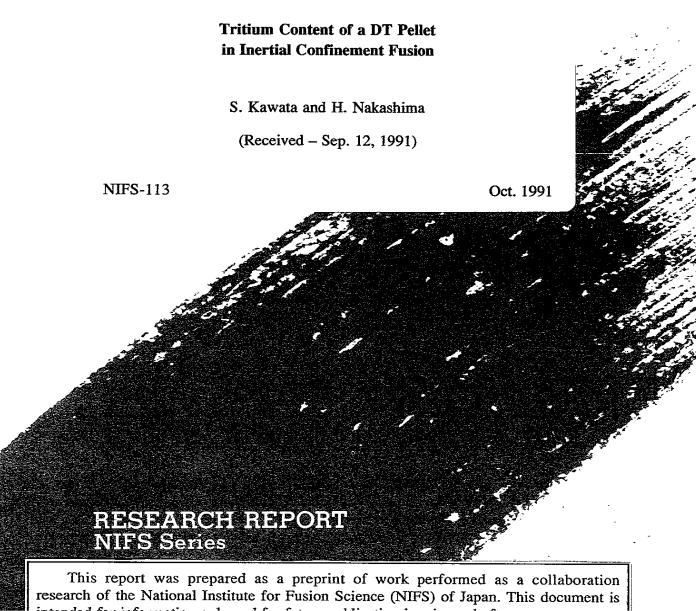
NATIONAL INSTITUTE FOR FUSION SCIENCE



intended for information only and for future publication in a journal after some rearrangements of its contents.

Inquiries about copyright and reproduction should be addressed to the Research Information Center, National Institute for Fusion Science, Nagoya 464-01, Japan.

Tritium Content of a DT Pellet in Inertial Confinement Fusion

S. KAWATA and *H. NAKASHIMA
Department of Electrical Engineering,
Nagaoka University of Technology, Nagaoka 940-21, Japan
*Department of Energy Conversion Engineering,
Kyusyu University, Kasuga, Fukuoka 816, Japan

Abstract

A numerical and an analytical estimations, and a one-dimensional hydrodynamic simulation present that the 30% order of reduction of the tritium content in a DT pellet can bring the sufficient energy output in a DT inertial confinement fusion reactor. In other words, the tritium content can be reduced significantly without a significant reduction in the DT fusion energy output. This new result also means that the tritium inventory can be reduced significantly before and during the reactor operation in the DT inertial confinement fusion. This result comes from the cotribution of a DD reaction to the DT reaction.

Key words: tritium content, inertial confinement fusion, tritium reduction, DT fusion, fusion reaction, pellet study

1.Introduction

In deuterium (D)-tritium(T) inertial confinement fusion (ICF) a pellet is compressed to about a 1000 times as high as the solid density[1] and burned with the peak temperature of $50 \sim 100 \text{keV}$ level. On the other hand, tritium is an unstable radioactive isotope of hydrogen. Therefore we have to breed tritium in a fusion reactor. Because the reaction rate is proportional to the product of the number density of deuterium and tritium, in the DT inertial confinement fusion it has been considered that the tritium content should be equal to that of deuterium in a DT fuel pellet, except for some researches[2] for the advanced fuel[3] in which a density-radius product of ρR is quite large compared with the typical value of $\rho R = 2 \sim 4g/cm^2$ in the DT ICF. However the deuterium-deuterium reaction which produces tritium (see (2)) at the high temperature of $50 \sim 100 \text{keV}$ level, also contributes to the real reaction process.

In this paper, a numerical and an analytical estimation, and a one-dimensional simulation present that a reduction of 30% order of tritium content in a DT fuel pellet can bring the sufficient fusion energy comparable to that in a pellet which contains the equal deuterium and tritium content. In our analyses we include the advanced fusion reactions (DD, DHe^3 , TT and THe^3) in addition to the DT reaction.

2. Reduction of Tritium Content

Let us first present the reactions[4] contributing to the *DT*-based reaction:

$$D + T \longrightarrow H_e^4 + n \tag{1}$$

$$D+D \longrightarrow T+P(50\%)$$

$$\searrow H_e^3 + n(50\%)$$
(2)

$$D + H_e^3 \longrightarrow H_e^4 + p \tag{3}$$

$$T + T \longrightarrow H_e^4 + 2n \tag{4}$$

$$T + H_e^3 \longrightarrow H_e^4 + D(43\%) \tag{5}$$

Here H_e^4 is Helium, H_e^3 an isotope of Helium, p a proton and n a neutron. The content of important elements in a fuel pellet is estimated by the following reaction equations:

$$dn_T/dt = -n_D n_T < \sigma v >_{DT} + n_D^2 < \sigma v >_{DD} / 4 -n_T^2 < \sigma v >_{TT} - n_T n_{H_2^2} < \sigma v >_{TH_2^2} - n_T / \tau$$
(6)

$$dn_D/dt = -n_D n_T < \sigma v >_{DT} - n_D n_{H_e^3} < \sigma v >_{DH_e^3} -n_D^2 < \sigma v >_{DD} +0.43 \times n_T n_{H_e^3} < \sigma v >_{TH_e^3} -n_D/\tau$$
(7)

$$dn_{H_{\epsilon}^{3}}/dt = n_{D}^{2} < \sigma v >_{DD} /4 - n_{D}n_{H_{\epsilon}^{3}} < \sigma v >_{DH_{\epsilon}^{3}} -n_{T}n_{H_{\epsilon}^{3}} < \sigma v >_{TH_{\epsilon}^{3}} -n_{H_{\epsilon}^{3}} /\tau$$
(8)

In the above equations $\langle \sigma v \rangle$, which depends on the temperature only, shows the reaction rate for each reaction and t the time. The relation of $\nabla \cdot (n_j v) \simeq n_j / \tau$ is also used, where j=D, T or H_e^3 , and τ is the confinement time which is $R/(3C_s)$. Here C_s is the sound speed and R is the radius of a spherical fuel pellet.

In the numerical estimation for the reduction of tritium content we compute the time sequence of n_T , n_D and n_{He}^3 by solving Eqs.(6)-(8) during the burning phase. Here we assume that the initial fuel density is a 1000 times of the solid density (n_s) . As a parameter the initial ρR is changed from 2 to 4 g/cm² which is the typical parameter range in order to burn the fuel in the DT inertial confinement fusion. Here ρ is the mass density. The temperature is estimated by the following equation[5] with the initial temperature of 10keV:

$$3n\frac{dT}{dt} = KT^{7/2}/R^2 - R_a n^2 T^{1/2} - nT/\tau + E_{s-h}$$
(9)

In the right hand side of the above equation the first term represents the electron heat conduction loss, the second the radiation (Bremsstrahlung) loss, the third the fuel-expansion loss and the last the self-heating by the fusion products except neutron. Here K and R_a are constant [6,7]. In our model being employed in this paper radiation and neutrons escape freely from the pellet. This is based on a fact that the Rosseland mean free path [7] l_R is much larger than the pellet radius: For a typical parameter set of T=10keV and $n=1000n_s$, $l_R=39.9\text{cm}\gg R$. Therefore this assumption is valid for a pellet without a heavy tamper. The contribution of neutron is also negligible bacause of the low ρR . Charged particles of fusion products deposit their whole energy given by the reactions instanteneously; Among these fusion products alpha particle being produced by the DT reaction is a dominant candidate for heating the fuel. In our parameter range the contribution of DT-alpha particle to heating the fuel is more than 90%. The most important reaction contributing to heating the fuel is (1) and the next is (2). Among the charged fusion products in (1) and (2) the most energetic one is proton in (2). However we neglect its transport effect on heating the fuel, because of the low contribution ratio compared with that of the DT-alpha particle. The mean free path l_{α} of DT-alpha particle is described by $0.04T(\text{keV})/\rho$ approximately[5]. For our parameter range, $l_{\alpha} < R$. Our assumption being employed is based on this fact. We also employ a one-temperature model; As is well known, ions are heated first by the fusion products during the burning phase[8]. The ion-electron energy-relaxation time τ_{ie} is estimated as follows: For 50keV of the ion temperature and 25keV of the electron temperature, τ_{ie} =11.0psec. This value of τ_{ie} is small compared with the whole burning time(see Fig.1). However it is not very small. In order to check the validity of this one-temperature model and to include the radial dependency of physical quantities precisely, we also perform the two-temperature one-dimensional hydrodynamic simulation at the end of this chapter.

Figure 1 presents the time sequence of the temperature TkeV and the fusion

output energy $E_f \text{J/cm}^3$ for $\rho R = 2.0 \text{ g/cm}^2$. In this particular case the tritium content is 42.5% in a fuel, that means 15% reduction of the tritium content. Figure 2 shows the fusion energy output E_f versus the reduction δ of tritium content, which is defined as follows:

$$n_{D0} = n_0/2 + \delta, n_{T0} = n_0/2 - \delta \tag{10}$$

Here n with a suffix of 0 represents the number density at time = 0 and n_0 means the initial total number density. In Fig.2 each pellet radius is slightly different from that of each other in order to keep the specified ρR exactly, when the tritium content is changed. Therefore the decrease in the tritium content introduces the decrease in ρ and the slight increase in the pellet radius. This increase in the pellet radius contributes to make the confinement time τ longer and then to make the energy output increase. In order to check this effect, we also performed the similar parameter study in which each pellet radius is fixed to that of the pellet containing the equal deuterium and tritium content in each specified ρR . The results are shown in Fig.3. Figure 3 also shows that the reduction of tritium content can reach 20 \sim 30% order. These figures present that the reduction of 30% order of tritium content in a fuel pellet can produce the sufficient fusion energy output comparable to that in a pellet which contains the equal deuterium and tritium content. The reduction of tritium can reach 30% without a significant decrease of the fusion energy output as shown in Figs.2 and 3. This fact comes from the contribution of the deuterium-deuterium reaction, which produces tritium, to the deuterium-tritium reaction. This result means that the fusion energy output is rather insensitive to the tritium content in the parameter range in the DT ICF

We also perform an analytical estimation for the optimum content of tritium. In order to do so, we use the simplified reaction equations with Eqs.(10):

$$dn_T/dt \simeq -n_D n_T < \sigma v >_{DT} + n_D^2 < \sigma v >_{DD} /4$$

$$\simeq n_T^2 (- < \sigma v >_{DT} + < \sigma v >_{DD} /4)$$
(11)

$$dn_D/dt \simeq -n_D n_T < \sigma v >_{DT} - n_D^2 < \sigma v >_{DD}$$

$$\simeq -n_D^2 (< \sigma v >_{DT} + < \sigma v >_{DD})$$
(12)

In this analysis we assume that the temperature is constant and $\delta << n_0$. Then we obtain

$$\delta/(n_0/2) = (5/32)(n_0\tau)^2 < \sigma v >_{DT} < \sigma v >_{DD} / (1 + n_0\tau < \sigma v >_{DT} / 2)^2$$
 (13)

for the optimum content of tritium in a DT ICF pellet from the condition of $dn_D n_T/d\delta$ = 0. This result shows that δ is positive and the reduction of tritium content brings sufficient fusion energy. For an example parameter set of T=100 keV, $n_0=4.5\times 10^{25}/\text{cm}^3$ and $\tau=100 \text{psec}$, $\delta/(n_0/2)=0.014$. This result suppports the result of the numerical analyses being presented above, although a number itself for the optimum δ is rather small compared with the numerical one because of its crude approximation.

In addition to these analyses, we also performed a one-dimensional hydrodynamic simulation [9], in order to confirm our results obtained by the above analyses. In this simulation the two-temperature model is employed with the following assumptions: Based on the above discussions about the assumptions being employed, neutrons and radiation escape freely, and charged particles of fusion products deposit their energy locally and instantaneously. The pellet has no tamper. As an initial condition we employ the constant-pressure model[10] in which the temperature of the inner hot spot is 5 keV, that of the outer part is 1 keV and the density ratio between these two parts is 1:5. In our simulation ρR of the hot spot is 0.5g/cm^2 and the density in the outer part is 1000n_s . The total ρR is 4g/cm^2 . The equation-of-state being employed is the ideal one, because of the high temperature during the burning phase. The pellet radius

is fixed to that of the pellet which contains the equal deuterium and tritium content. The simulation results are presented in Fig.3 by square signs. These results reproduce and confirm well the estimation results presented above.

3.Summary

In this paper, we presented the study for the tritium content in a DT-ICF fuel pellet, including the advanced fusion reactions in addition to the DT reaction. We performed this work numerically and analytically. We found that in the DT ICF a reduction of 30% order of tritium content in a fuel pellet can realize the sufficient fusion energy comparable to that in a pellet which contains the equal deuterium and tritium content. This result is also supported by the one-dimensional hydrodynamic simulation. The purpose of this paper is to point out this result and to point out that accordingly the tritium inventory can be reduced significantly before and during the operation of the fusion reactor.

Acknowledgements

We are grateful to Prof.K.Niu for many useful discussions and to M.Ichioka for his valuable help for the computation. This work is partly supported by the Scientific Research Fund of the Ministry of Education and Culture in Japan, and also partly supported by the cooperation program in the National Institute for Fusion Science in Japan.

References

- [1] For example, Duderstadt, J.J. and Moses, G.A., "Inertial Confinement Fusion", (John Wiley and Sons, New York 1981); Kawata, S. and Niu, K., J.Phys.Soc.Jpn., **53**, 3416(1984).
- [2] Skupsky, S., Nuclear Fusion 18, 843(1978).

- [3] For example, Kawata, S., Takase, H. and Niu, K., J.Phys.Soc.Jpn., 51, 3018(1982);
 Takase, H., Kawata, S. and Niu, K., J.Phys.Soc.Jpn., 52, 3400(1983).
- [4] Book, D.L., "NRL Plasma Formula", (Naval Res. Lab., Washington D.C., 1983).
- [5] Takabe, T., Mima, K., and Nakai, S., Laser and Particle Beams 7, 175(1989).
- [6] Spitzer, L., "Physics of Fully Ionized Gases", (Wiley, New York, 1962).
- [7] Zel'Dovich, Ya.B. and Raizer, Y.P., "Physics of Schock Wave and High Temperature Hydrodynamic Phenomena", (Academic Press, New York, 1966).
- [8] Fraley, G.S., Linnebur, E.J., Mason, R.J., and Morse, R.L., Phys. Fluids 17, 474(1974).
- [9] Shiba, T., Nakashima, T., Yoshioka, Y., and Kanda, Y., Nuclear Fusion 28, 699(1988).
- [10] Meyer-ter-Vehn, J., Nuclear Fusion 22, 56(1982).

Figure Captions

- Fig. 1. The time t sequence of temepature T (a solid line) and fusion output energy E_f (a dotted line) for $\rho R=2g/\text{cm}^2$. In this example case the tritium content is 42.5% in a fuel, that means 15% reduction of the tritium content.
- Fig. 2. The fusion energy output E_f versus the tritium content. In this analysis each pellet radius is changed slightly in order to keep the specified ρR exactly, when the tritium content is changed. About 30% reduction of tritium can still bring the sufficient fusion energy in the DT-ICF parameter range.
- Fig. 3. The fusion energy output E_f versus the tritium content. In this case the pellet radius is fixed to that of the pellet which has the equal deuterium-tritium content in each specified ρR . About 30% reduction of tritium can also bring the sufficient fusion energy. Squre signs show the results for $\rho R=4g/cm^2$ by a one-dimensional hydrodynamic simulation. Figure 3 presents that the results by the numerical estimation and by the one-dimensional simulation agree well with each other.

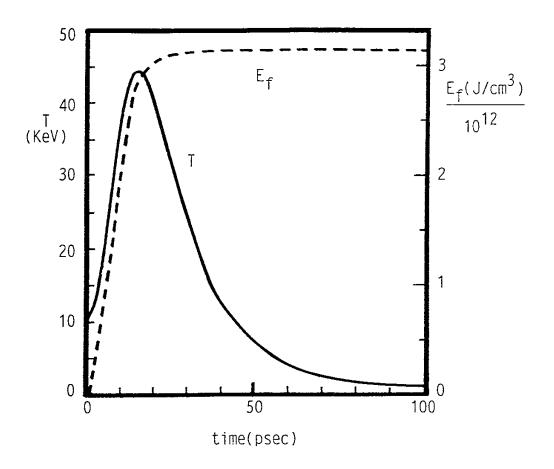


Fig.1

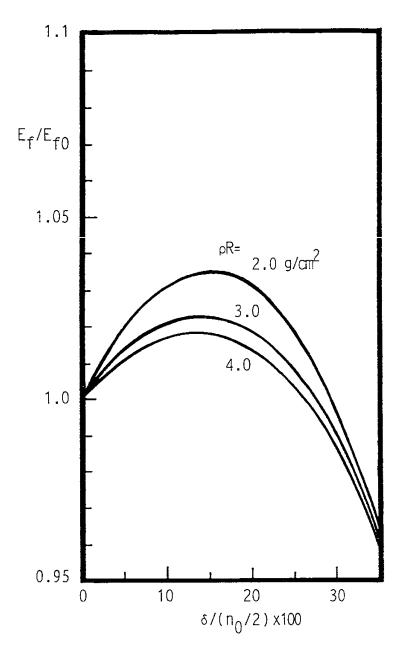


Fig.2

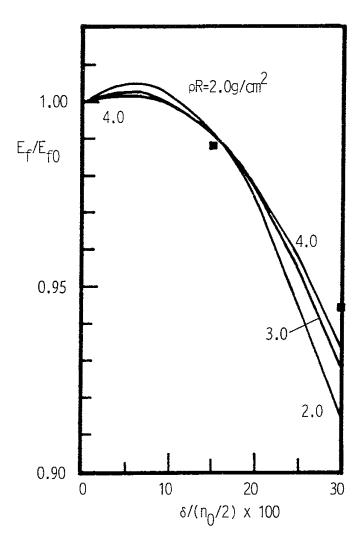


Fig.3

Recent Issues of NIFS Series

NIFS-52	T.Sato, R.Horiuchi, K. Watanabe, T. Hayashi and K.Kusano, Self-Organizing Magnetohydrodynamic Plasma; Sep. 1990
NIFS-53	M.Okamoto and N.Nakajima, <i>Bootstrap Currents in Stellarators and Tokamaks</i> ; Sep. 1990
NIFS-54	K.Itoh and SI.Itoh, Peaked-Density Profile Mode and Improved Confinement in Helical Systems; Oct. 1990
NIFS-55	Y.Ueda, T.Enomoto and H.B.Stewart, Chaotic Transients and Fractal Structures Governing Coupled Swing Dynamics; Oct. 1990
NIFS-56	H.B.Stewart and Y.Ueda, Catastrophes with Indeterminate Outcome; Oct. 1990
NIFS-57	SI.ltoh, H.Maeda and Y.Miura, Improved Modes and the Evaluation of Confinement Improvement; Oct. 1990
NIFS-58	H.Maeda and SI.Itoh, The Significance of Medium- or Small-size Devices in Fusion Research; Oct. 1990
NIFS-59	A.Fukuyama, SI.Itoh, K.Itoh, K.Hamamatsu, V.S.Chan, S.C.Chiu, R.L.Miller and T.Ohkawa, <i>Nonresonant Current Drive by RF Helicity Injection</i> ; Oct. 1990
NIFS-60	K.Ida, H.Yamada, H.Iguchi, S.Hidekuma, H.Sanuki, K.Yamazaki and CHS Group, <i>Electric Field Profile of CHS Heliotron/Torsatron Plasma with Tangential Neutral Beam Injection</i> ; Oct. 1990
NIFS-61	T.Yabe and H.Hoshino, Two- and Three-Dimensional Behavior of Rayleigh-Taylor and Kelvin-Helmholz Instabilities; Oct. 1990
NIFS-62	H.B. Stewart, Application of Fixed Point Theory to Chaotic Attractors of Forced Oscillators; Nov. 1990
NIFS-63	K.Konn., M.Mituhashi, Yoshi H.Ichikawa, Soliton on Thin Vortex Filament; Dec. 1990
NIFS-64	K.Itoh, SI.Itoh and A.Fukuyama, Impact of Improved Confinement on Fusion Research; Dec. 1990
NIFS -65	A.Fukuyama, SI.Itoh and K. Itoh, A Consistency Analysis on the Tokamak Reactor Plasmas; Dec. 1990
NIFS-66	K.Itoh, H. Sanuki, SI. Itoh and K. Tani, Effect of Radial Electric Field on α-Particle Loss in Tokamaks; Dec. 1990
NIFS-67	K.Sato,and F.Miyawaki, Effects of a Nonuniform Open Magnetic Field on the Plasma Presheath; Jan. 1991
NIFS-68	K.Itoh and SI.Itoh, On Relation between Local Transport Coefficient and Global Confinement Scaling Law; Jan. 1991

- NIFS-69 T.Kato, K.Masai, T.Fujimoto,F.Koike, E.Källne, E.S.Marmar and J.E.Rice, He-like Spectra Through Charge Exchange Processes in Tokamak Plasmas; Jan.1991
- NIFS-70 K. Ida, H. Yamada, H. Iguchi, K. Itoh and CHS Group, *Observation of Parallel Viscosity in the CHS Heliotron/Torsatron*; Jan.1991
- NIFS-71 H. Kaneko, Spectral Analysis of the Heliotron Field with the Toroidal Harmonic Function in a Study of the Structure of Built-in Divertor; Jan. 1991
- NIFS-72 S. -l. Itoh, H. Sanuki and K. Itoh, Effect of Electric Field Inhomogeneities on Drift Wave Instabilities and Anomalous Transport; Jan. 1991
- NIFS-73 Y.Nomura, Yoshi.H.lchikawa and W.Horton, Stabilities of Regular Motion in the Relativistic Standard Map; Feb. 1991
- NIFS-74 T.Yamagishi, Electrostatic Drift Mode in Toroidal Plasma with Minority Energetic Particles, Feb. 1991
- NIFS-75 T.Yamagishi, Effect of Energetic Particle Distribution on Bounce Resonance Excitation of the Ideal Ballooning Mode, Feb. 1991
- NIFS-76 T.Hayashi, A.Tadei, N.Ohyabu and T.Sato, Suppression of Magnetic Surface Breading by Simple Extra Coils in Finite Beta Equilibrium of Helical System; Feb. 1991
- NIFS-77 N. Ohyabu, *High Temperature Divertor Plasma Operation*; Feb. 1991
- NIFS-78 K.Kusano, T. Tamano and T. Sato, Simulation Study of Toroidal Phase-Locking Mechanism in Reversed-Field Pinch Plasma; Feb. 1991
- NIFS-79 K. Nagasaki, K. Itoh and S. -I. Itoh, Model of Divertor Biasing and Control of Scrape-off Layer and Divertor Plasmas; Feb. 1991
- NIFS-80 K. Nagasaki and K. Itoh, Decay Process of a Magnetic Island by Forced Reconnection: Mar. 1991
- NIFS-81 K. Takahata, N. Yanagi, T. Mito, J. Yamamoto, O.Motojima and LHDDesign Group, K. Nakamoto, S. Mizukami, K. Kitamura, Y. Wachi, H. Shinohara, K. Yamamoto, M. Shibui, T. Uchida and K. Nakayama, Design and Fabrication of Forced-Flow Coils as R&D Program for Large Helical Device; Mar. 1991
- NIFS-82 T. Aoki and T. Yabe, Multi-dimensional Cubic Interpolation for ICF Hydrodynamics Simulation; Apr. 1991
- NIFS-83 K. Ida, S.-I. Itoh, K. Itoh, S. Hidekuma, Y. Miura, H. Kawashima, M. Mori, T. Matsuda, N. Suzuki, H. Tamai, T.Yamauchi and JFT-2M Group, Density Peaking in the JFT-2M Tokamak Plasma with Counter Neutral Beam Injection; May 1991

- NIFS-84 A. liyoshi, Development of the Stellarator/Heliotron Research; May 1991
- NIFS-85 Y. Okabe, M. Sasao, H. Yamaoka, M. Wada and J. Fujita, Dependence of Au⁻ Production upon the Target Work Function in a Plasma-Sputter-Type Negative Ion Source; May 1991
- NIFS-86 N. Nakajima and M. Okamoto, Geometrical Effects of the Magnetic Field on the Neoclassical Flow, Current and Rotation in General Toroidal Systems; May 1991
- NIFS-87 S. -I. Itoh, K. Itoh, A. Fukuyama, Y. Miura and JFT-2M Group, ELMy-H mode as Limit Cycle and Chaotic Oscillations in Tokamak Plasmas; May 1991
- NIFS-88 N.Matsunami and K.Kitoh, *High Resolution Spectroscopy of H*+ *Energy Loss in Thin Carbon Film*; May 1991
- NIFS-89 H. Sugama, N. Nakajima and M.Wakatani, Nonlinear Behavior of Multiple-Helicity Resistive Interchange Modes near Marginally Stable States; May 1991
- NIFS-90 H. Hojo and T.Hatori, Radial Transport Induced by Rotating RF Fields and Breakdown of Intrinsic Ambipolarity in a Magnetic Mirror; May 1991
- NIFS-91 M. Tanaka, S. Murakami, H. Takamaru and T.Sato, Macroscale Implicit, Electromagnetic Particle Simulation of Inhomogeneous and Magnetized Plasmas in Multi-Dimensions; May 1991
- NIFS-92 S.-I. Itoh, H-mode Physics, -Experimental Observations and Model Theories-, Lecture Notes, Spring College on Plasma Physics, May 27 June 21 1991 at International Centre for Theoretical Physics (IAEA UNESCO) Trieste, Italy; Jun. 1991
- NIFS-93 Y. Miura, K. Itoh, S. I. Itoh, T. Takizuka, H. Tamai, T. Matsuda, N. Suzuki, M. Mori, H. Maeda and O. Kardaun, Geometric Dependence of the Scaling Law on the Energy Confinement Time in H-mode Discharges; Jun. 1991
- NIFS-94 H. Sanuki, K. Itoh, K. Ida and S. I. Itoh, On Radial Electric Field Structure in CHS Torsatron / Heliotron; Jun. 1991
- NIFS-95 K. Itoh, H. Sanuki and S. I. Itoh, *Influence of Fast Ion Loss on Radial Electric Field in Wendelstein VII-A Stellarator;* Jun. 1991
- NIFS-96 S. I. Itoh, K. Itoh, A. Fukuyama, *ELMy-H mode as Limit Cycle and Chaotic Oscillations in Tokamak Plasmas*; Jun. 1991
- NIFS-97 K. Itoh, S. I. Itoh, H. Sanuki, A. Fukuyama, An H-mode-Like Bifurcation in Core Plasma of Stellarators; Jun. 1991
- NIFS-98 H. Hojo, T. Watanabe, M. Inutake, M. Ichimura and S. Miyoshi, Axial Pressure Profile Effects on Flute Interchange Stability in the Tandem Mirror GAMMA 10; Jun. 1991

- NIFS-99 A. Usadi, A. Kageyama, K. Watanabe and T. Sato, A Global Simulation of the Magnetosphere with a Long Tail: Southward and Northward IMF; Jun. 1991
- NIFS-100 H. Hojo, T. Ogawa and M. Kono, Fluid Description of Ponderomotive Force Compatible with the Kinetic One in a Warm Plasma; July 1991
- NIFS-101 H. Momota, A. Ishida, Y. Kohzaki, G. H. Miley, S. Ohi, M. Ohnishi K. Yoshikawa, K. Sato, L. C. Steinhauer, Y. Tomita and M. Tuszewski Conceptual Design of D-3He FRC Reactor "ARTEMIS"; July 1991
- NIFS-102 N. Nakajima and M. Okamoto, Rotations of Bulk Ions and Impurities in Non-Axisymmetric Toroidal Systems; July 1991
- NIFS-103 A. J. Lichtenberg, K. Itoh, S. I. Itoh and A. Fukuyama, *The Role of Stochasticity in Sawtooth Oscillation*; Aug. 1991
- NIFS-104 K. Yamazaki and T. Amano, *Plasma Transport Simulation Modeling for Helical Confinement Systems*; Aug. 1991
- NIFS-105 T. Sato, T. Hayashi, K. Watanabe, R. Horiuchi, M. Tanaka, N. Sawairi and K. Kusano, *Role of Compressibility on Driven Magnetic Reconnection*; Aug. 1991
- NIFS-106 Qian Wen Jia, Duan Yun Bo, Wang Rong Long and H. Narumi, Electron Impact Excitation of Positive Ions - Partial Wave Approach in Coulomb - Eikonal Approximation; Sep. 1991
- NIFS-107 S. Murakami and T. Sato, Macroscale Particle Simulation of Externally Driven Magnetic Reconnection; Sep. 1991
- NIFS-108 Y. Ogawa, T. Amano, N. Nakajima, Y. Ohyabu, K. Yamazaki, S. P. Hirshman, W. I. van Rij and K. C. Shaing, *Neoclassical Transport Analysis in the Banana Regime on Large Helical Device (LHD) with the DKES Code;* Sep. 1991
- NIFS-109 Y. Kondoh, Thought Analysis on Relaxation and General Principle to Find Relaxed State; Sep. 1991
- NIFS-110 H. Yamada, K. Ida, H. Iguchi, K. Hanatani, S. Morita, O. Kaneko, H. C. Howe, S. P. Hirshman, D. K. Lee, H. Arimoto, M. Hosokawa, H. Idei, S. Kubo, K. Matsuoka, K. Nishimura, S. Okamura, Y. Takeiri, Y. Takita and C. Takahashi, *Shafranov Shift in Low-Aspect-Ratio Heliotron / Torsatron CHS*; Sep 1991
- NIFS-111 R. Horiuchi, M. Uchida and T. Sato, Simulation Study of Stepwise Relaxation in a Spheromak Plasma; Oct. 1991
- NIFS-112 M. Sasao, Y. Okabe, A. Fujisawa, H. Iguchi, J. Fujita, H. Yamaoka and M. Wada, Development of Negative Heavy Ion Sources for Plasma Potential Measurement; Oct. 1991