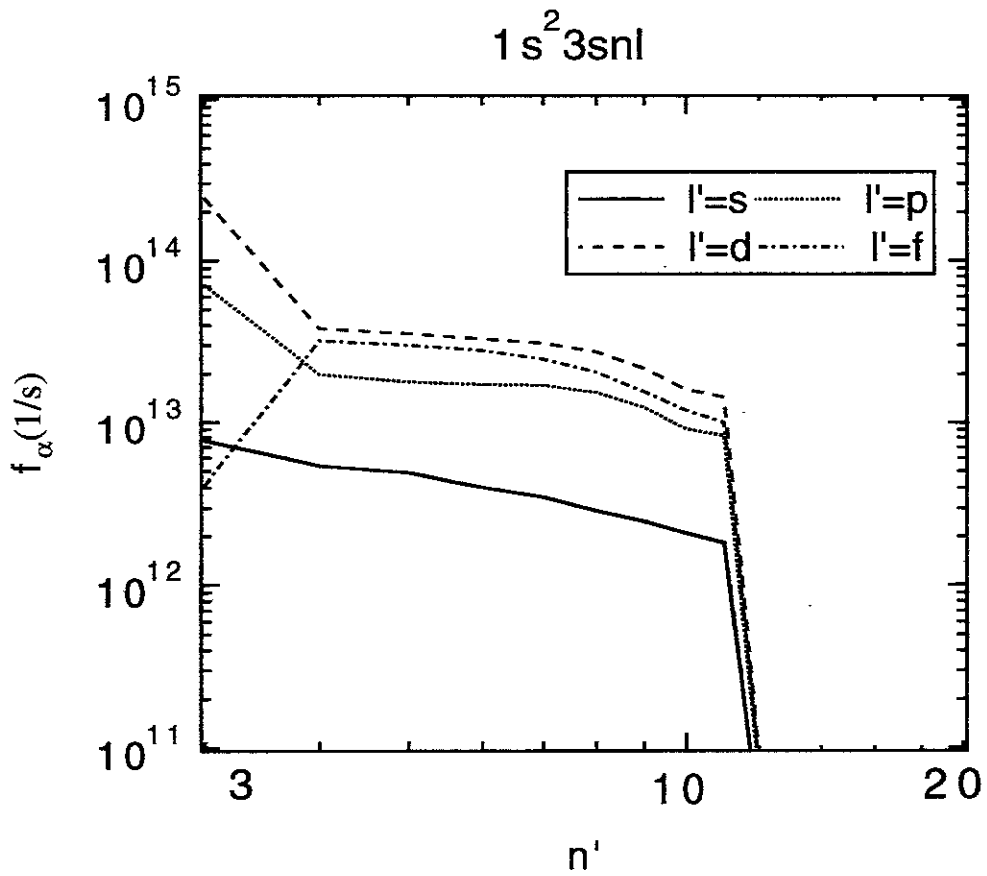


Fig.9(c)

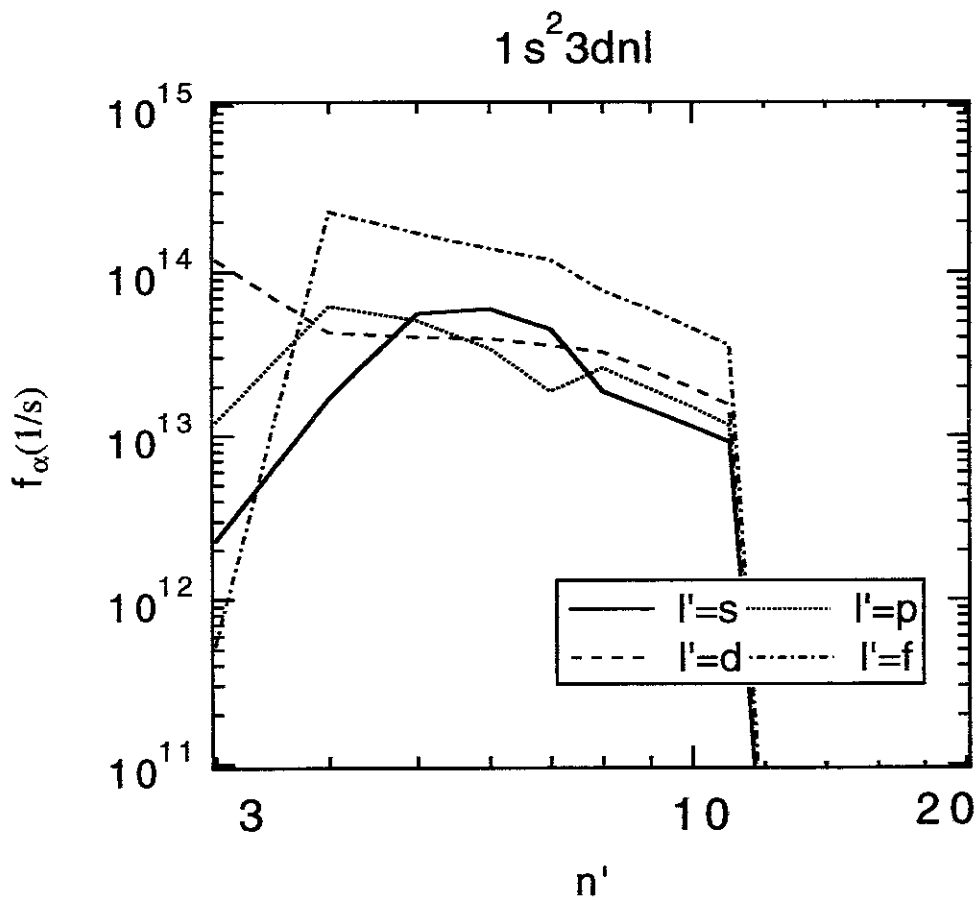


Fig.9(d)

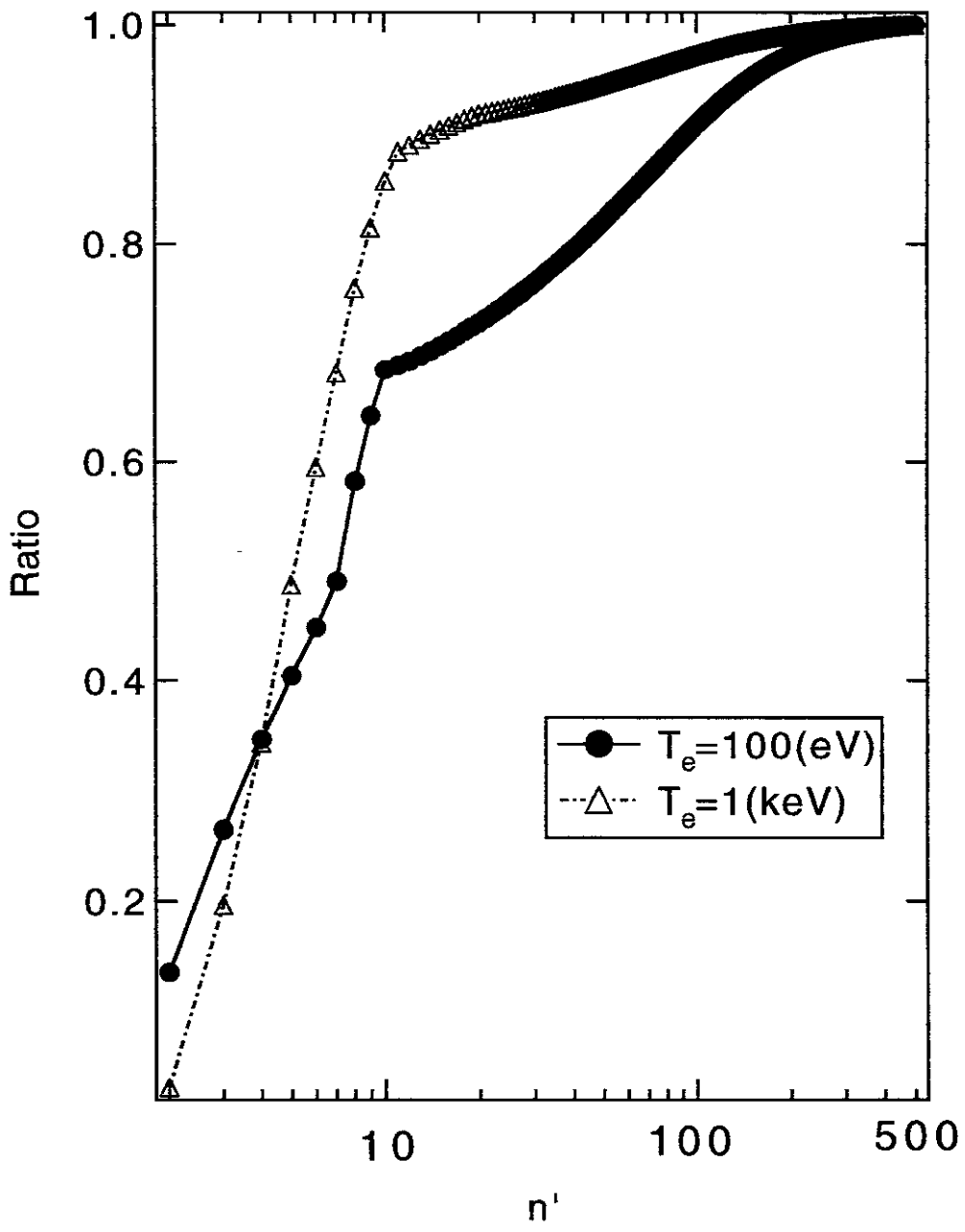


Fig.10



## Publication List of NIFS-DATA Series

- NIFS-DATA-1 Y. Yamamura, T. Takiguchi and H. Tawara,  
*Data Compilation of Angular Distributions of Sputtered Atoms*;  
Jan. 1990
- NIFS-DATA-2 T. Kato, J. Lang and K. E. Berrington,  
*Intensity Ratios of Emission Lines from OV Ions for Temperature  
and Density Diagnostics* ; Mar. 1990 [ *At Data and Nucl Data Tables*  
44(1990)133]
- NIFS-DATA-3 T. Kaneko,  
*Partial Electronic Straggling Cross Sections of Atoms for Protons*  
;Mar. 1990
- NIFS-DATA-4 T. Fujimoto, K. Sawada and K. Takahata,  
*Cross Section for Production of Excited Hydrogen Atoms  
Following Dissociative Excitation of Molecular Hydrogen by  
Electron Impact* ; Mar. 1990
- NIFS-DATA-5 H. Tawara,  
*Some Electron Detachment Data for H<sup>-</sup> Ions in Collisions with  
Electrons, Ions, Atoms and Molecules –an Alternative Approach to  
High Energy Neutral Beam Production for Plasma Heating–*;  
Apr. 1990
- NIFS-DATA-6 H. Tawara, Y. Itikawa, H. Nishimura, H. Tanaka and Y. Nakamura,  
*Collision Data Involving Hydro-Carbon Molecules*; July 1990  
[Supplement to *Nucl. Fusion* 2(1992)25]
- NIFS-DATA-7 H.Tawara,  
*Bibliography on Electron Transfer Processes in Ion-  
Ion/Atom/Molecule Collisions –Updated 1990–*; Aug. 1990
- NIFS-DATA-8 U.I.Safronova, T.Kato, K.Masai, L.A.Vainshtein and A.S.Shlyapzeva,  
*Excitation Collision Strengths, Cross Sections and Rate  
Coefficients for OV, SiXI, FeXXIII, MoXXXIX by Electron Impact  
(1s<sup>2</sup>2s<sup>2</sup> - 1s<sup>2</sup>2s2p - 1s<sup>2</sup>2p<sup>2</sup> Transitions)* Dec.1990
- NIFS-DATA-9 T.Kaneko,  
*Partial and Total Electronic Stopping Cross Sections of Atoms and  
Solids for Protons*; Dec. 1990
- NIFS-DATA-10 K.Shima, N.Kuno, M.Yamanouchi and H.Tawara,  
*Equilibrium Charge Fraction of Ions of Z=4-92 (0.02-6 MeV/u) and  
Z=4-20 (Up to 40 MeV/u) Emerging from a Carbon Foil*; Jan.1991  
[*AT.Data and Nucl. Data Tables* 51(1992)173]

- NIFS-DATA-11 T. Kaneko, T. Nishihara, T. Taguchi, K. Nakagawa, M. Murakami, M. Hosono, S. Matsushita, K. Hayase, M. Moriya, Y. Matsukuma, K. Miura and Hiro Tawara,  
*Partial and Total Electronic Stopping Cross Sections of Atoms for a Singly Charged Helium Ion: Part I*; Mar. 1991
- NIFS-DATA-12 Hiro Tawara,  
*Total and Partial Cross Sections of Electron Transfer Processes for  $Be^{q+}$  and  $B^{q+}$  Ions in Collisions with H,  $H_2$  and He Gas Targets - Status in 1991-*; June 1991
- NIFS-DATA-13 T. Kaneko, M. Nishikori, N. Yamato, T. Fukushima, T. Fujikawa, S. Fujita, K. Miki, Y. Mitsunobu, K. Yasuhara, H. Yoshida and Hiro Tawara,  
*Partial and Total Electronic Stopping Cross Sections of Atoms for a Singly Charged Helium Ion : Part II*; Aug. 1991
- NIFS-DATA-14 T. Kato, K. Masai and M. Arnaud,  
*Comparison of Ionization Rate Coefficients of Ions from Hydrogen through Nickel* ; Sep. 1991
- NIFS-DATA-15 T. Kato, Y. Itikawa and K. Sakimoto,  
*Compilation of Excitation Cross Sections for He Atoms by Electron Impact*; Mar. 1992
- NIFS-DATA-16 T. Fujimoto, F. Koike, K. Sakimoto, R. Okasaka, K. Kawasaki, K. Takiyama, T. Oda and T. Kato,  
*Atomic Processes Relevant to Polarization Plasma Spectroscopy* ; Apr. 1992
- NIFS-DATA-17 H. Tawara,  
*Electron Stripping Cross Sections for Light Impurity Ions in Colliding with Atomic Hydrogens Relevant to Fusion Research*; Apr. 1992
- NIFS-DATA-18 T. Kato,  
*Electron Impact Excitation Cross Sections and Effective Collision Strengths of N Atom and N-Like Ions -A Review of Available Data and Recommendations-* ; Sep. 1992
- NIFS-DATA-19 Hiro Tawara,  
*Atomic and Molecular Data for  $H_2O$ , CO &  $CO_2$  Relevant to Edge Plasma Impurities* , Oct. 1992
- NIFS-DATA-20 Hiro. Tawara,  
*Bibliography on Electron Transfer Processes in Ion-Ion/Atom/Molecule Collisions -Updated 1993-*; Apr. 1993

- NIFS-DATA-21 J. Dubau and T. Kato,  
*Dielectronic Recombination Rate Coefficients to the Excited States of C I from C II*; Aug. 1994
- NIFS-DATA-22 T. Kawamura, T. Ono, Y. Yamamura,  
*Simulation Calculations of Physical Sputtering and Reflection Coefficient of Plasma-Irradiated Carbon Surface*; Aug. 1994
- NIFS-DATA-23 Y. Yamamura and H. Tawara,  
*Energy Dependence of Ion-Induced Sputtering Yields from Monoatomic Solids at Normal Incidence*; Mar. 1995
- NIFS-DATA-24 T. Kato, U. Safronova, A. Shiyaptseva, M. Cornille, J. Dubau,  
*Comparison of the Satellite Lines of H-like and He-like Spectra*; Apr. 1995
- NIFS-DATA-25 H. Tawara,  
*Roles of Atomic and Molecular Processes in Fusion Plasma Researches - from the cradle (plasma production) to the grave (after-burning) -*; May 1995
- NIFS-DATA-26 N. Toshima and H. Tawara  
*Excitation, Ionization, and Electron Capture Cross Sections of Atomic Hydrogen in Collisions with Multiply Charged Ions*; July 1995
- NIFS-DATA-27 V.P. Shevelko, H. Tawara and E. Salzborn,  
*Multiple-Ionization Cross Sections of Atoms and Positive Ions by Electron Impact*; July 1995
- NIFS-DATA-28 V.P. Shevelko and H. Tawara,  
*Cross Sections for Electron-Impact Induced Transitions Between Excited States in He:  $n, n'=2,3$  and 4*; Aug. 1995
- NIFS-DATA-29 U.I. Safronova, M.S. Safronova and T. Kato,  
*Cross Sections and Rate Coefficients for Excitation of  $\Delta n = 1$  Transitions in Li-like Ions with  $6 < Z < 42$* ; Sep. 1995
- NIFS-DATA-30 T. Nishikawa, T. Kawachi, K. Nishihara and T. Fujimoto,  
*Recommended Atomic Data for Collisional-Radiative Model of Li-like Ions and Gain Calculation for Li-like Al Ions in the Recombining Plasma*; Sep. 1995
- NIFS-DATA-31 Y. Yamamura, K. Sakaoka and H. Tawara,  
*Computer Simulation and Data Compilation of Sputtering Yield by Hydrogen Isotopes ( $^1\text{H}^+$ ,  $^2\text{D}^+$ ,  $^3\text{T}^+$ ) and Helium ( $^4\text{He}^+$ ) Ion Impact from Monoatomic Solids at Normal Incidence*; Oct. 1995

- NIFS-DATA-32 T. Kato, U. Safronova and M. Ohira,  
*Dielectronic Recombination Rate Coefficients to the Excited States of CII from CIII*; Feb. 1996
- NIFS-DATA-33 K.J. Snowdon and H. Tawara,  
*Low Energy Molecule-Surface Interaction Processes of Relevance to Next-Generation Fusion Devices*; Mar. 1996
- NIFS-DATA-34 T. Ono, T. Kawamura, K. Ishii and Y. Yamamura,  
*Sputtering Yield Formula for B<sub>4</sub>C Irradiated with Monoenergetic Ions at Normal Incidence*; Apr. 1996
- NIFS-DATA-35 I. Murakami, T. Kato and J. Dubau,  
*UV and X-Ray Spectral Lines of Be-Like Fe Ion for Plasma Diagnostics*; Apr. 1996
- NIFS-DATA-36 K. Moribayashi and T. Kato,  
*Dielectronic Recombination of Be-like Fe Ion*; Apr. 1996