§36. Characteristic Evolution of Electron and Ion Temperatures in a Plasma with Non-Monotonic Rotational Transform Profile

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In a tokamak plasma, the reversed magnetic shear configuration where the global magnetic shear changes negative to positive toward the edge having a zero shear layer has received much interest because an internal transport barrier is easily formed and it is a promising candidate for steady-state operation in ITER [1]. It is interesting and significant to investigate whether or not a configuration with the reversed magnetic shear can be realized in stellarator /helical plasmas and how the plasma in the configuration behaves.

This configuration was realized in an LHD plasma where large counter plasma current was driven by neutral beam injection (NBI)[2, 3]. A typical discharge waveform is shown in Fig.1, where the toroidal field is  $B_t$ =1.3T and the magnetic axis position of the vacuum field is  $R_{ax}$ =3.75m. The plasma current reaches ~140 kA by two counter NBI systems having about 2MW absorption power and about 180 keV hydrogen beams. The line averaged electron density is kept constant ( $< n_e > \approx 0.6 \times 10^{19} \text{ m}^{-3}$ ). Appreciable amount of neon is puffed into a plasma by a short pulse (~70ms) at t~0.8s to enhance NB power absorption and minimize electron return current in NB current drive, where the effective charge  $Z_{eff} \sim 6$ . Electron temperature profile is a parabolic shape having  $\sim 1.3$  keV at the plasma center in the constant  $\langle n_e \rangle$  phase. Electron density profile is considerably hollow, where  $n_e(0)/n_{emax}$   $\sim 0.5$ . The beta value of beam component  $\langle \beta_b \rangle$  is comparable or higher than that of bulk plasma ( $\sim 0.5\%$ ). That is, energetic ion effect is expected to be significant in these plasmas. The rotational transform  $\iota/2\pi$ -profile measured by motional-stark-effect (MSE) spectroscopy is non-monotonic, as shown in Fig.2. In the time window shown in Fig.1, the position of the minimum  $\iota/2\pi$  moves inward from  $\rho \sim 0.7$  to  $\rho \sim 0.4$ , and the minimum value gradually decreases, passing through the rational values  $\iota/2\pi=1/2$  and 1/3. At these timings, electron temperature measured by electron cyclotron emission (ECE)  $T_e^{ECE}$  exhibits sharp drops at  $t=t_1$  and  $t=t_2$  in Fig.1.

In a similar shot with that shown in Fig.1, the central ion temperature measured by soft X-ray crystal spectroscopy has increased linearly for more than global energy confinement time from the time when the minimum of the rotational transform approaches or passes  $v/2\pi=1/3$ , as shown in Fig.3. The mechanism of this ion temperature rise is under investigation.

- [1] Editors of "Progress in the ITER Physics Basis", Nucl. Fusion 47, S1 (2007).
- [2] Toi, K., et al., 10<sup>th</sup> IAEA Technical Meeting on Energetic Particles in Magnetic Confinement

Systems, Kloster Seeon (Germany), 8-10 Oct., 2007. [3] Toi, K., et al., 35th EPS Plasma Physics Conference, Crete (Greece), 9-13 June 2008, paper No.P1.054.

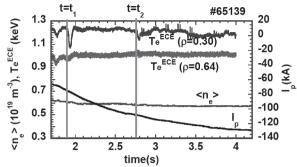


Fig.1 Typical waveform of a RS-plasma of LHD.  $T_e^{ECE}$  has sharp drops when the minimum of the rotational transform has passed the rational values 1/2(at  $t=t_1$ ) and 1/3(at  $t=t_2$ )(vertical lines).

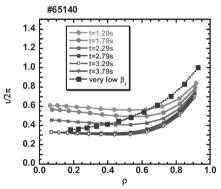


Fig.2 Time evolution of the rotationalTransform profile measured by MSE. The dotted curve indicates the profile in the vacuum.

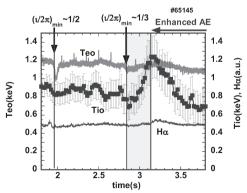


Fig.3 Time evolutions of  $T_{io}$  and  $T_{eo}$  (ECE) in a low density RS plasma of  $< n_e > = 0.5 \times 10^{19} \text{ m}^{-3}$ .

H $\alpha$  emission indicates an enhanced loss from  $t \sim 3.1 \text{s}$ .