## **ECRH Superposition on Linear Cylindrical Helicon Plasma in LMD-U**

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The electron cyclotron resonance heating (ECRH) is superimposed on a linear cylindrical helicon plasma in LMD-U to investigate the characteristics of a combined ECRH-helicon plasma. The LMD-U has a vacuum vessel, 3.74 m long and 0.44 m in diameter. At the one end of the device, a conventional helicon source (with a frequency of 7 MHz and rf power of 3kW) is installed for plasma production. At the opposite end of the device, the microwave launcher (with a frequency of 2.45 GHz and the microwave power of 0.8 kW) is installed. The magnetic field configuration is adjusted so that the resonance layer is around the axis position z = 2.3 m. Here, z = 0 shows the position of the boundary of the cylindrical vacuum chamber and the helicon source [1, 2]. The magnetic configuration is different from previous works [1,2] in LMD-U, and thus plasma features are different. Typical time evolution of ion saturation current  $I_{is}$  is shown in Fig. (a). The radial profiles of electron density and electron temperature are measured with the double probe system. The maximum electron density increases from 0.9  $\times 10^{19}$  m<sup>-3</sup> to 1  $\times 10^{19}$  m<sup>-3</sup> and the electron temperature changes from 2.4 eV to 2.5 eV during ECRH injection. Figure (b) shows the typical auto-power spectra of  $I_{is}$ for both helicon discharge phase and ECRH injection phase. A peak at 1 kHz observed in both phases has a m = 1 structure propagating in the ion diamagnetic direction, where m is poloidal

mode number. In the helicon discharge phase, three peaks located near 3.1, 5.3, 6.3 kHz have m = 1, 3, 2 structures propagating in the electron diamagnetic direction. In the ECRH phase, frequencies of these mode structure are changed to 2.7, 4.1, 5.1 kHz, respectively. By the bi-spectrum analysis, the nonlinear coupling between a 1 kHz coherent mode and high frequency broadband components (f > 10 kHz) is observed in the helicon discharge phase. On the other hand, such nonlinear coupling is not observed in ECRH injection phase. This work is partly supported by the Grant-in-Aids for Scientific Research (S) from JSPS (21224014), by collaboration programs of RIAM of Kyushu University and NIFS, and by Research Fellowships of the JSPS for Young Scientists.

[1] K. Terasaka *et al.*, Plasma Fusion Res. **2** (2007) 031.

[2] H. Arakawa *et al.*, Plasma Phys. Control. Fusion **51** (2009) 085001.



Figure (a) time evolution of ion saturation current, (b) the power spectrum in the pure-helicon and ECRH injection phases at r = 4cm, z = 1.885 m.