## §4. Behaviors of He II in Two-dimensional Channels Filled with He II

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Two-dimensional channels composed of the good conductor are common in a He II cooled apparatus. In the 2D-channels, the heat transfer from to He II below the  $\lambda$ -point pressure  $P_{\lambda}$  is largely influenced by the superheating in He II (<sup>S</sup>He II). So far, we have reported that the heated surface of a heated conductor is partially insulated with the superheated He I (<sup>S</sup>He I) in He II below the  $\lambda$ -point pressure  $P_{\lambda}$ [1]. That is, <sup>S</sup>He I appears forming an intermediate state in heat transfer characteristics such that the transition at  $Q_{\lambda}$  is gradual like in the pressurized He II. In the present paper, we confirm that the intermediate state consists of the superheated phases. The appearance of the intermediate state gives a considerable margin for the stabilization of the He II cooled apparatus.

The experiment was performed by the use of a copper disk with the radius of 10 mm with the thickness of 5 mm. An insulator disk was placed in parallel to a plane surface of the copper disk as shown in Fig. 1. The heat flux Q was applied from a thermo-foil in a vacuum can. To measure the temperature of helium in the channel, chipresistors  $T_{1-6}$  were arrayed on the insulator disk at intervals of 2 mm. The temperature of the copper disk  $T_C$  was measured with thermometer  $T_C$ . Pressures in the channel  $P_c$  and in the bath  $P_b$  were measured with in-situ pressure gauges. To measure the stability of superheated states, a pendulum with a metal ball was prepared to disturb the cryostat mechanically. The threshold potential energy  $E_M$  of the pendulum to break the metastability phases was taken as an index of the stability.

It is estimated from Fig. 2 that not the bubble but <sup>S</sup>He I nucleates in the hottest area on the heated surface covered with <sup>S</sup>He II, when  $T_1$  crosses the  $\lambda$ -line extended below  $P_{\lambda}$ . The small heat conduction of <sup>S</sup>He I drives a part of Q downstream through the conductor without the sharp transition at  $Q_{\lambda}$ .  $T_1$  increases over  $T_{\lambda}$  on an isobar without boiling until  $Q_n$  reaches as shown in Fig. 2b, that is,  $\Delta P = 0$ , where  $\Delta P = P_c - P_b$ . When Q is decreased from above  $Q_n$ ,  $T_1$  jumps beyond  $T_{\lambda}$  at a critical heat flux of recovery  $Q_r$ . Both the isobaric behaviors up to  $Q_n$  and the sudden changes in  $\Delta P$  at  $Q_n$  and at  $Q_r$  suggest that the boiling does not occur in the intermediate state. This also means that <sup>S</sup>He II layer surrounds coaxially the <sup>S</sup>He I layer in the intermediate state.

The disturbance with  $E_{\rm M}$  breaks momentarily superheated states.  $E_{\rm M}$  can be an indirect index of the metastability. However, the metastable phases, <sup>S</sup>He II and <sup>S</sup>He I reappear in a second after the transient collapse (the inset of Fig. 3). That is, <sup>S</sup>He I and <sup>S</sup>He II are apparently stabilized such that they return immediately after a mechanical shock disrupts them. The stability decreases with increasing *Q*.

The intermediate state where <sup>S</sup>He I coexists with <sup>S</sup>He II is succeeded by the mixed state where the alternation of superheating and boiling is sustained. Above  $P_{\lambda}$ , by contrast, the temperature of the conductor in 2D-channel rises steadily due to the stable and viscous subcooled He I.

1). Kobayashi H. et al. Proc ICEC21 2006:389-392

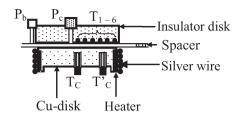


Fig. 1. Schematics of tested 2D-channel The 2D-channel is formed between the insulator disk and the copper disk.  $T_{1-6}$ ,  $T_C$ : the thermometer,  $T'_C$ : the spare.  $P_c$ ,  $P_b$ : In-situ pressure gauges. The thermo-foil is fastened to the wall of the copper disk with the silver wire.

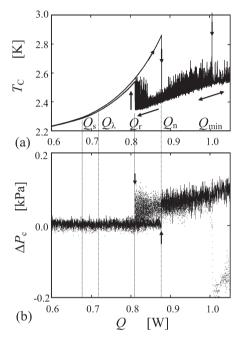


Fig. 2. Heat transfer characteristics and the pressure change

(a)  $T_C - Q$ , (b)  $\Delta P - Q$ ,  $\Delta P = P_c - P_b$  $T_b = 1.95$  K,  $P_b = 4.50$  kPa, gap distance: 0.15 mm

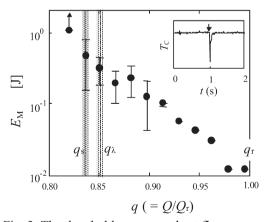


Fig. 3. The threshold energy vs. heat flux  $E_{\rm M}$ : the threshold potential energy to breaks the superheated states. The heat flux is normalized with  $Q_{\rm r}$ . The arrow mark in the inset: the moment of hitting, the bars: standard errors,  $T_{\rm b} = 1.95$  K,  $P_{\rm b} = 4.25$  kPa,  $Q_{\lambda} = 0.68$  W, the gap distance: 0.15 mm