

# **Formation and sustainment of tokamak equilibrium with a current hole in JT-60U**

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# Introduction

JT-60U

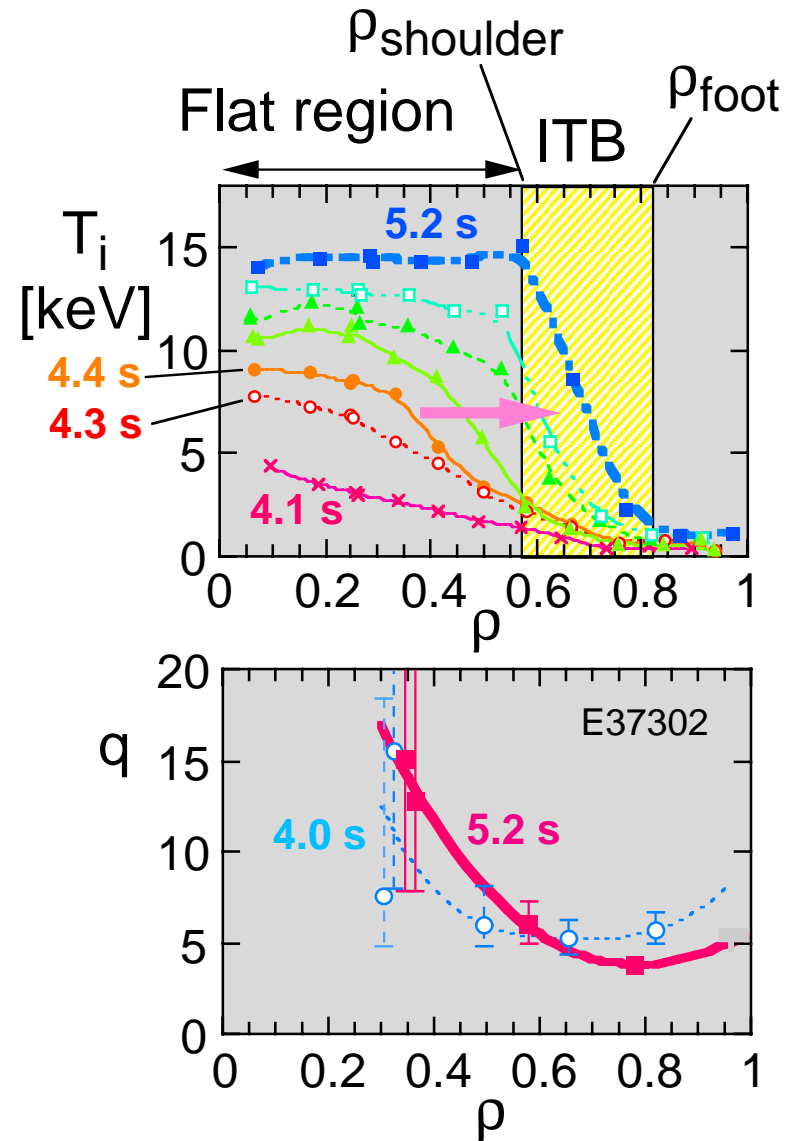
In JT-60U reversed shear plasmas,

The MSE data indicate  $q(0)$  is very high ( $B_p$  is very low in the central region).

Since the confinement of particles in toroidal systems depends on the poloidal field, it was an important issue to evaluate how high  $q(0)$  was.

However, the resolution of MSE was insufficient to address this issue, because the error in  $q$  becomes large by nature near the axis where  $B_p/B_t$  is small.

The MHD equilibrium code was unable to deal with a high  $q(0)$  equilibrium.



➔ Improvement in MSE and equilibrium code

# Improvement in MSE

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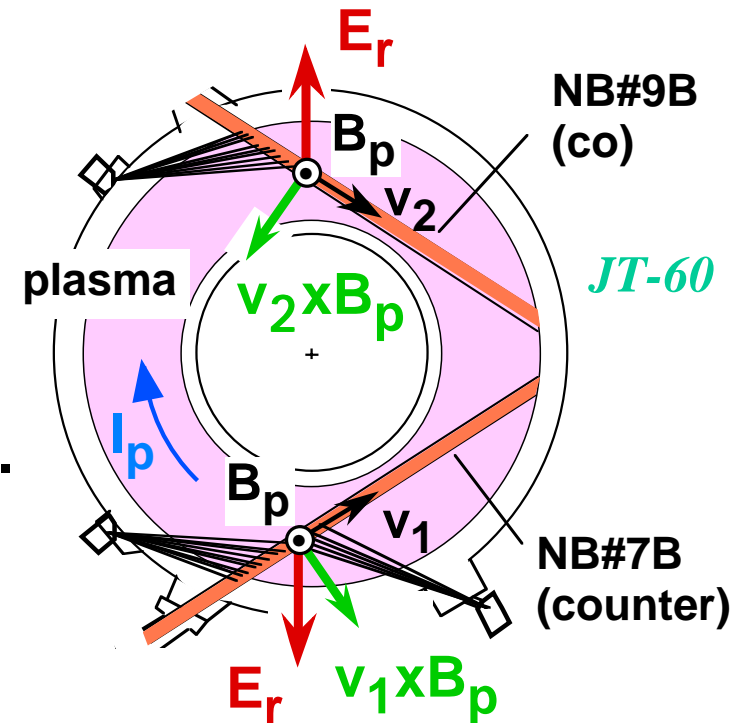
## (1) Correction of $E_r$ effects ( $E_r$ measurement)

MSE measures the direction of  
 $E_{\text{tot}} = E_r + v \times B$

Conventional MSE views one of counter-tangential beams.

New MSE system viewing one of co-tangential beams was installed in 2000.

From  $E_r + v_1 \times B$  and  $E_r + v_2 \times B$ ,  
 $E_r$  and  $B$  ( $B_p$ ) can be separated.



## (2) Improvement in calibration (correction of changes in mirror properties)

Dielectric multilayer mirrors are used in MSE.

Their polarization properties depends the temperature, which is raised to  
 $\sim 70$  degC during the experiment (300 degC VV).

Calibration was done heating up the mirrors to same temperature.

# Improvement in MHD equilibrium code

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Generation of various current profiles has become possible to fit MSE data by following improvements.

- (1) Function for  $j(r)$  ( $dp/d\psi$ ,  $FdF/d\psi$ ) is changed from a polynomial of  $\psi$  (poloidal flux) to **a spline of  $\rho$  (normalized radius)**.
- (2) Introducing a deceleration factor in the iteration of Grad-Shafranov equation solver.

Grad-Shafranov equation;

$$-\Delta^*\Psi = \mu_0 R j(R, \Psi)$$

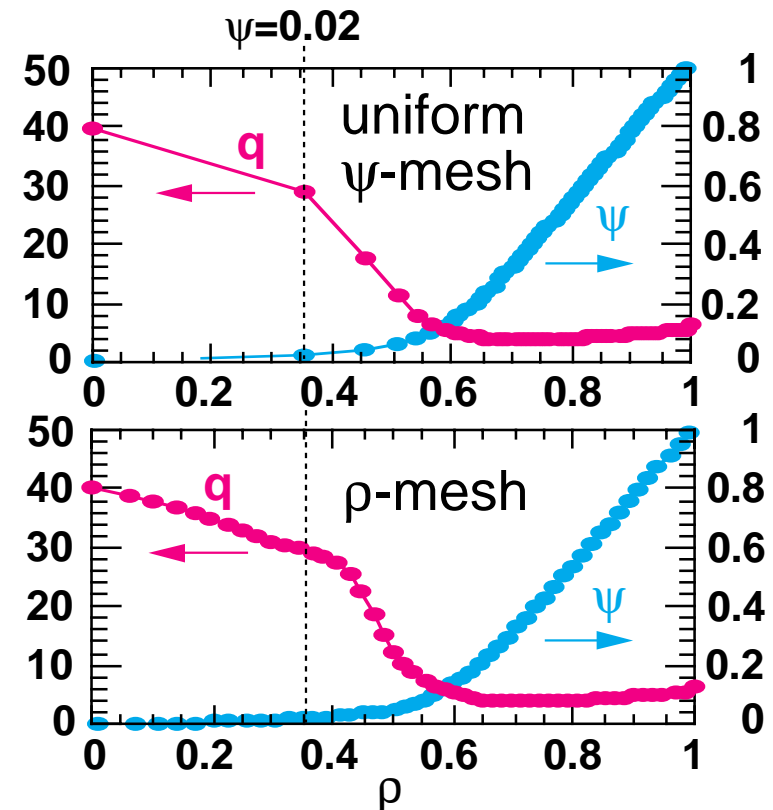
$$j(R, \Psi) = R p'(\Psi) + (1/\mu_0 R) * F(\Psi) F'(\Psi).$$

In an equilibrium with high  $q(0)$ ,  $\Psi$  is not a proper parameter to specify a position.

$Y(\Psi)$  [=  $p'(\Psi)$  or  $F(\Psi)F'(\Psi)$ ] is represented by a third-order spline function of  $\rho$ ,  $g(\rho)$ ;

$$Y(\Psi) = g(\rho(\Psi)).$$

$Y(\Psi)$  or  $\rho(\Psi)$  is evaluated on a mesh in a real space ( $\sim\rho$ ) NOT in a  $\Psi$  space.



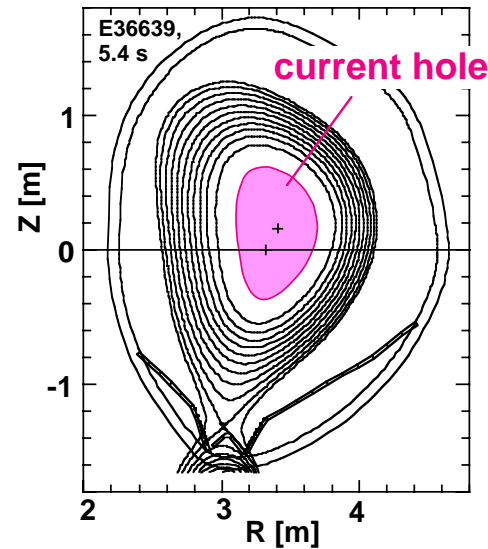
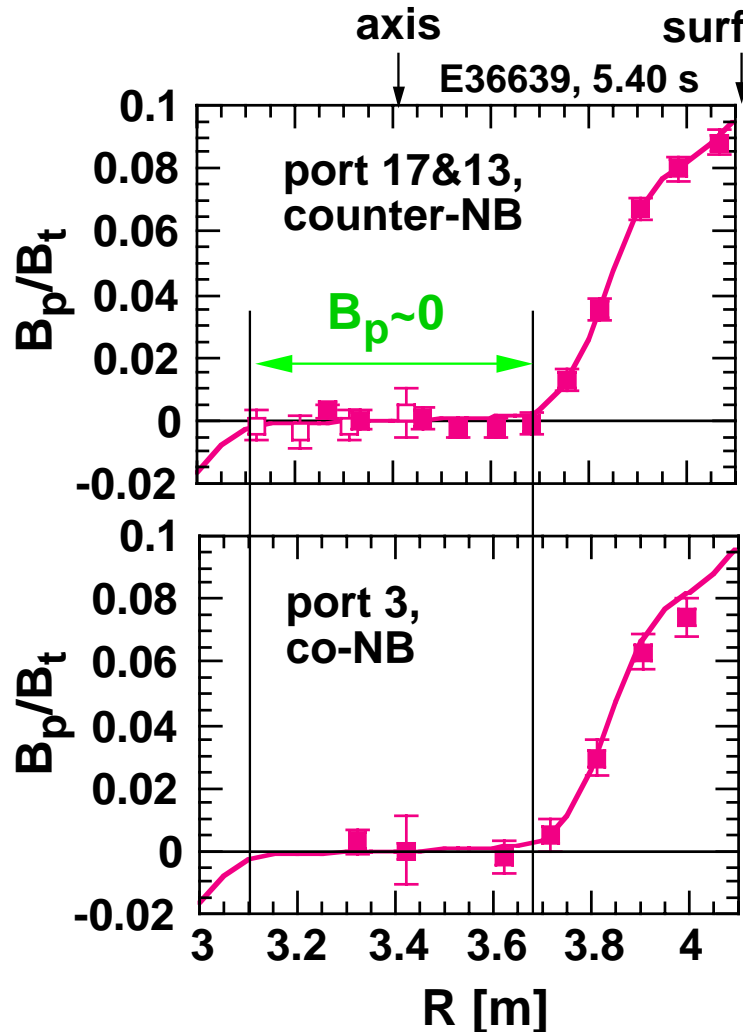
# $B_p \sim 0$ was observed in a central region

- Projected angle was  $\sim 0$  in a central region for both MSE viewing co and counter beams.

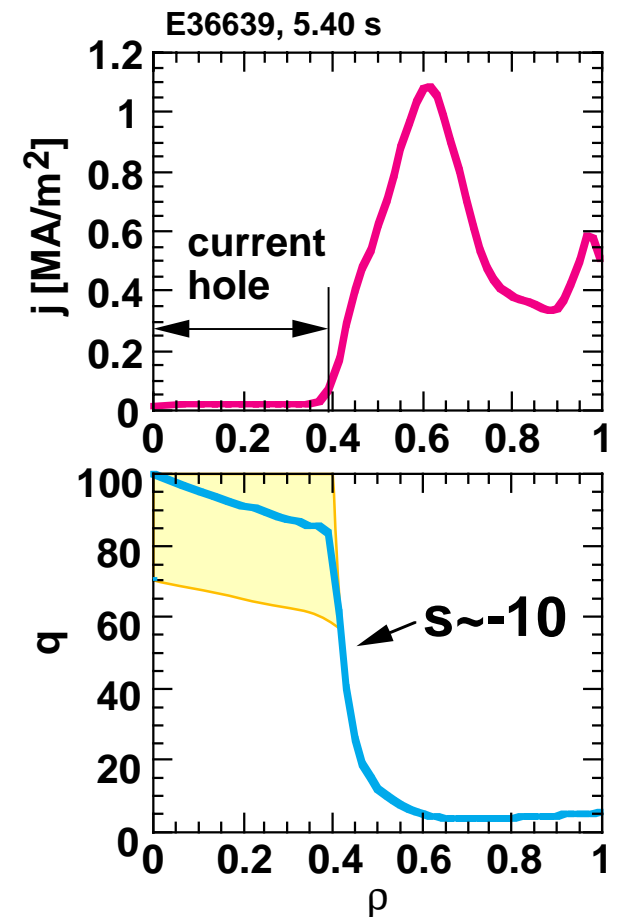
→  $E_r$  effect was small and  $B_p \sim 0$ .

- Equilibrium with  $q(0) \sim 100$  almost agrees with MSE data.
- Very small current in a central region;  $|j(0)| < \sim 0.07 \langle j \rangle$

→ "Current hole"



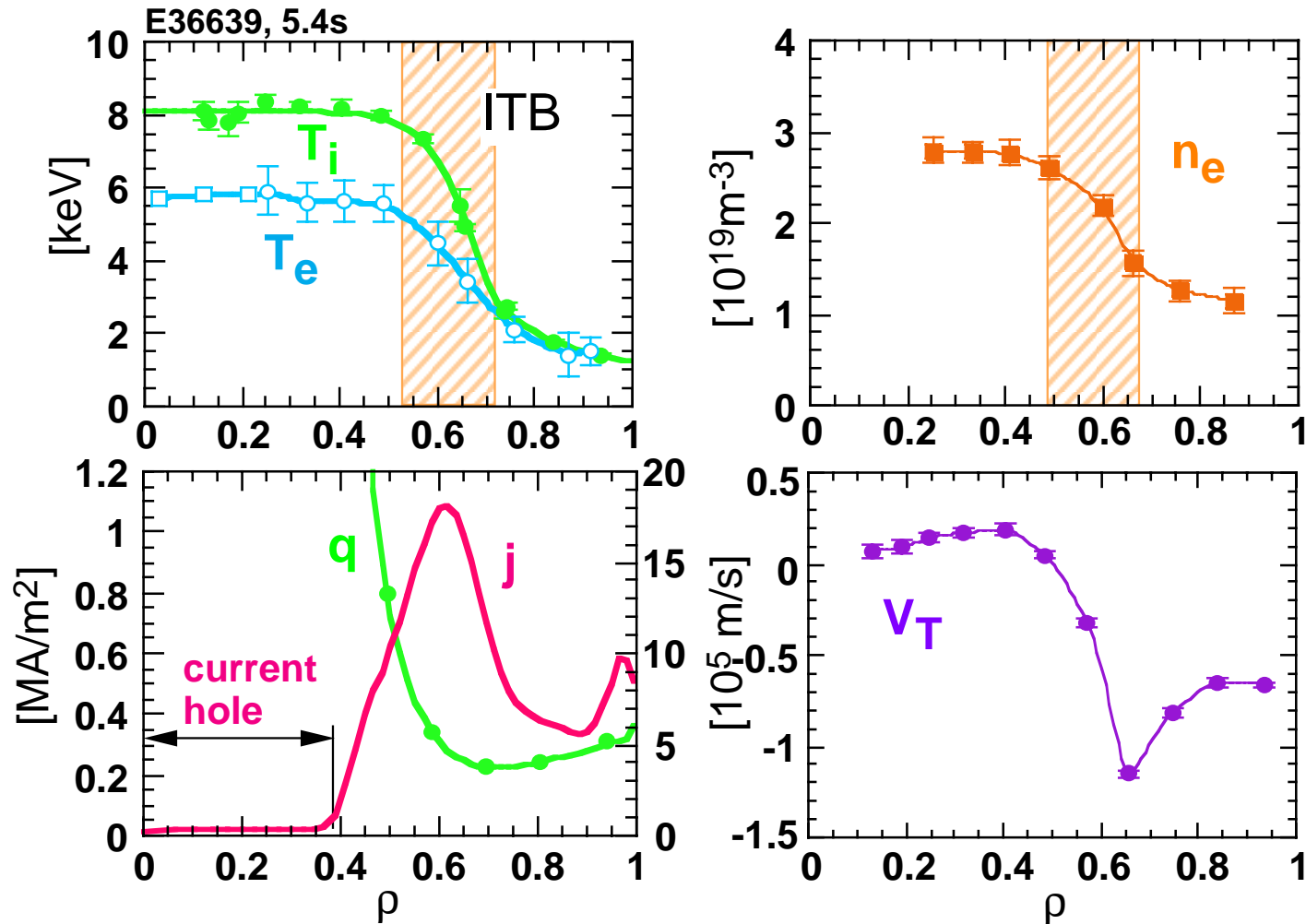
$B_t = 3.7$  T  
 $I_p = 1.35$  MA  
 $q_{95} = 5.2$



# High temperature plasma in the current hole

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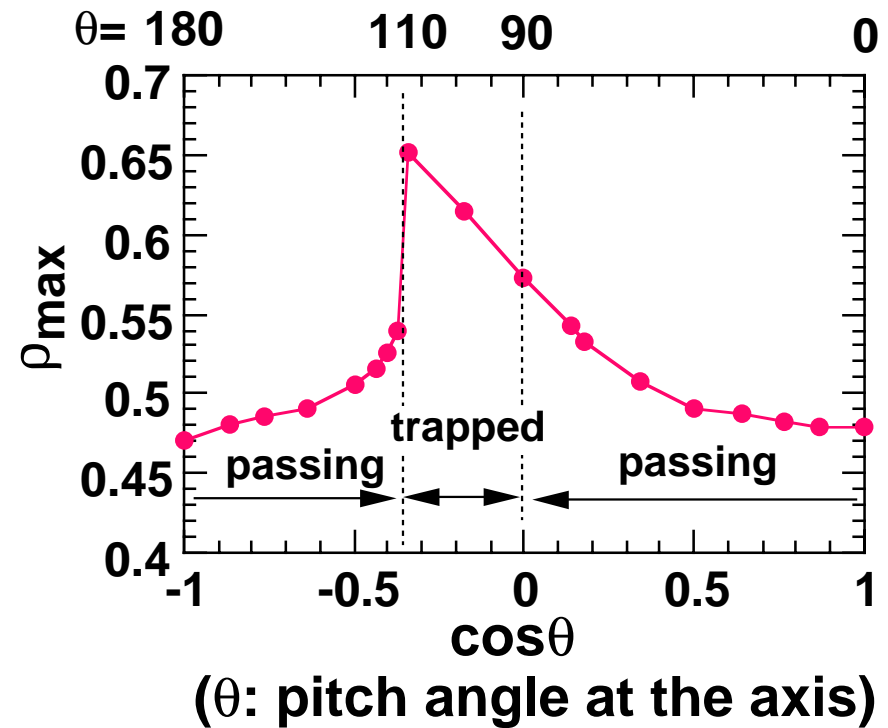
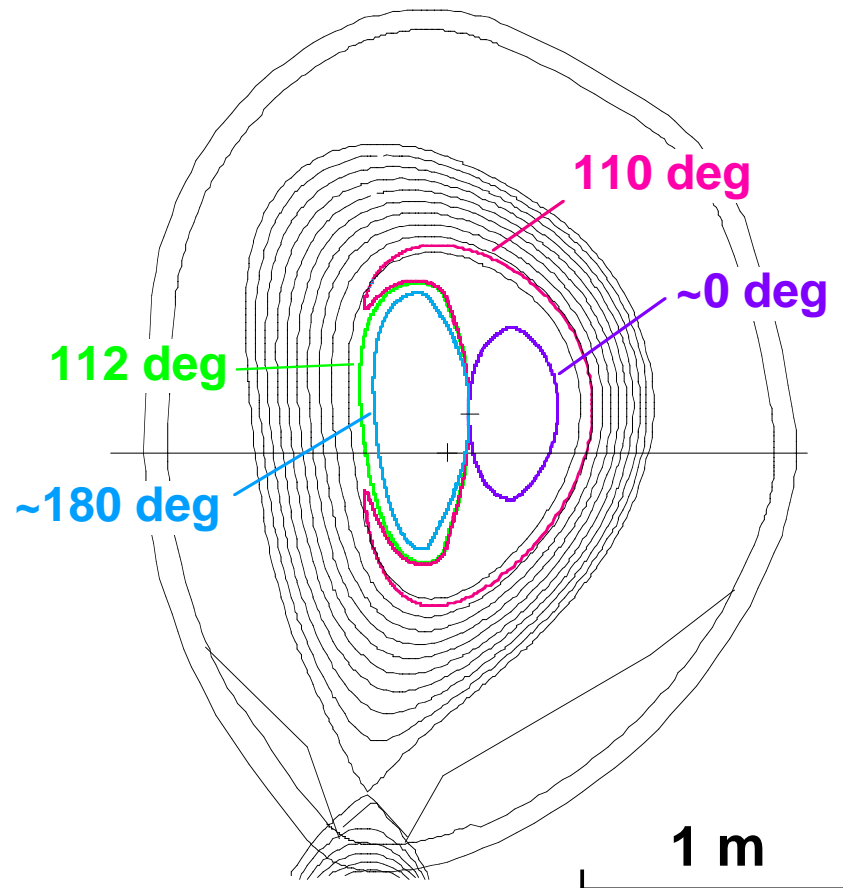
- $T(r)$  and  $n(r)$  are flat in the current hole, but steep gradients (ITBs) are formed outside the current hole where  $j(r)$  is peaked.
- High temperature plasmas can be confined in the central region only by  $B_p$  at a half radius.



# Orbit of thermal ions

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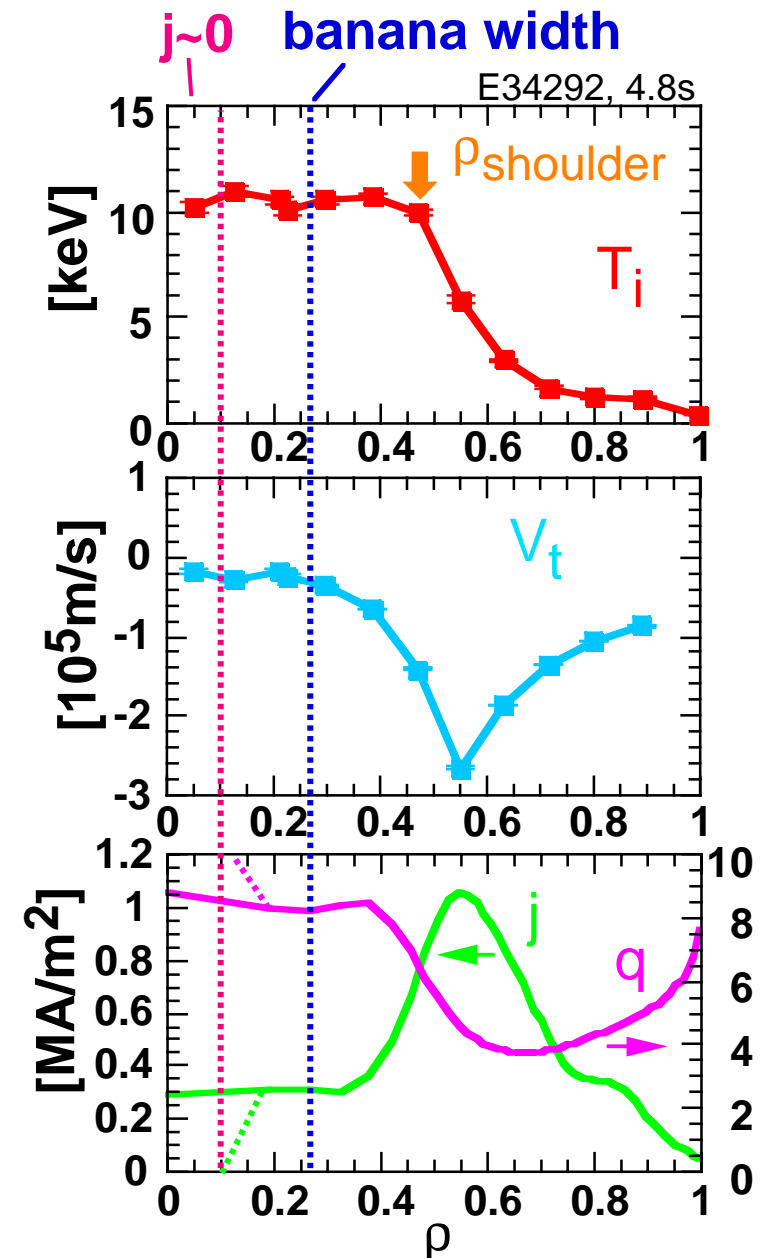
- Orbit of thermal ions with 8 keV ( $=T_i(0)$ ), passing through the axis, was calculated using the equilibrium with  $q(0)=100$ .
- The largest banana extends to  $\rho \sim 0.65$ .
- Even passing ions drift to  $\rho > \sim 0.47$  ( $\sim \rho_{\text{shoulder}}$ )



# Shoulder radius can be larger than banana width

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- No current hole or small current hole with  $\rho < 0.1$  if any.
- The width of largest banana of thermal ions is  $\rho \sim 0.25$  while  $\rho_{\text{shoulder}} \sim 0.45$ .
- The radius of flat region of  $V_t$  is  $\rho \sim 0.25$  and equal to the banana width.

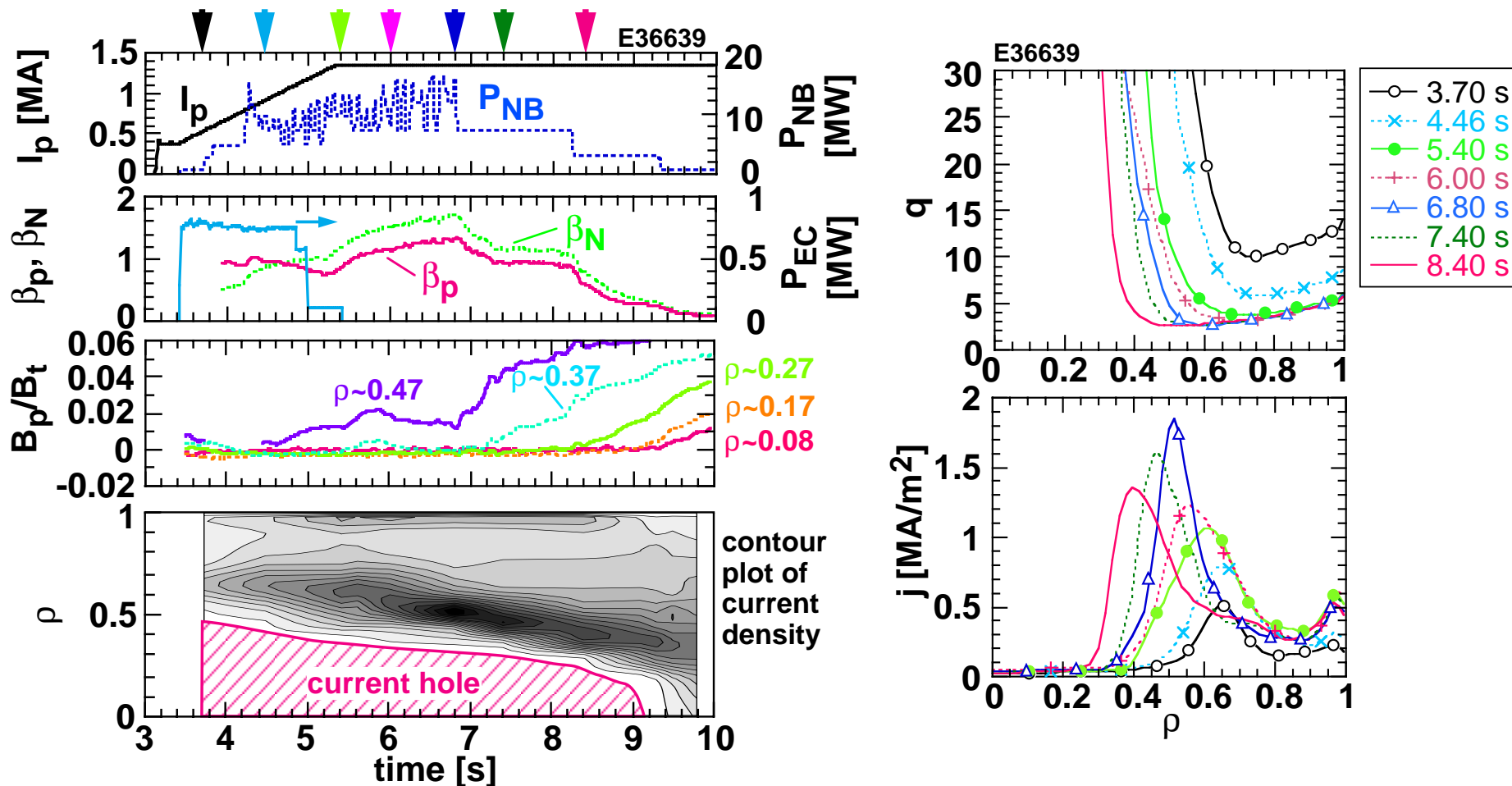




# Current hole was sustained stably

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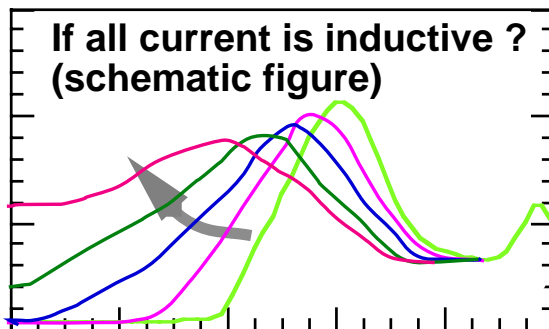
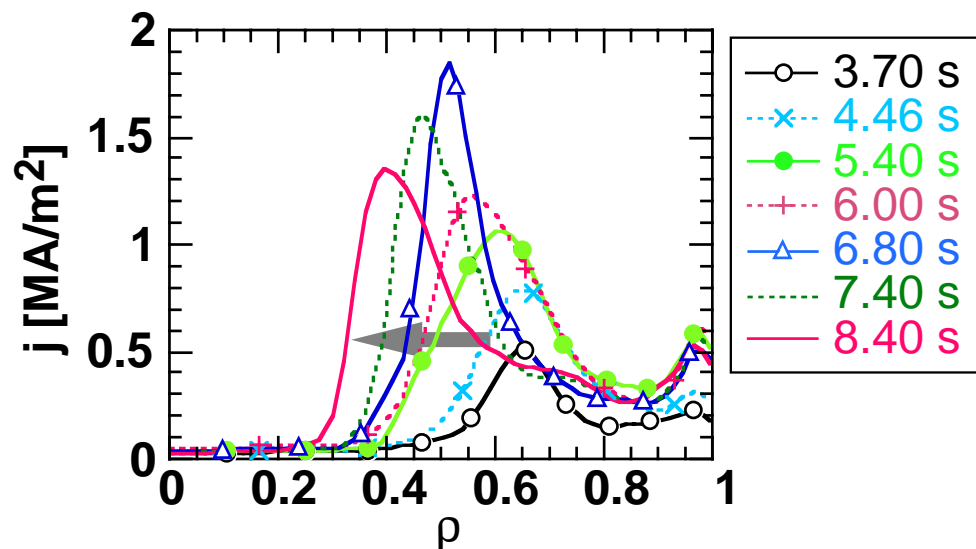
- The current hole was sustained for ~5 seconds without any global instabilities though its radius continued to shrink due to the penetration of inductive current.
- High confinement ( $HH_{98y2} < \sim 1.5$ ) and moderate beta ( $\beta_N < \sim 1.7$ ) were obtained, suggesting possibility of stable operation of tokamak reactors without on-axis current.



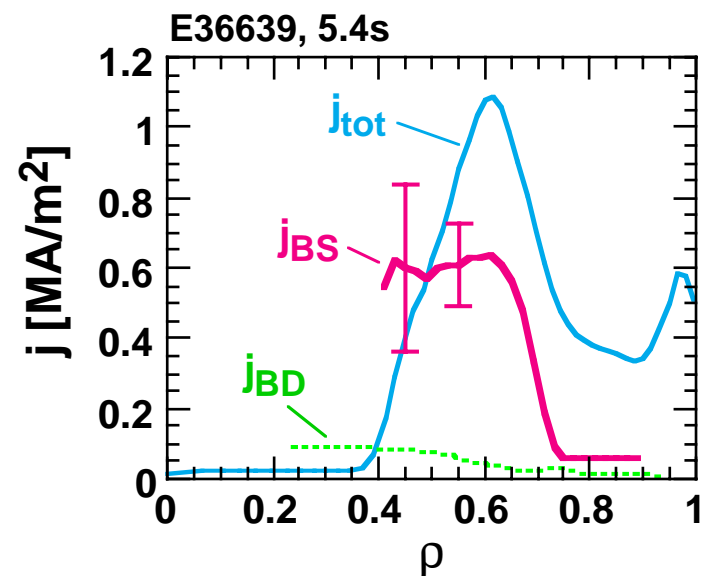
# Bootstrap current is dominant around the current hole

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- While the peak of  $j$  moves inward, a steep gradient in  $j$  (scale length  $\sim 10$  cm or 10% of minor radius) was sustained around the current hole.
- Since the current diffusion time for 10 cm is  $\sim 1.5$  s, the gradient should decrease if all current is inductive one.  $\rightarrow$  non-inductive current



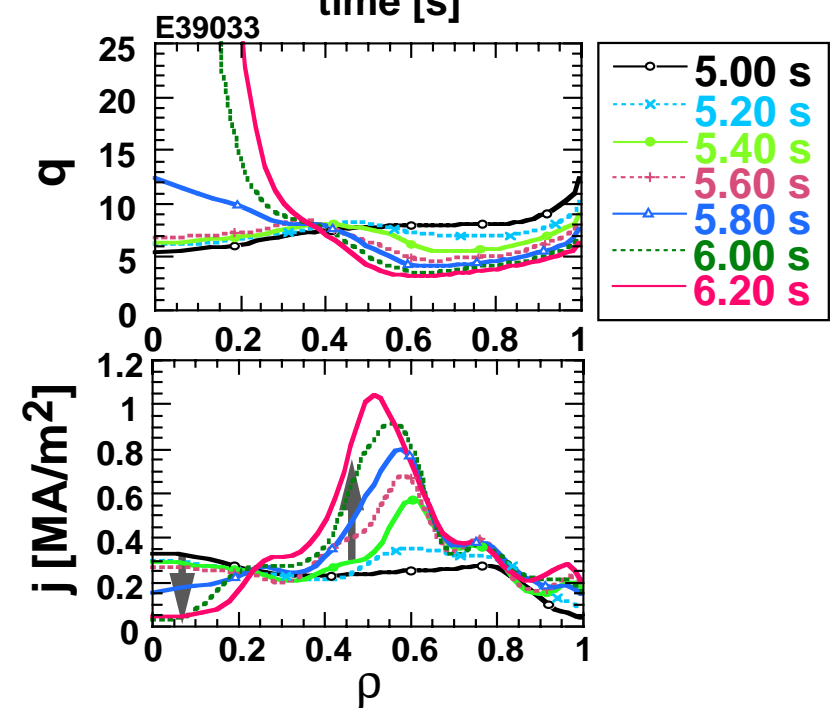
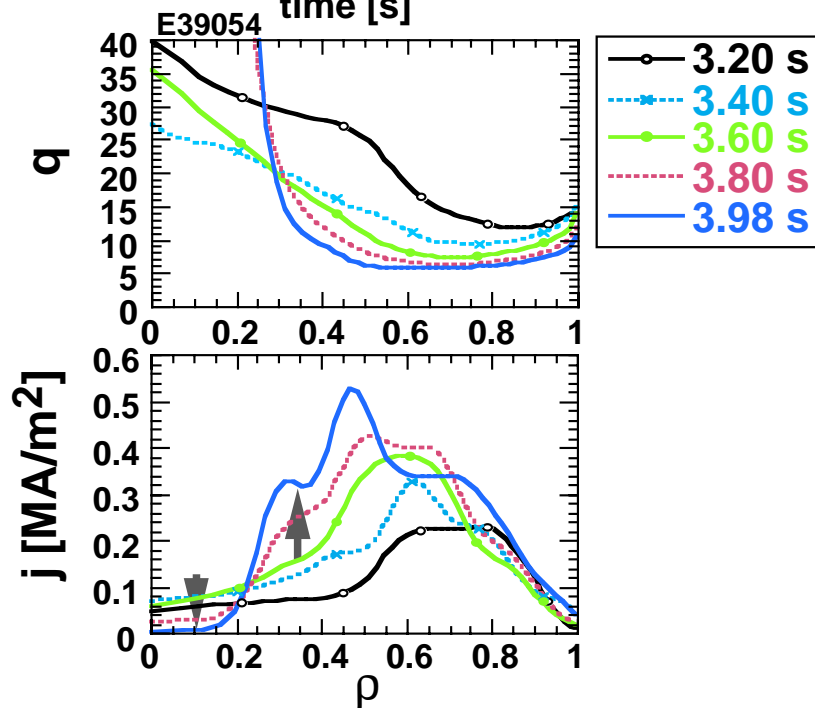
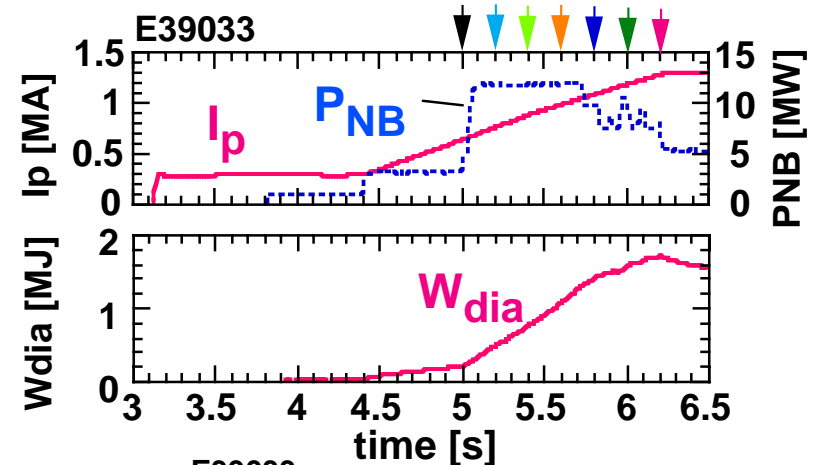
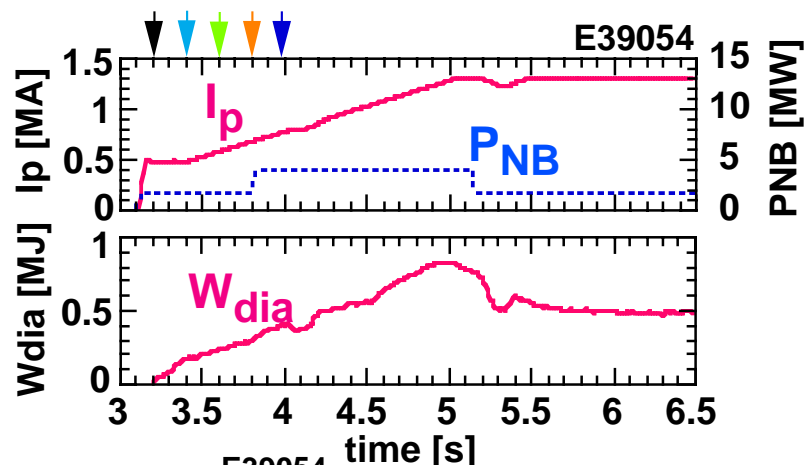
- Since  $B_p$  is small around the current hole,  $j_{BS}$  is large and comparable to the total current, while  $j_{BD}$  is small due to balanced NB injection.



# Formation of current hole (1)

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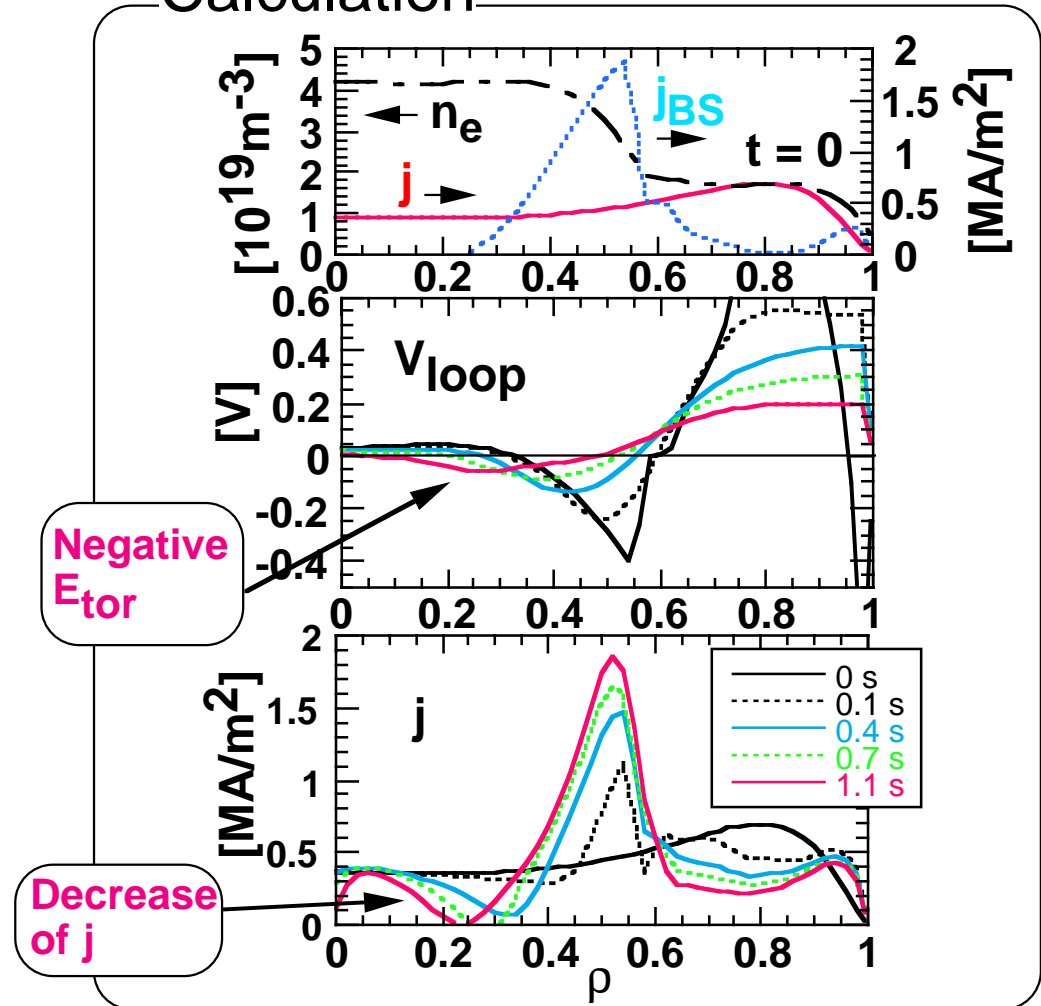
- Central current started to decrease after the growth of off-axis current.
- No counter current drive is expected due to balanced NB injection.



# Formation of current hole (2)

- Increase of off-axis non-inductive current generates (transient) negative  $E_{\text{toroidal}}$ , which penetrates into the central region and reduces  $j_{\text{OH}}$ .
- Since  $j_{\text{NBCD}}$  was small,  $j_{\text{BS}}$  seems to be a cause. (In JET,  $j_{\text{LH}}$  produces current hole.)  
[N. C. Hawkes et al., Phys. Rev. Lett. 87, 115001 (2001)]

## Calculation



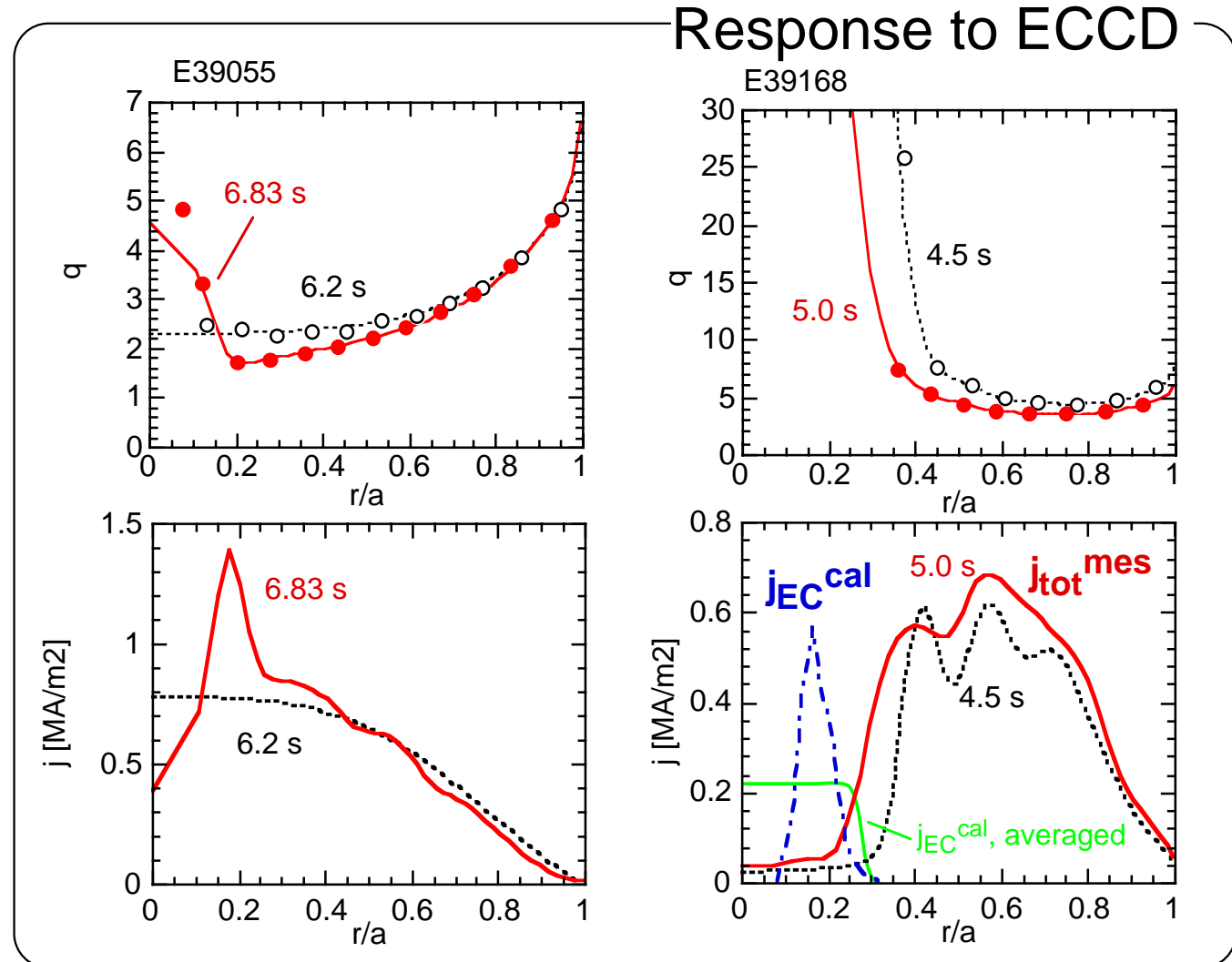
# The central current is clamped at zero ?

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The negative  $j(0)$  may be produced through the transient negative  $E_{\text{tor}}(0)$ , but was not observed.

$j(0)$  seems to be clamped at zero.

No driven-current was observed when EC was injected into the current hole.



The peaked  $j_{\text{EC}}$  profile may not be generated in a flat  $T_e$  region (no confinement of electrons). However, if a uniform  $j_{\text{EC}}$  is generated in the current hole, it should be detected (green curve).

# What happens in the current hole ?

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Observation of **clamping the central current density at zero** for a several seconds seems to exclude

A)  $j_{OH} = \sigma E_{tor} \sim 0$  and  $j_{NI} \sim 0$  ( $j_{NI}$ : non-inductive current)

but suggest existence of **some mechanism to keep  $j(0) \sim 0$**  under following conditions

B)  $\sigma E_{tor} + j_{NI} \sim 0$

( $j_{BS}$  is automatically adjusted to cancel  $j_{OH}$ .)

C)  $\sigma E_{tor} + j_{NI} < 0$  but  $j_{tot} \sim 0$  due to instabilities

(eg. Huysmans et al, PRL 87, 245002 (2001).)

D)  $\sigma E_{tor} + j_{NI} \neq 0$  but  $j_{tot} \sim 0$  once  $j_{tot}$  becomes zero.

(current cannot be driven in a torus without  $B_p$  or  $B_v$ )

To clarify which is the case (B or C or D), analysis of internal voltage or  $E_{tor}(r)$  is important and essential. However, even if  $E_{tor}(r)$  is known, it is not so easy to distinguish B/C/D because we have a large ambiguity in  $j_{BS}$  in a very low  $B_p$  region.

# Summary

- **Equilibrium with nearly zero toroidal current in a central region or "current hole" has been observed to persist stably for several seconds.**
- **The radius of current hole extended up to 40% of plasma minor radius.**
- **A high temperature plasma was confined in the current hole and good confinement was obtained.**
- **The current hole was generated by a negatively induced toroidal electric field through the increase of off-axis bootstrap current.**
- **The current hole seems to be one of causes for formation of ITB shoulder, but other cause also exist.**
- **Clamping at zero current for several seconds and no response to ECCD suggest existence of some mechanism; anomalous resistivity or formation of self-organizing structure (reaching equilibrium limit).**