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Electron Impact Excitation Cross Sections and Effective Collision Strengths of N Atom and N-Like Ions — A Review of Available Data and Recommendations —

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Electron impact excitation cross sections and effective collision strengths
of N atom and N-like ions
- A review of available data and recommendations -

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Abstract

The cross sections and/or the effective collision strengths for excitation of N atom and N-like ions by electron impact are reviewed. Main sources for data are listed and the recommended data sources are selected for each ion. The comparison of the cross sections is shown in Graphs. The fit parameters are given for the recommended effective collision strengths in a table; the effective collision strengths calculated therefrom are shown graphically.

Key words

excitation cross section, electron impact, N atom, N-like ions
recommended data

Contents

Introduction

Explanation of Tables

Explanation of Graphs

Tables

- I. Available data sources
- II. Sources for recommended data
- III. Fit parameters for recommended data

Graphs

- I. Comparison of the cross sections
- II. Recommended effective collision strengths

Introduction

The cross sections and/or the effective collision strengths for excitation of N-like ions by electron impact are reviewed. Main sources for data are listed in Table I along the published year for each atom and the recommended data sources are selected for each ion in Table II. The comparison of the cross sections is shown in Graph I. In order to obtain the effective collision strength, the energy dependence of the collision strength in a wide energy region needs to be known. Since most of the results are given as collision strengths, it is necessary the collision strength is first fitted to an analytical function before the effective collision strengths are calculated. The data points are not sufficient to derive the effective collision strength for N-like ions. Then we give here the only original effective collision strength for N atom and Fe XX ions as recommended data.

The effective collision strengths for N atoms and Fe XX ions are fitted to an analytical formula assuming that the collision strength Ω is expressed by the formula

$$\Omega = A + B/X + C/X^2 + D/X^3 + E \ln X, \quad (1)$$

where $X = E_0/\Delta E$, E_0 the energy of the incident electron, ΔE the excitation energy. Correspondingly then, the effective collision strength γ is given as

$$\gamma = y \{ (A/y + C) + (D/2)(1 - y) + e^y E_1(y)(B - Cy + Dy^2/2 + E/y) \} \quad (2)$$

where $y = \Delta E/kT_e$, T_e is the electron temperature, and

$$E_1(y) = \int_y^{\infty} \frac{e^{-t}}{t} dt \quad (3)$$

For the dipole allowed transitions, we fit γ with a fixed value $E = 4\omega_i f_{ij}/\Delta E$ for the parameter E in eq.(2) where f_{ij} is the absorption oscillator strength and ΔE the excitation energy in Rydbergs, ω_i the statistical weight of the initial state¹. The fit parameters are given in Table II with the excitation energy and the effective collision strengths γ are plotted in Graph II.

The rate coefficient is derived with γ as

$$C = \frac{8.010 \times 10^{-8}}{\omega T_e(\text{eV})} e^{-\Delta E/kT_e} \gamma \text{ cm}^3 \text{s}^{-1} \quad (4)$$

The first excited configuration of N atoms is $2s^2 2p^2 3s$. The levels with the configuration $2s2p^4$ appear as the first excited states for OII and FIII ions below the energy levels from the $2s^2 2p^2 3s$ configuration. For ions heavier than NaV the first, second and third excited configurations are $2s2p^4$, $2p^5$ and $2s^2 2p^2 3s$. Thus the cross sections for $2s^2 2p^2 3s$ are often given for N atoms whereas $2s2p^4$ and $2p^5$ for heavier ions, besides those among the levels of the ground state configuration.

1.1. N atoms

Doering and Goembel² measured the absolute differential and integral cross section for the $2s^2 2p^3 4S - 2p^2 3s 4P$ transition with electron scattering techniques. Their values are lower by a factor of two than the fluorescence cross section of Stone and Zipf^{3,4} corrected to reflect the recent change in the H₂ dissociative excitation reference cross section. The values by ref.4 and 5 include the contribution of cascade. Spence and Burrow⁵ measured the cross sections for the $2s^2 2p^3 4S - 2p^2 3s 4P$, $2s^2 p^4 4P$ transition by the electron energy loss spectra near threshold. They observed a strong peak near threshold for $2s^2 2p^3 4S - 2p^2 3s 4P$ transition. The only available theoretical data for this transition are by Ormond et al⁶ and their values are smaller than the experimental values.

Most of the theoretical calculations give the data for transitions among the ground state $2s^2 2p^3$ 4S , 2D and 2P . Berrington et al.⁷ calculated cross sections using R-matrix method with 8 states for $n = 2 - 2$ transitions. Ref.5 used the data of ref. 7 for the $2s^2 2p^3$ 4S - $2s2p^4$ 4P transition to obtain the absolute values for the $2s^2 2p^3$ 4S - $2p^2 3s$ 4P transition. The effective collision strengths by R-matrix calculations are given in ref. 8 for the transitions between the $2s^2 2p^3$ 4S , 2D and 2P in the temperature range from 500 to 10^5 K. The data by ref. 8 are fitted analytical formula of eq.(2) and the fitting parameters are given in Table II. They are plotted in Graph II.

1. 2. O II ions

McLaughlin et al.⁹ calculated the cross sections for the $2p^3$ 4S , 2D , 2P - $2p^2 3s$ 4P , 2P transitions using both a 9 state and a 34 state R-matrix approximation. The collision strengths exhibit strong resonance structure in the energy region between the $2s^2 2p^2 3s$ 2P and $2s2p^4$ 2P thresholds. The importance of the resonances will decrease as the nuclear charge increases. Itikawa et al¹⁰ gave the recommended data and recommended values are fitted to an analytical formula for carbon and oxygen ions. They included the calculations of Henry¹¹, Ho and Henry¹², Pradhan^{13,14}. Ref.12 calculated the collision strengths from the ground state to the $2s2p^4$ 4P and $2s^2 2p^2 3s$ 4P excited states in a two state close coupling approximation. Ref.9 obtains the higher values than ref.12 near the threshold, although the two results at 26 eV are almost similar. Effective collision strengths are obtained using the close coupling approximation by ref.14 for the fine structure transitions within the ground configuration. The results in ref. 13 and 14 are in good agreement with those of Martins et al.¹⁵ for energies above the 2P threshold but about 10% higher at the 2D threshold¹³. The results given by Saraph et al¹⁶ agree within 10 % with those by ref.13.

1.3. F III

There are few results for F III ions. ref.16 gave numerical calculations for the first four ions for the configurations

$2s^2 2p^q$ ($q = 1 - 5$) in the approximation of neglecting coupling to configurations other than $2s^2 2p^q$. Martins et al¹⁵ gave the collision strengths among fine structure levels including resonances converging onto the 2P level by similar method as ref.16. We recommend data by ref..15 and 16, since their data for OII ions are in agreement within 10 % comparing to the recent calculations as shown in the previous section, .

1.5 .Ne IV

Giles¹⁷ presented the rate coefficients for transitions between fine structure levels in the ground configurations. The inclusion of the next excited configuration in the close coupling expansion gives rise to near threshold resonances comparing to those of ref.16. Bhatia and Kastner¹⁸ calculated the collision strength between the 72 levels of the six configurations $2s^2 2p^3$ - $2s^2 2p^3$, $2s2p^4$, $2p^5$, $2p^2 3s$, $2p^2 3p$, $2p^2 3d$. The numerical values of the collision strength for each transitions are given for only one incident electron energy, 7 Ryd. It is required to give the energy dependence of the collision strengths.

1.5. Na V

Same as F III ions.

1.6.Mg VI

Bhatia and Mason¹⁹ calculated collision strengths including the $2s^2 2p^3$, $2s2p^4$ and $2p^5$ configurations for Mg VI, Si VIII, S X, Ar XII and Ca XIV using a computer package developed at University College, London. The data for the configurations $2s^2 2p^3$, $2s2p^4$ are given at 10, 15 and 20 Ryd. incident electron energies and the results connected with the $2p^5$ term and within the $2s2p^4$ term are available on request. According to their paper, their results are in good agreement with the previous calculation by Davis et al²⁰ .

1.7. Si VIII

The DW results by ref.19, 20 and 21 are available.

1.8. S X

The data by Ref.19 is the only available source.

1.9 .Ar XII

Bhatia et al²² calculated the collision strengths including 72 levels of $2s^2 2p^3$, $2s2p^4$, $2p^5$, $2s^2 2p^3 3s$, $2s^2 2p^2 3p$ and $2s^2 2p^2 3d$ for Ar XII, Ti XVI, Fe XX, Zn XXIV and Kr XXX ions. The values of the collision strengths for one impact electron energy point are given for a limited number of transitions for which the radiative transition probabilities exceed a certain value. The collision strengths for spin exchange transitions are not given.

1.10. Ca XIV

Same as S X.

1.11. Ti XVI

Bhatia et al²³ presented the calculations of the collision strength for the Li- through F - like titanium ions. For Ti XVI ions, the collision strengths at 15, 20 and 25 Ryd. are given for all the transitions among $2s^2 2p^3$, $2s2p^4$, $2p^5$ fine structure levels.

1.12.Mn XIX

Bhatia²⁴ calculated the collision strengths for the configurations $2s^2 2p^k$, $2s2p^{k+1}$ and $2s2p^{k+2}$ of Mn ions of the Li - through F - like sequences as calculated for Ti ions in ref.23. For N-like Mn XIX ions, the collision strengths for the impact electron energies of 20, 40 and 60 Ryd. are given for the transitions among 15 levels with the configurations of $2s^2 2p^3$, $2s2p^4$, $2p^5$.

1.13.Fe XX

Ref.25 presented the collision strengths of Fe XX ions for the fine structure transitions among $2s^2 2p^3$ and $2s2p^4$ including the configurations $2s^2 2p^3$, $2s2p^4$ and $2p^5$ in the calculation. Mason and Bhatia²⁶ extended these calculations to include transitions from $2s^2 2p^3$ up to $2s^2 2p^2 3s$ and

$2s^2 2p^2 3d$. Kingston and Lennon²⁷ presented polynomial fit coefficients for recommended rate coefficient of iron ions based on the data of Robb(1980)²⁸ and Bhatia. Sometimes their rate coefficients derived fit parameters give wrong values due to the mis-printing of the fit coefficients. Bhatia et al²² calculated new data including 72 levels. The values of radiative decay and the collision strength by ref.26 and ref.22 are almost similar. But some differences are found due to the difference of the order of the energy levels; e.g. the order of the energy levels of $2p^2 3d\ ^2P_{1/2}$ and $2p^2 3d\ ^4P_{1/2}$. This discrepancy is due to the naming of the levels. The effective collision strengths²⁸ based on ref.25 are fitted to analytical formula and plotted in Graph II. The cross sections are not plotted for all the data but only for comparison in Graph I.

1.14 Summary

For the N atom and O II ions there are few results involving transitions to $n = 3$ levels. For ions from Ne IV to Kr XXX, DW calculations have been carried out for $n = 2$ (Mg VI, Si VIII, S X, Ca XIV) and $n = 2$ and 3 (Ne IV, Ar XII, Ti XVI, Mn XIX, Fe XX, Zn XXIV, Kr XXX) for the transitions among fine structure levels.

As the next step, it is necessary to calculate the effective collision strengths from the recommended collision strengths. DW calculations by Bhatia et al. give the collision strengths for no more than three incident electron energies. It is required to calculate the cross sections for more several energy points in order to know the energy dependence.

Since there are so many transitions among fine structure levels, it will be useful to consider a standard database for excitation processes. The problems include the format how to store the configuration of the levels, identification of the transition, fitting formulae and parameters, transition energies etc.

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Explanation of Tables

Table I. Available data sources.

Author(year)	The author of the reference and the published year
Transition	Transition from the initial(left) to the final(wright) state.
(fine str.)	transitions among fine structure levels
Data	
σ	Cross section
Ω	Collision strength
γ	Effective collision strength
C	Rate coefficient
(par. fit)	Fit parameters are given
Energy(Temp.)	The energy or temperature range spanned by the data
(thre.)	Threshold energy
Method	
R-matrix	R-matrix calculation
CC	Close-coupling calculation
DW	Distorted wave calculation
beam	Experiment using electron beam
(electron)	Electron energy loss measurement
(optical)	Optical line emission measurement
Evaluation	Data are evaluated

Table II Sources and energy(temperature) ranges for recommended data.

Table III Fit parameters for recommended effective collision strengths.

Fit parameters are tabulated for each excitation process of N atom⁷ and Fe XX²⁹. The notation, $2s^2 2p^3 4S^0 - 2p^2 2p^3 2D^0$, for example, means the excitation from the lower state to the upper state.

ΔE	Excitation energy in eV
A,B,C,D,E	Fit parameters for the effective collision strengths γ in eq.(2).

Explanation of Graphs

Graph I Comparison of the cross sections.

The excitation cross sections (cm^2) are plotted as a function of energy (eV) for each transitions. A letter "T" or "E" (after the authors and the year of publication in the legend) indicates a theoretical or an experimental paper, respectively. Since there are many transitions among fine structure levels, only part of them are shown.

Graph II The recommended effective collision strengths.

The effective collision strengths are plotted as a function of temperature T_e in Kelvin degree for N atom and Fe XX ions. Their fit parameters are listed in Table III. The symbols and solid lines indicate the original points from ref. 4 and ref. 29, and calculated values from the fit parameters in Table III, respectively.

References for Table I and Graph I

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Table I. Available Data Sources

Author(year)	Transition	Data, Energy(Temp.) range	Method
N atom			
Smith et al(1967)	among $2s^2 2p^3 4S, 2D, 2P$	σ (threshold - 4Ryd.)	CC
Henry and Williams(1968)	among $2s^2 2p^3 4S, 2D, 2P$	γ (500 - 10^5 K, par. fit)	CC
Henry et al(1969)	among $2s^2 2p^3 4S, 2D, 2P$	σ (thre. - 50eV)	CC
Stone and Zipf(1971)	$2s^2 2p^3 4S - 2p^2 3s 4P$	σ (threshold - 15eV)	beam(optical)
Stone and Zipf(1973)	$2s^2 2p^3 4S - 2s2p^4 (4P)$, $2s^2 2p^3 - 2p^2 3s, 3p$	σ (17 - 32eV)	beam(optical)
Ganas(1973)	$2p - ns$ ($n=7$), $2p-np$ ($n=4$), $2p-nd$ ($n=5$)	σ (thre.-1000eV)	Born
Ormond et al(1973)	among $2s^2 2p^3 4S, 2D, 2P$ $2s^2 2p^3 - 2p^2 3s 4P$	Ω (0.3-1.4 Ryd.) Ω (1.0 - 1.4 Ryd.)	CC
Davis et al(1975)	$2s^2 2p^3 - 2p^2 3s$	σ (1.14 - 6.08 Ryd.)	DW
Thomas and Nesbet(1975)	among $2s^2 2p^3 4S, 2D, 2P$	σ (3.4 - 11.0 eV)	matrix variational
Le Dourneuf(1975)	$2s^2 2p^3 2D - 2P$	σ	R-matrix
Berrington et al(1975)	among $2s^2 2p^3 4S, 2D, 2P$ $2s^2 2p^3 - 2s2p^4 4P, 2D, 2S, 2P$	σ (thre. - 2.5 Ryd.)	R-matrix
Spence and Burrow(1980)	$2s^2 2p^3 - 2s^2 2p^3 4S, 2D, 2P$	σ (10.3-11.1 eV)	beam(electron)
Berrington and Burke(1981)	among $2s^2 2p^3 4S, 2D, 2P$	γ (500 - 10^5 K, par. fit)	R-matrix
Doering and Goemel(1991)	$2s^2 2p^3 4S - 2p^2 3s 4P$	σ (30-100 eV)	beam(electron)
O II ions			
Henry et al(1969)	among $2s^2 2p^3 4S, 2D, 2P$	σ (thre. - 50eV)	CC
Saraph et al(1969)	among $2s^2 2p^3 4S, 2D, 2P$ $2D - 2D, 2P$ (fine str.)	Ω (thre.-.91Ryd.)	exact resonance, DW
Martins et al(1969)	among $2s^2 2p^3$ (fine str.)	Ω (0.3 - 0.47 Ryd.)	exact resonance
Ormond et al(1973)	$2s^2 2p^3 4S - 2s2p^4 4P$	Ω (1.16 - 2.6 Ryd.)	CC
Davis et al(1975)	$2s^2 2p^3 - 2s2p^4, 2p^2 3s$	σ (1.1-13.5 Ryd.)	DW
Pradhan(1976a)	among $2s^2 2p^3 4S, 2D, 2P$ $2D - 2D$ (fine structures)	Ω (thre.)	CC
Pradhan(1976b)	among $2s^2 2p^3 4S, 2D, 2P$ $2D - 2D, 2P - 2P$ (fine str.)	Ω (thre.), γ (5×10^3 - 2×10^4 K) γ (5×10^3 - 2×10^4 K)	CC
Ho and Henry(1983)	$2s^2 2p^3 4S -$ $2s2p^4 4P, 2p^2 3s 4P$	Ω, γ (par. fit, $< 10^6$ K)	CC

Table I. Available Data Sources (continued)

Author(year)	Transition	Data, Energy(Temp.) range	Method
Itikawa et al(1985)	among $2s^2 2p^3$ 4S , 2D , 2P $2s^2 2p^3$ 4S - $2s2p^4$ 4P , $2p^2 3s$ 4P	γ (0.3 - 100eV, par.fit)	Evaluation
McLaughlin et al(1987)	$2p^3$ 4S , 2D , 2P - $2p^2 3s$ 4P , 2P	Ω (thre. -2 Ryd.)	R-matrix
F III			
Saraph et al(1969)	among $2s^2 2p^3$ 4S , 2D , 2P 2D - 2D , 2P (fine str.)	Ω	exact resonance, DW
Martins et al(1969)	among $2s^2 2p^3$ (fine str.)	Ω (0.1 - 0.22 Ryd.)	exact resonance
Ne IV			
Saraph et al(1969)	among $2s^2 2p^3$ 4S , 2D , 2P 2D - 2D , 2P (fine str.)	Ω	exact resonance, DW
Martins et al(1969)	among $2s^2 2p^3$ (fine str.)	Ω (0.08 - 0.2 Ryd)	exact resonance
Giles(1981)	among 4S , 2D , 2P (fine str.)	γ (6×10^3 - 2×10^4 K)	CC
Bhatia and Kastner(1988)	$2s^2 2p^3$ - $2s^2 2p^3$, $2s2p^4$, - $2p^5$, $2p^2 3s$, $2p^2 3p$, $2p^2 3d$ (fine str.)	Ω (7Ryd.)	DW
Na V			
Saraph et al(1969)	among $2s^2 2p^3$ 4S , 2D , 2P 2D - 2D , 2P (fine str.)	Ω	exact resonance, DW
Martins et al(1969)	among $2s^2 2p^3$ (fine str.)	Ω (0.05 - 1.7 Ryd.)	exact resonance
Mg VI			
Saraph et al(1969)	among $2s^2 2p^3$ 4S , 2D , 2P 2D - 2D , 2P (fine str.)	Ω	exact resonance, DW
Davis et al(1976)	$2s^2 2p^3$ $^4S_{3/2}$ - $2s2p^4$ $^4P_{1/2}$ $2s^2 2p^3$ $^2P_{3/2}$ - $2s2p^4$ $^2P_{1/2}$	C(par. fit)	DW
Bhatia and Mason(1980a)	$2s^2 2p^3$ - $2s^2 2p^3$, $2s2p^4$ (fine str.)	Ω (10, 15, 20 Ryd.)	DW

Table I. Available Data Sources (continued)

Author(year)	Transition	Data, Energy(Temp.) range	Method
Si VIII			
Davis et al(1976)	$2s^2 2p^3 \ ^4S_{3/2} - 2s2p^4 \ ^4P_{1/2}$	C(par. fit)	DW
	$2s^2 2p^3 \ ^2D_{3/2} - 2s2p^4 \ ^2D_{3/2}$		
Davis et al(1977)	$2s^2 2p^3 \ ^4S - 2p^2 3s \ ^4P, 2p^2 3d \ ^4P$	C(par. fit), Ω	DW
Bhatia and Mason(1980a)	$2s^2 2p^3 - 2s^2 2p^3, 2s2p^4$ (fine str.)	$\Omega(10, 15, 20)$ Ryd.	DW
S X.			
Bhatia and Mason(1980a)	$2s^2 2p^3 - 2s^2 2p^3, 2s2p^4$ (fine str.)	$\Omega(15, 20, 30)$ Ryd.	DW
Ar XII			
Bhatia and Mason(1980a)	$2s^2 2p^3 - 2s^2 2p^3, 2s2p^4$ (fine str.)	$\Omega(15, 30, 45)$ Ryd.	DW
Bhatia et al(1989)	among $2s^2 2p^3, 2s2p^4, 2p^5,$ $2p^2 3s, 2p^2 3p, 2p^2 3d$ (fine str.)	$\Omega(32)$ Ryd. DW,CB	
Ca XIV			
Bhatia and Mason(1980a)	$2s^2 2p^3 - 2s^2 2p^3, 2s2p^4$	$\Omega(15, 30, 45)$ Ryd.	DW
Ti XVI			
Bhatia et al(1980c)	among $2s^2 2p^3, 2s2p^4, 2p^5$ (fine str.)	$\Omega(15, 20, 25)$ Ryd.	DW
Bhatia et al(1989)	among $2s^2 2p^3, 2s2p^4, 2p^5,$ $2p^2 3s, 2p^2 3p, 2p^2 3d$ (fine str.)	$\Omega(32)$ Ryd.	DW,CB
Mn XIX			
Bhatia(1982)	among $2s^2 2p^3, 2s2p^4, 2p^5$ (fine str.)	$\Omega(20, 40, 60)$ Ryd.	DW
Fe XX			
Davis et al(1976)	$2s^2 2p^3 - 2s2p^4, 2p^3 - 2p^2 3d,$ $2s2p^4 - 2p^5$	C(par. fit)	DW
Bhatia and Mason(1980b)	$2s^2 2p^3 - 2s2p^4$ (fine str.)	$\Omega(20, 50, 100)$ Ryd.	DW
Mason and Bhatia(1983)	$2s^2 2p^3 - 2p^2 3s, 2p^2 3d$ (fine str.)	$\Omega(80)$ Ryd.	DW
Mann(1983)	$2s^2 2p^3 (^4S) - 2s2p^4 (^4P, ^2D, ^2S, ^2P)$	$\Omega(x=1 - 50, \text{par. fit})$	DW
Kingston and Lennon(1987)	among $2s^2 2p^3, 2s2p^4$ (fine str.)	C(par. fit)	Evaluation
Bhatia et al(1989)	among $2s^2 2p^3, 2s2p^4,$ $2p^5, 2p^2 3s, 2p^2 3p, 2p^2 3d$ (fine str.)	$\Omega(80)$ Ryd.	DW,CB
Zn XXIV			
Bhatia et al(1989)	among $2s^2 2p^3, 2s2p^4,$ $2p^5, 2p^2 3s, 2p^2 3p$ and $2p^2 3d$ (fine str.)	$\Omega(110)$ Ryd.	DW,CB
Kr XXX			
Bhatia et al(1989)	among $2s^2 2p^3, 2s2p^4,$ $2p^5, 2p^2 3s, 2p^2 3p$ and $2p^2 3d$ (fine str.)	$\Omega(160)$ Ryd.	DW,CB

Table II Sources for Recommended Data

Transition	Data, Energy(Temp.) range	References
N atom		
among $2s^2 2p^3$ $^4S, ^2D, ^2P$	γ (500 - 105 K, par. fit)	Berrington and Burk(1981)
$2s^2 2p^3$ $^4S, ^2D, ^2P$ - $2s2p^4$ $^4P, ^2D, ^2S, ^2D$	σ (thre. - 2.5 Ryd..)	Berrington et al(1975)
$2s^2 2p^3$ 4S - $2p^2 3s$ 4P	σ (30 - 100 eV)	Doering and Goembel(1991)
O II		
$2p^3$ ($^4S, ^2D, ^2P$) - $2p^2 3s$ ($^4P, ^2P$)	Ω (thre. - 2 Ryd.)	McLaughlin et al(1987)
$2s^2 2p^3$ 4S - $2p^2 3s$ $^4P, 2s2p^4$ 4P	γ (0.3 - 100eV)	Itikawa et al(1985)
among $2s^2 2p^3$ $^4S, ^2D, ^2P$	γ (0.3 - 100eV)	Itikawa et al(1985)
$^2D - ^2D, ^2P - ^2P$ (fine structures)	$\gamma(5 \times 10^3 - 2 \times 10^4 K)$	Pradhan(1976b)
F III		
among $2s^2 2p^3$ $^4S, ^2D, ^2P$	Ω	Saraph et al(1969)
among $2s^2 2p^3$ $^4S, ^2D, ^2P$ (fine structure)	$\Omega(0.1 - 0.22 \text{ Ryd.})$	Martins et al(1969)
Ne IV		
among $2s^2 2p^3$ $^4S, ^2D, ^2P$ (fine structure)	$\gamma(6 \times 10^3 - 2 \times 10^4 K)$	Giles(1981)
$2s^2 2p^3$ - $2s2p^4, 2p^5, 2p^2 3s, 2p^2 3p, 2p^2 3d$ (fine structure)	$\Omega(7 \text{ Ryd.})$	Bhatia and Kastner(1988)
Na V		
among $2s^2 2p^3$ $^4S, ^2D, ^2P$	Ω	Saraph et al(1969)
among $2s^2 2p^3$ $^4S, ^2D, ^2P$ (fine structure)	$\Omega(0.027 - 0.15 \text{ Ryd.})$	Martins et al(1969)
Mg VI		
$2s^2 2p^3$ - $2s^2 2p^3, 2s2p^4$ (fine structures)	$\Omega(10, 15, 20 \text{ Ryd.})$	Bhatia and Mason(1980a)
Si VIII		
$2s^2 2p^3$ - $2s^2 2p^3, 2s2p^4$ (fine structures)	$\Omega(10, 15, 20 \text{ Ryd.})$	Bhatia and Mason(1980a)
S X		
$2s^2 2p^3$ - $2s^2 2p^3, 2s2p^4$ (fine structures)	$\Omega(15, 20, 30 \text{ Ryd.})$	Bhatia and Mason(1980a)
Ar XII		
$2s^2 2p^3$ - $2s^2 2p^3, 2s2p^4$ (fine structure)	$\Omega(15, 30, 45 \text{ Ryd.})$	Bhatia and Mason(1980a)
among $2s^2 2p^3, 2s2p^4, 2p^5, 2s^2 2p^3 3s,$	$\Omega(32 \text{ Ryd.})$	Bhatia et al(1989)
$2s^2 2p^2 3p$ and $2s^2 2p^2 3d$ (fine structure)		

Table II Sources for Recommended data (continued)

Transition	Data, Energy(Temp.) range	References
Ca XIV		
$2s^2 2p^3 - 2s^2 2p^3, 2s2p^4$ (fine structure)	$\Omega(15, 30, 45$ Ryd.)	Bhatia and Mason(1980a)
Ti XVI		
among $2s^2 2p^3, 2s2p^4, 2p^5$ (fine structure)	$\Omega(15, 30, 45$ Ryd.)	Bhatia et al(1980)
among $2s^2 2p^3, 2s2p^4, 2p^5, 2p^2 3s,$ $2p^2 3p, 2p^2 3d$ (fine structure)	$\Omega(55$ Ryd.)	Bhatia et al(1989)
Mn XIX		
among $2s^2 2p^3, 2s2p^4, 2p^5$ (fine structure)	$\Omega(20, 40, 60$ Ryd.)	Bhatia(1982)
Fe XX		
among $2s^2 2p^3 - 2s2p^4$ (fine structure), including forbidden transitions	$\Omega(20, 50$ and 100 Ryd.)	Bhatia and Mason(1980b)
$2s^2 2p^3 - 2p^2 3s, 2p^2 3d$ (fine structure)	$\Omega(80$ Ryd.)	Mason and Bhatia(1983)
among $2s^2 2p^3, 2s2p^4, 2p^5, 2p^2 3s,$ $2p^2 3p$ and $2p^2 3d$ (fine structure)	$\Omega(80$ Ryd.)	Bhatia et al(1989)

TABLE III. Fit Parameters for the recommended effective collision strength

N		ΔE (eV)	A	B	C	D	E	
$2s^2 2p^3 4S^0$	-	$2s^2 2p^3 2zD^0$	2.38E+00	-2.07E+00	1.04E+01	-1.53E+01	6.94E+00	1.55E+00
$2s^2 2p^3 4S^0$	-	$2s^2 2p^3 2P^0$	3.57E+00	-2.52E+00	8.74E+00	-1.05E+01	4.30E+00	1.16E+00
$2s^2 2p^3 2D^0$	-	$2s^2 2p^3 2P^0$	1.19E+00	1.73E+00	-5.93E+00	8.23E+00	-3.96E+00	1.11E+00
 F e¹⁹⁺								
$2s^2 2p^3 4S_{3/2}$	-	$2s^2 2p^3 2D_{3/2}$	1.74E+01	4.07E-02	2.87E-01	-7.68E-01	5.12E-01	-5.96E-03
$2s^2 2p^3 4S_{3/2}$	-	$2s^2 2p^3 2D_{5/2}$	2.22E+01	5.10E-02	1.49E-01	-3.75E-01	2.42E-01	-7.46E-03
$2s^2 2p^3 4S_{3/2}$	-	$2s^2 2p^3 2P_{1/2}$	3.19E+01	5.03E-03	1.37E-01	-3.16E-01	2.00E-01	-8.17E-04
$2s^2 2p^3 4S_{3/2}$	-	$2s^2 2p^3 2P_{3/2}$	3.97E+01	9.62E-03	8.47E-02	-1.85E-01	1.13E-01	-1.78E-03
$2s^2 2p^3 4S_{3/2}$	-	$2s^2 p^4 4P_{5/2}$	9.29E+01	2.88E-01	-3.91E-02	5.65E-01	-3.94E-01	1.25E-01
$2s^2 2p^3 4S_{3/2}$	-	$2s^2 p^4 4P_{3/2}$	1.01E+02	1.94E-01	2.50E-02	2.46E-01	-1.74E-01	8.94E-02
$2s^2 2p^3 4S_{3/2}$	-	$2s^2 p^4 4P_{1/2}$	1.04E+02	1.16E-01	-1.17E-01	4.30E-01	-2.75E-01	4.63E-02
$2s^2 2p^3 4S_{3/2}$	-	$2s^2 p^4 2D_{3/2}$	1.30E+02	8.00E-03	6.66E-03	3.69E-03	-4.78E-03	3.78E-03
$2s^2 2p^3 4S_{3/2}$	-	$2s^2 p^4 2D_{5/2}$	1.32E+02	-1.44E-05	1.21E-03	-1.13E-03	4.19E-04	5.71E-05
$2s^2 2p^3 4S_{3/2}$	-	$2s^2 p^4 2S_{1/2}$	1.49E+02	2.25E-03	1.57E-02	-2.14E-02	1.07E-02	1.37E-03
$2s^2 2p^3 4S_{3/2}$	-	$2s^2 p^4 2P_{3/2}$	1.55E+02	9.18E-03	5.00E-02	-6.83E-02	3.44E-02	5.74E-03
$2s^2 2p^3 4S_{3/2}$	-	$2s^2 p^4 2P_{1/2}$	1.67E+02	4.77E-04	4.43E-03	-2.48E-03	1.36E-04	6.30E-05
$2s^2 2p^3 2D_{3/2}$	-	$2s^2 2p^3 2D_{5/2}$	4.76E+00	7.69E-02	3.91E-01	-1.94E+00	1.77E+00	-9.02E-03
$2s^2 2p^3 2D_{3/2}$	-	$2s^2 2p^3 2P_{1/2}$	1.45E+01	2.45E-02	4.23E-02	-1.22E-01	8.68E-02	-1.04E-03
$2s^2 2p^3 2D_{3/2}$	-	$2s^2 2p^3 2P_{3/2}$	2.22E+01	3.02E-02	9.95E-02	-2.52E-01	1.61E-01	-4.36E-03
$2s^2 2p^3 2D_{3/2}$	-	$2s^2 p^4 4P_{5/2}$	7.55E+01	2.59E-02	4.85E-02	-5.33E-02	2.32E-02	1.03E-02
$2s^2 2p^3 2D_{3/2}$	-	$2s^2 p^4 4P_{3/2}$	8.36E+01	4.60E-03	4.10E-02	-7.24E-02	3.87E-02	-8.19E-04
$2s^2 2p^3 2D_{3/2}$	-	$2s^2 p^4 4P_{1/2}$	8.63E+01	3.38E-03	1.79E-02	-3.12E-02	1.61E-02	-5.42E-04
$2s^2 2p^3 2D_{3/2}$	-	$2s^2 p^4 2D_{3/2}$	1.12E+02	3.11E-01	1.67E-01	1.74E-01	-1.60E-01	1.51E-01
$2s^2 2p^3 2D_{3/2}$	-	$2s^2 p^4 2D_{5/2}$	1.15E+02	1.08E-03	2.80E-02	-4.22E-02	2.12E-02	1.15E-04
$2s^2 2p^3 2D_{3/2}$	-	$2s^2 p^4 2S_{1/2}$	1.31E+02	1.11E-01	-5.01E-02	2.95E-01	-1.95E-01	4.98E-02
$2s^2 2p^3 2D_{3/2}$	-	$2s^2 p^4 2P_{3/2}$	1.38E+02	6.14E-02	-3.03E-03	1.05E-01	-7.19E-02	3.01E-02
$2s^2 2p^3 2D_{3/2}$	-	$2s^2 p^4 2P_{1/2}$	1.50E+02	4.85E-02	-3.34E-02	1.56E-01	-1.02E-01	2.22E-02

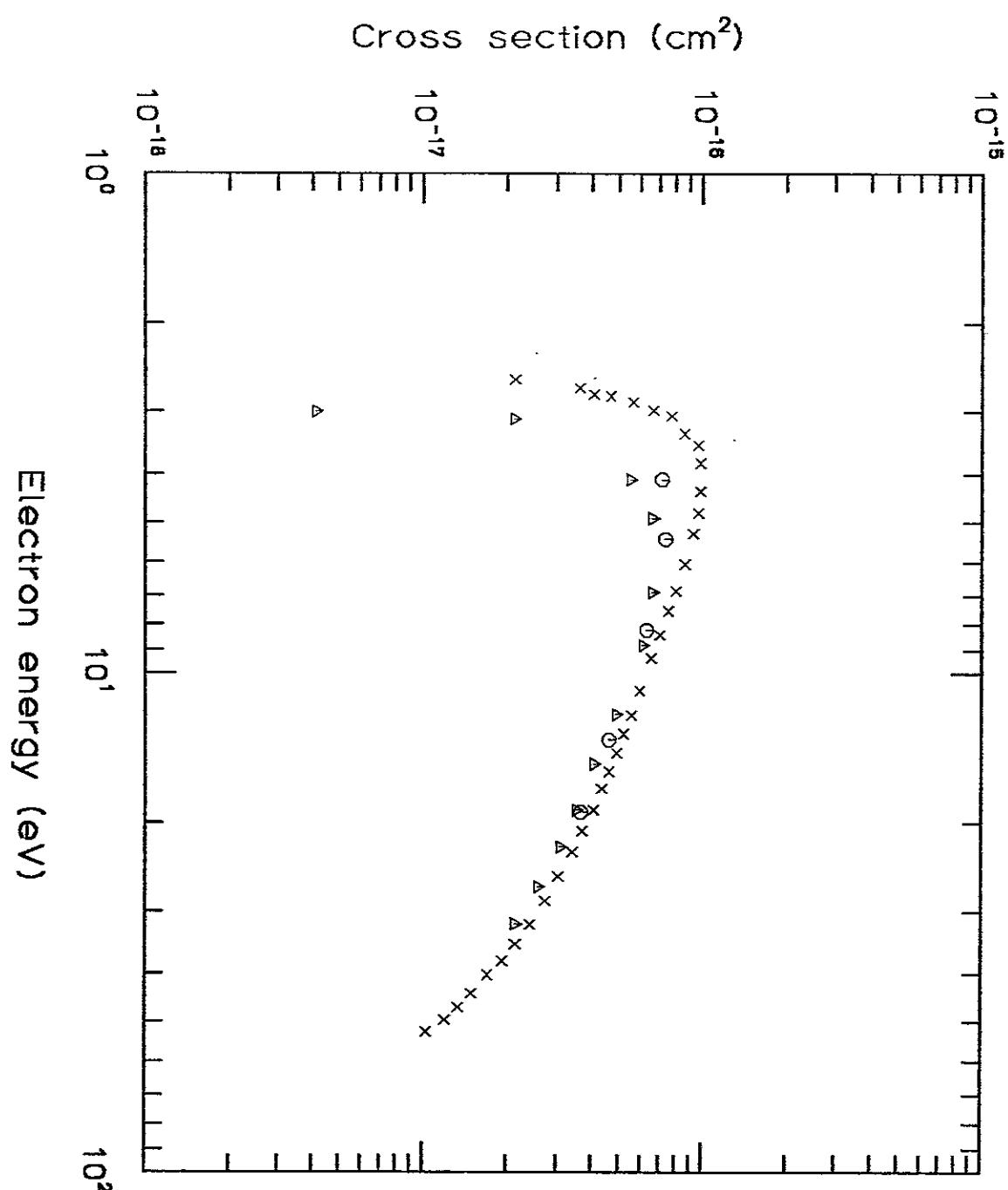
TABLE III. (continued)

Fe¹⁹⁺	ΔE (eV)	A	B	C	D	E
$2s^2 2p^3 \ ^2D_{5/2} - 2s^2 2p^3 \ ^2P_{1/2}$	9.74E+00	2.74E-02	5.08E-02	-1.84E-01	1.45E-01	-1.82E-03
$2s^2 2p^3 \ ^2D_{5/2} - 2s^2 2p^3 \ ^2P_{3/2}$	1.75E+01	6.24E-02	7.53E-02	-2.08E-01	1.42E-01	-3.79E-03
$2s^2 2p^3 \ ^2D_{5/2} - 2s2p^4 \ ^4P_{5/2}$	7.07E+01	1.17E-02	1.16E-01	-2.16E-01	1.21E-01	5.40E-03
$2s^2 2p^3 \ ^2D_{5/2} - 2s2p^4 \ ^4P_{3/2}$	7.88E+01	7.02E-04	2.73E-02	-3.65E-02	1.54E-02	4.90E-04
$2s^2 2p^3 \ ^2D_{5/2} - 2s2p^4 \ ^4P_{1/2}$	8.15E+01	2.98E-04	2.85E-03	-4.86E-03	2.54E-03	-5.62E-05
$2s^2 2p^3 \ ^2D_{5/2} - 2s2p^4 \ ^2D_{3/2}$	1.08E+02	2.34E-03	2.82E-02	-3.07E-02	1.10E-02	2.01E-04
$2s^2 2p^3 \ ^2D_{5/2} - 2s2p^4 \ ^2D_{5/2}$	1.10E+02	4.05E-01	-1.77E-02	7.10E-01	-5.02E-01	1.89E-01
$2s^2 2p^3 \ ^2D_{5/2} - 2s2p^4 \ ^2S_{1/2}$	1.26E+02	3.64E-05	2.51E-03	-3.93E-03	1.95E-03	-8.97E-06
$2s^2 2p^3 \ ^2D_{5/2} - 2s2p^4 \ ^2P_{3/2}$	1.33E+02	4.86E-01	-3.07E-01	1.45E+00	-9.48E-01	2.21E-01
$2s^2 2p^3 \ ^2D_{5/2} - 2s2p^4 \ ^2P_{1/2}$	1.45E+02	5.40E-05	4.09E-03	-5.69E-03	2.65E-03	-1.63E-05
$2s^2 2p^3 \ ^2P_{1/2} - 2s^2 2p^3 \ ^2P_{3/2}$	7.73E+00	2.48E-02	2.35E-01	-8.50E-01	6.81E-01	-2.72E-03
$2s^2 2p^3 \ ^2P_{1/2} - 2s2p^4 \ ^4P_{5/2}$	6.10E+01	4.13E-04	7.00E-03	-1.42E-02	8.29E-03	-6.91E-05
$2s^2 2p^3 \ ^2P_{1/2} - 2s2p^4 \ ^4P_{3/2}$	6.91E+01	8.24E-04	1.89E-02	-3.06E-02	1.58E-02	6.01E-05
$2s^2 2p^3 \ ^2P_{1/2} - 2s2p^4 \ ^4P_{1/2}$	7.18E+01	2.23E-03	5.95E-02	-1.18E-01	6.75E-02	1.61E-03
$2s^2 2p^3 \ ^2P_{1/2} - 2s2p^4 \ ^2D_{3/2}$	9.79E+01	4.00E-02	-2.78E-03	7.68E-02	-5.48E-02	1.64E-02
$2s^2 2p^3 \ ^2P_{1/2} - 2s2p^4 \ ^2D_{5/2}$	1.00E+02	1.72E-03	1.88E-02	-3.14E-02	1.60E-02	-3.29E-04
$2s^2 2p^3 \ ^2P_{1/2} - 2s2p^4 \ ^2S_{1/2}$	1.17E+02	1.41E-01	-1.44E-02	2.78E-01	-1.96E-01	6.07E-02
$2s^2 2p^3 \ ^2P_{1/2} - 2s2p^4 \ ^2P_{3/2}$	1.23E+02	5.61E-02	-1.32E-02	1.24E-01	-8.51E-02	2.48E-02
$2s^2 2p^3 \ ^2P_{1/2} - 2s2p^4 \ ^2P_{1/2}$	1.35E+02	1.06E-02	-6.26E-03	2.85E-02	-1.85E-02	4.71E-03
$2s^2 2p^3 \ ^2P_{3/2} - 2s2p^4 \ ^4P_{5/2}$	5.32E+01	3.36E-03	2.35E-02	-5.57E-02	3.41E-02	-5.95E-04
$2s^2 2p^3 \ ^2P_{3/2} - 2s2p^4 \ ^4P_{3/2}$	6.14E+01	2.32E-02	-6.16E-03	1.41E-02	-8.23E-03	6.86E-04
$2s^2 2p^3 \ ^2P_{3/2} - 2s2p^4 \ ^4P_{1/2}$	6.41E+01	2.13E-03	1.35E-02	-8.00E-03	1.95E-04	1.90E-04
$2s^2 2p^3 \ ^2P_{3/2} - 2s2p^4 \ ^2D_{3/2}$	9.02E+01	1.09E-02	4.80E-02	-7.16E-02	3.68E-02	5.08E-03
$2s^2 2p^3 \ ^2P_{3/2} - 2s2p^4 \ ^2D_{5/2}$	9.23E+01	1.50E-01	-5.98E-02	3.70E-01	-2.49E-01	6.01E-02
$2s^2 2p^3 \ ^2P_{3/2} - 2s2p^4 \ ^2S_{1/2}$	1.09E+02	1.28E-02	3.42E-02	-3.44E-02	1.40E-02	5.95E-03
$2s^2 2p^3 \ ^2P_{3/2} - 2s2p^4 \ ^2P_{3/2}$	1.16E+02	7.20E-02	-7.14E-03	1.36E-01	-9.54E-02	3.21E-02
$2s^2 2p^3 \ ^2P_{3/2} - 2s2p^4 \ ^2P_{1/2}$	1.27E+02	2.68E-01	-1.10E-01	7.22E-01	-4.93E-01	1.21E-01

Graphs I.

Comparison of the cross sections

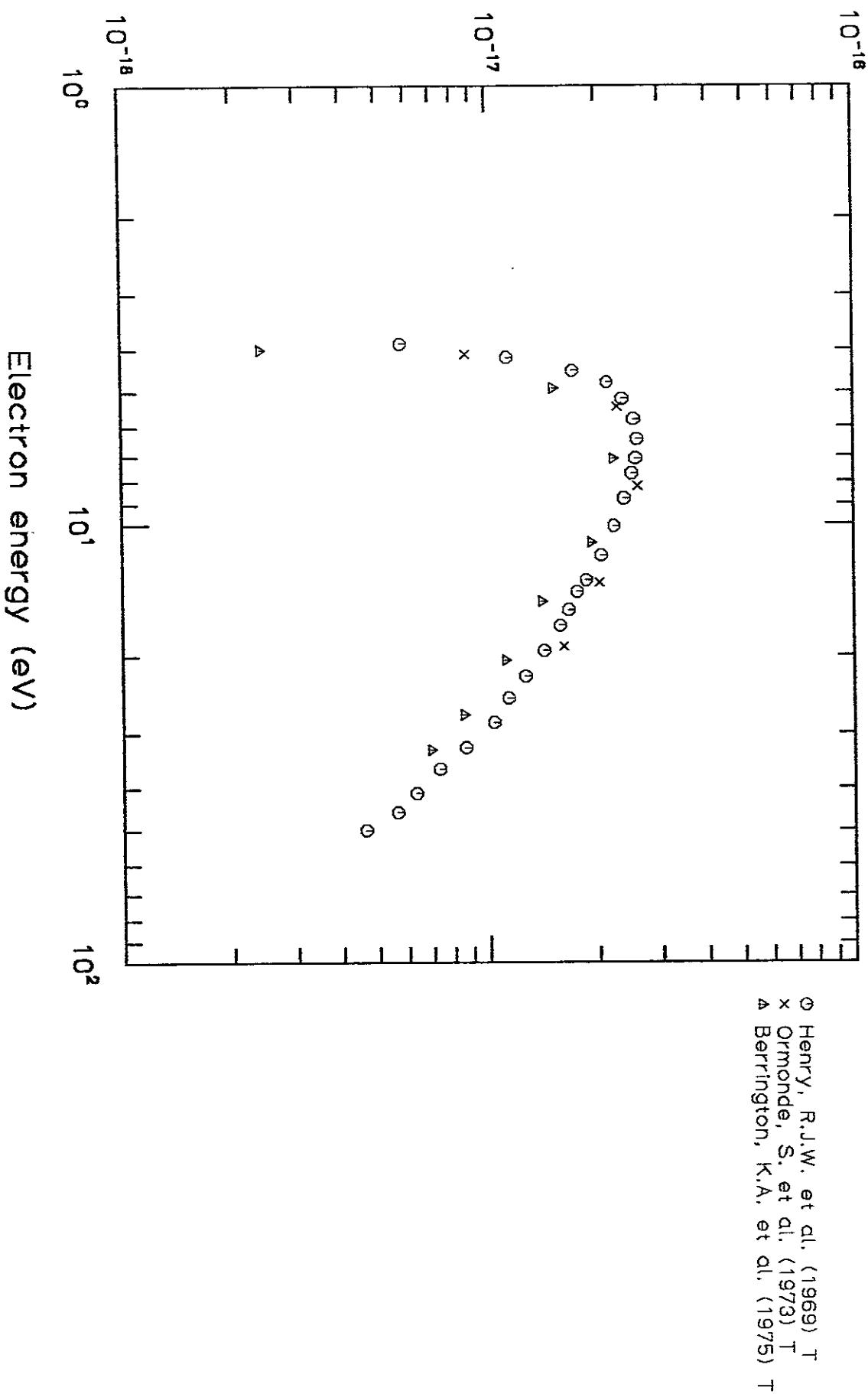
$N \{ 2s^2 2p^3 \text{ } ^4S \rightarrow 2s^2 2p^3 \text{ } ^2D \}$



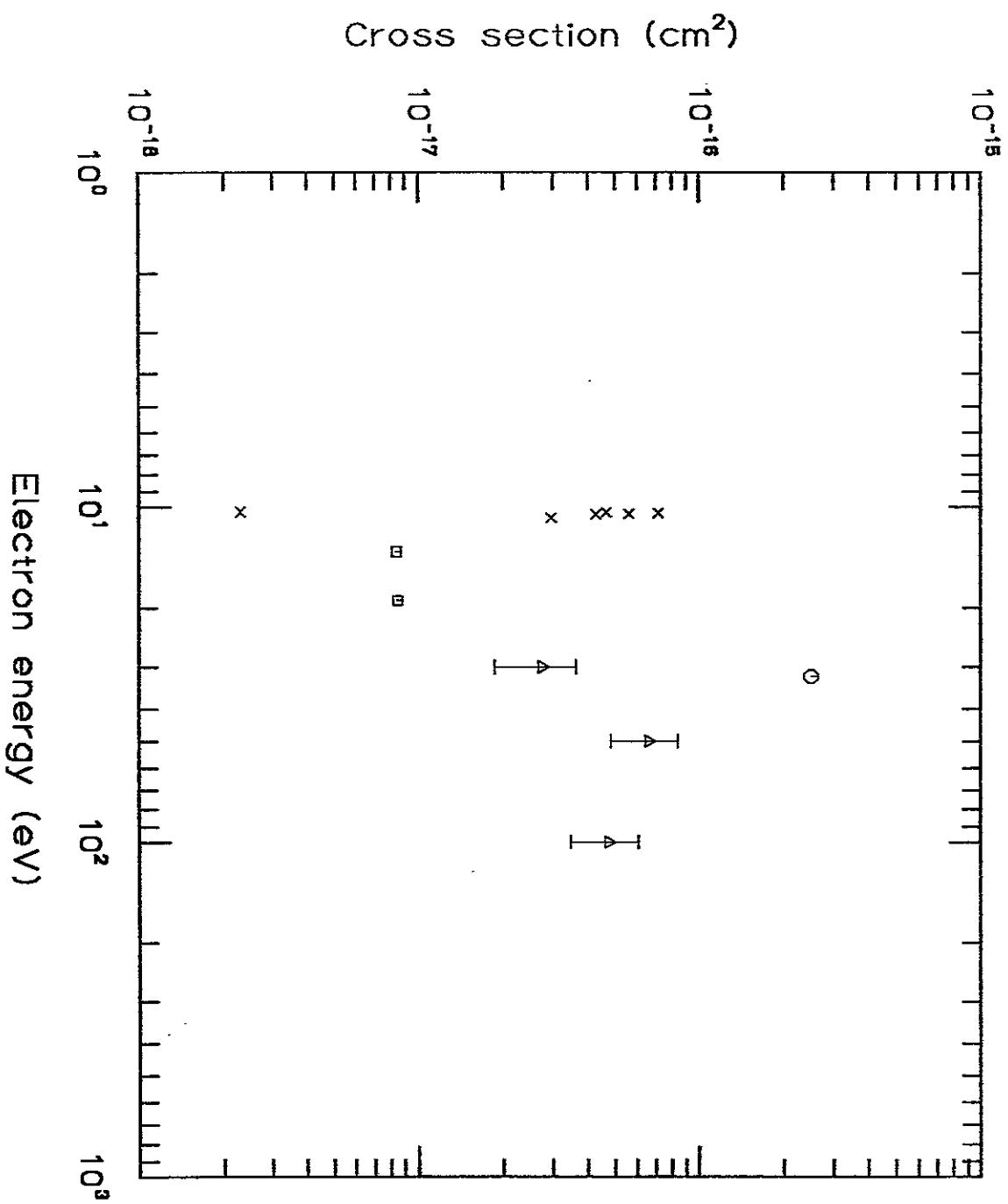
- 22 -

$N \{ 2s^2 2p^3 \text{ } ^4S \rightarrow 2s^2 2p^3 \text{ } ^2P \}$

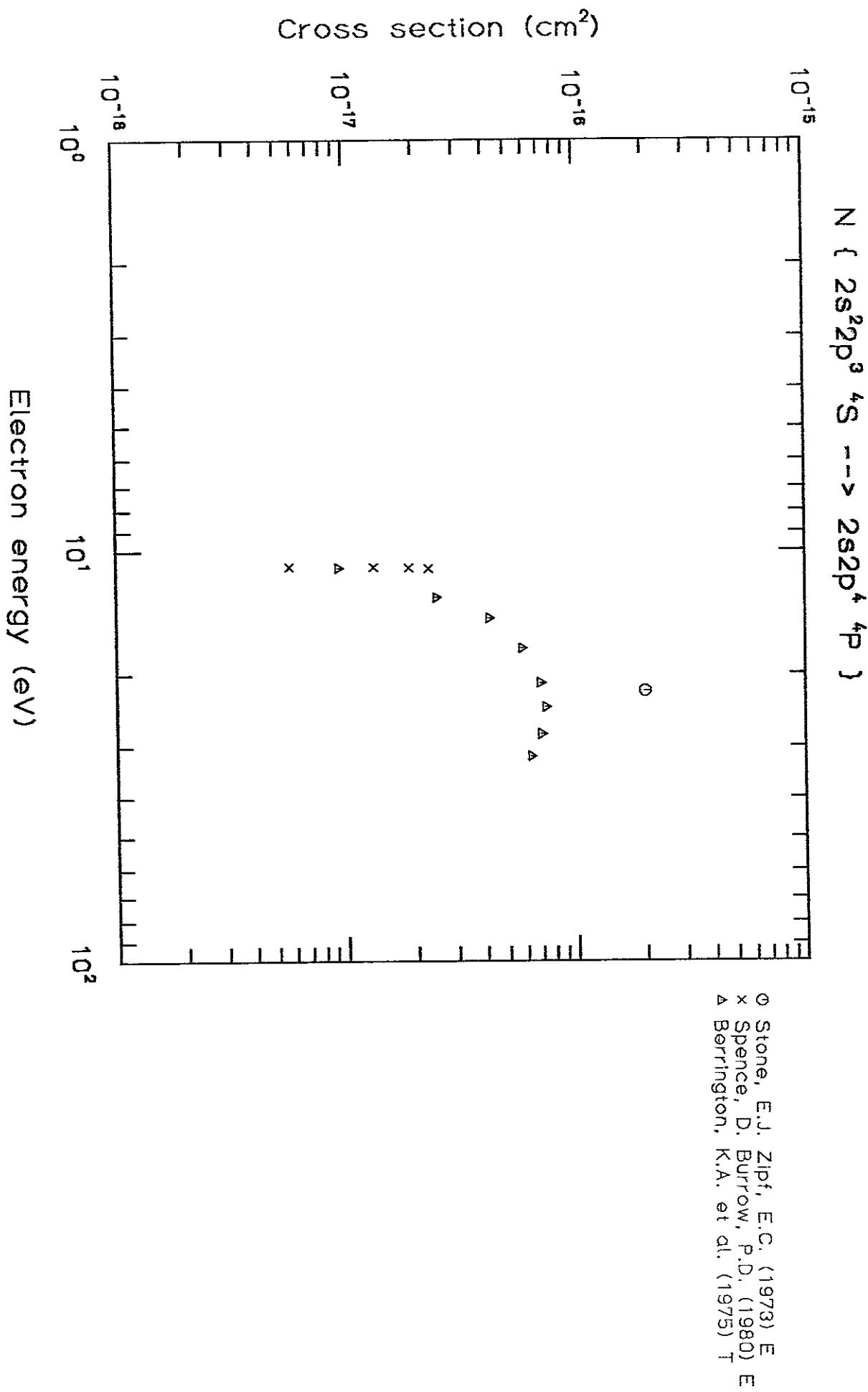
Cross section (cm^2)



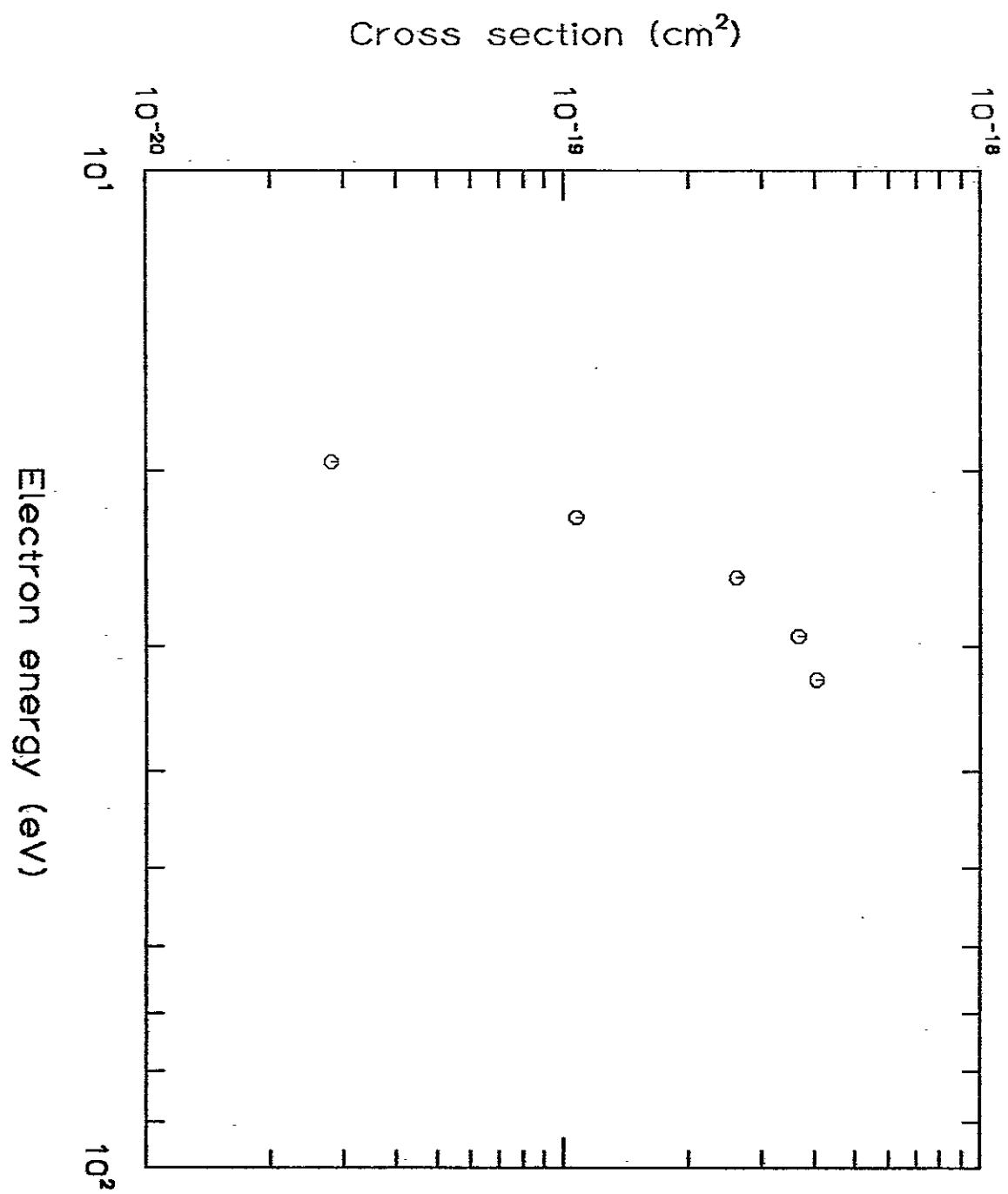
$N \{ 2s^2 2p^3 \text{ } ^4S \rightarrow 2s^2 2p^2 3s \text{ } ^4P \}$



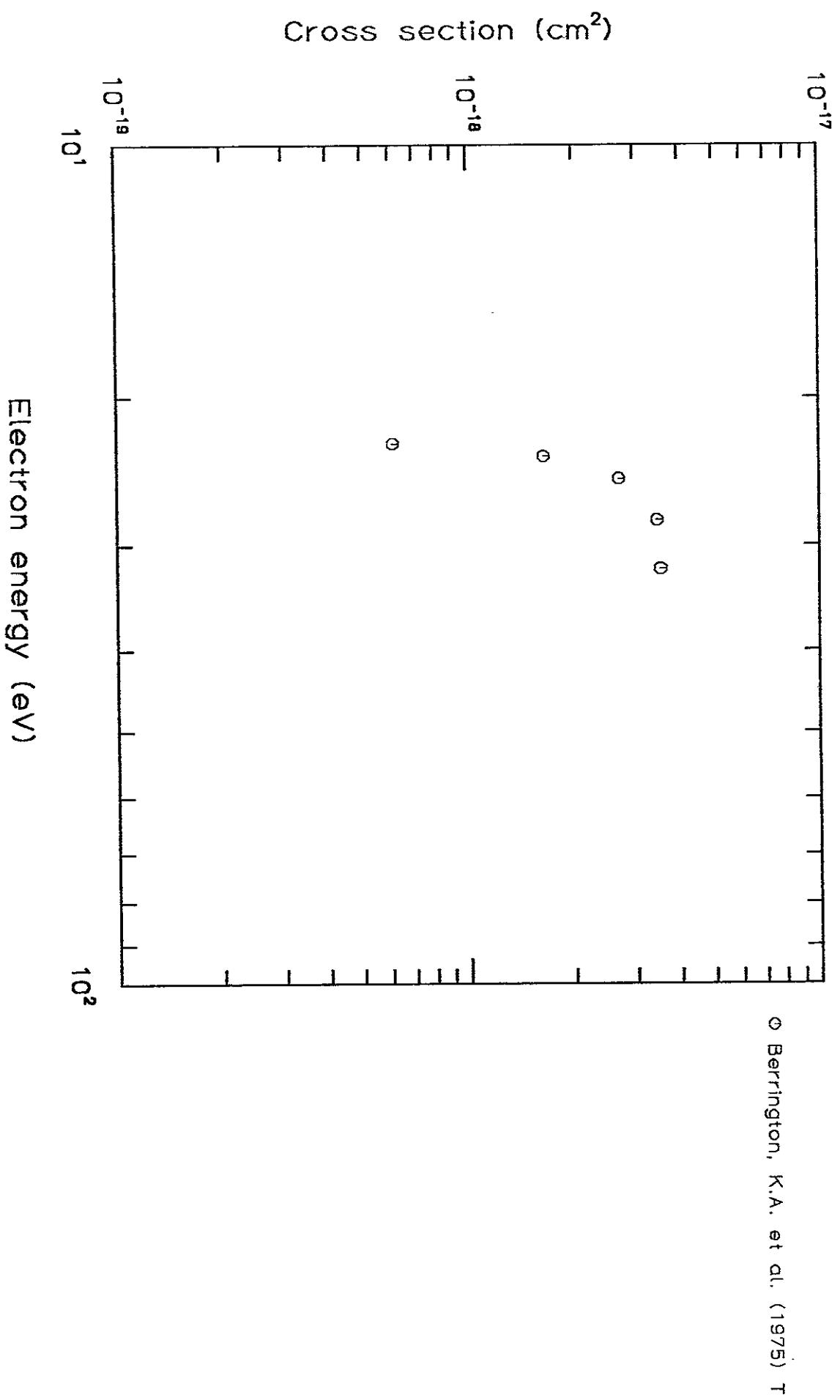
○ Stone, E.J. Zippf, E.C. (1973) E
× Spence, D. Burrow, P.D. (1980) E
▲ Doering, J.P. Goemel, L. (1991) E
■ Ormonde, S. et al. (1973) T

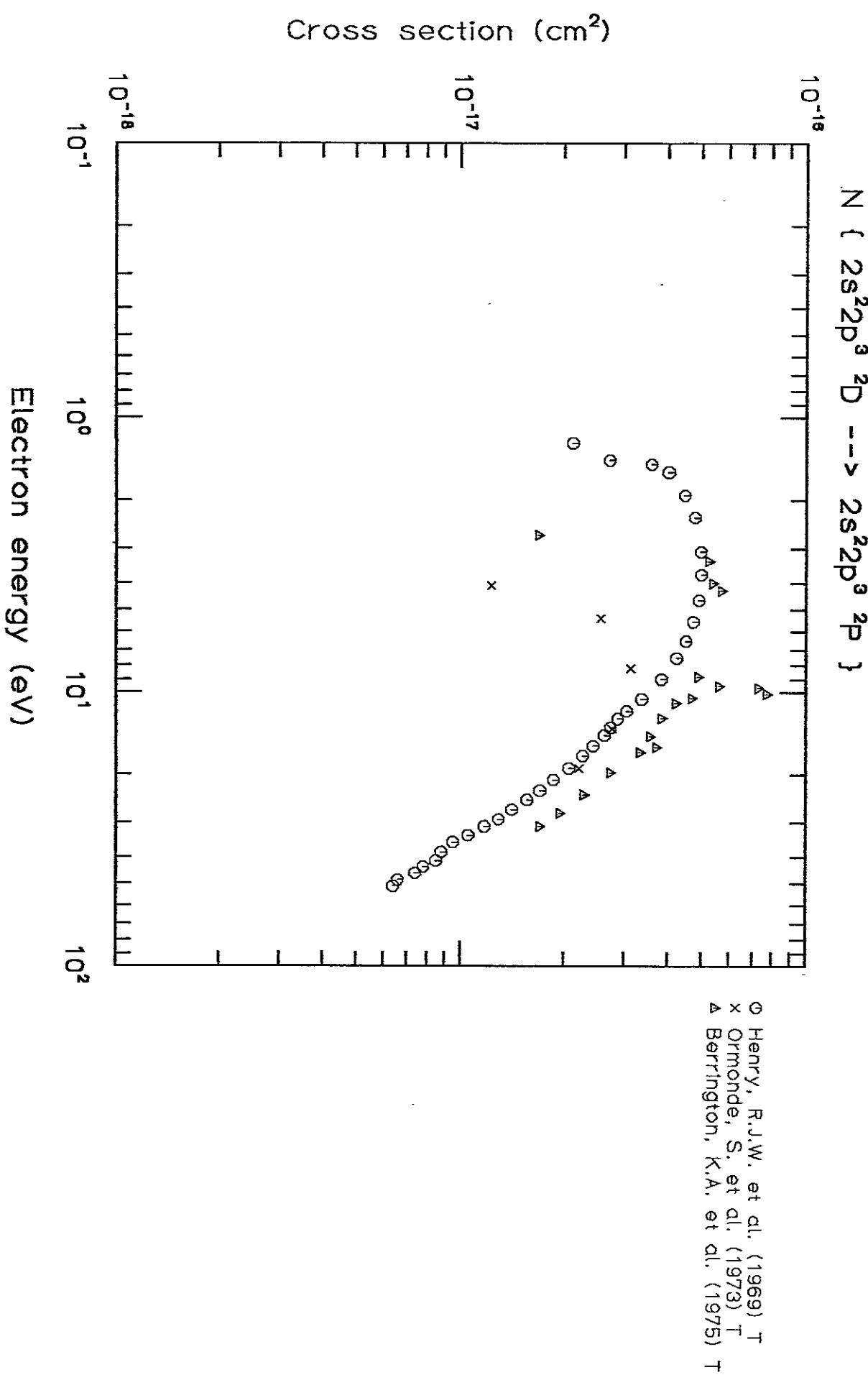


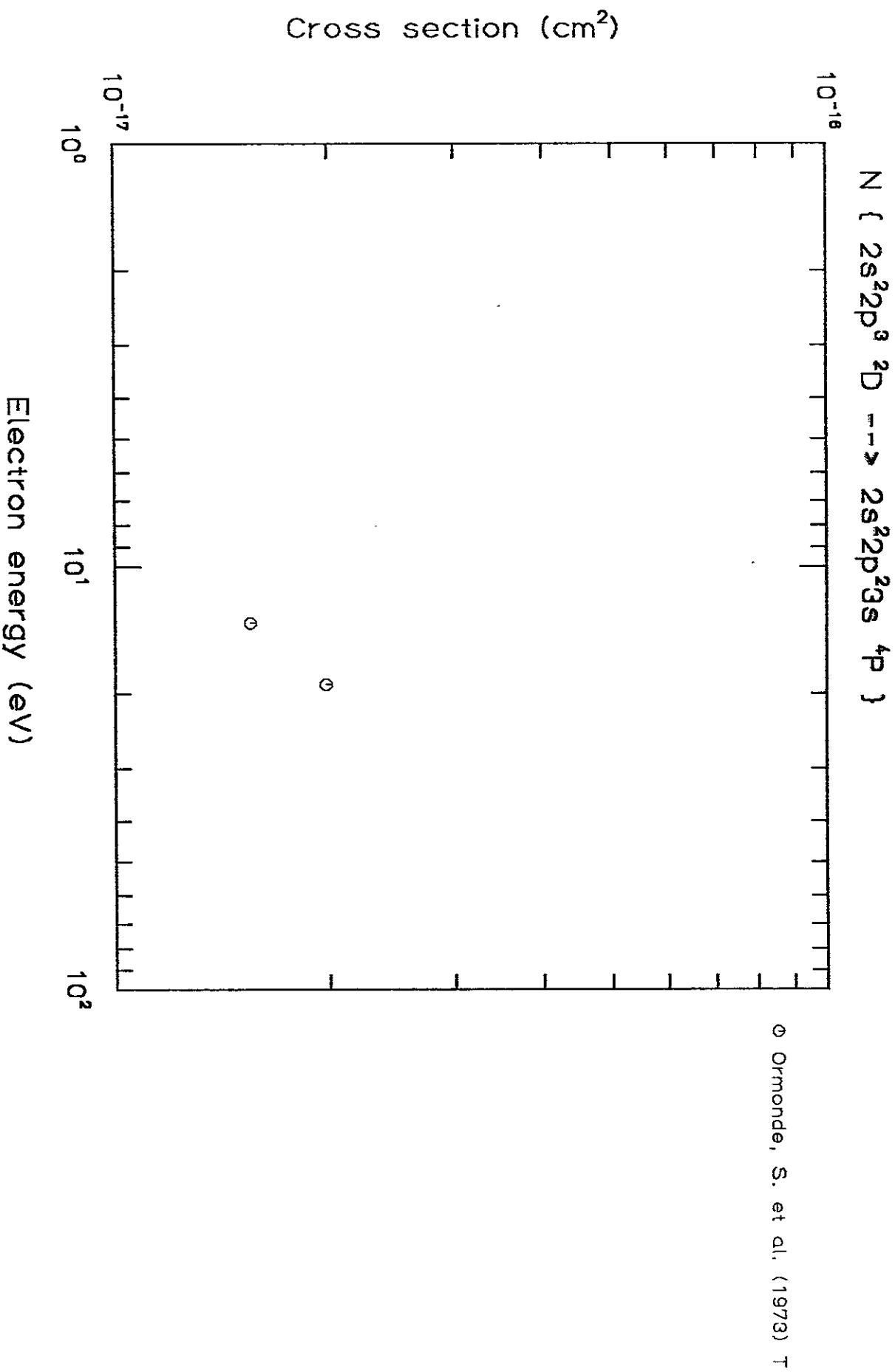
$N \{ 2s^2 2p^3 \text{ } ^4S \rightarrow 2s2p^4 \text{ } ^2D \}$



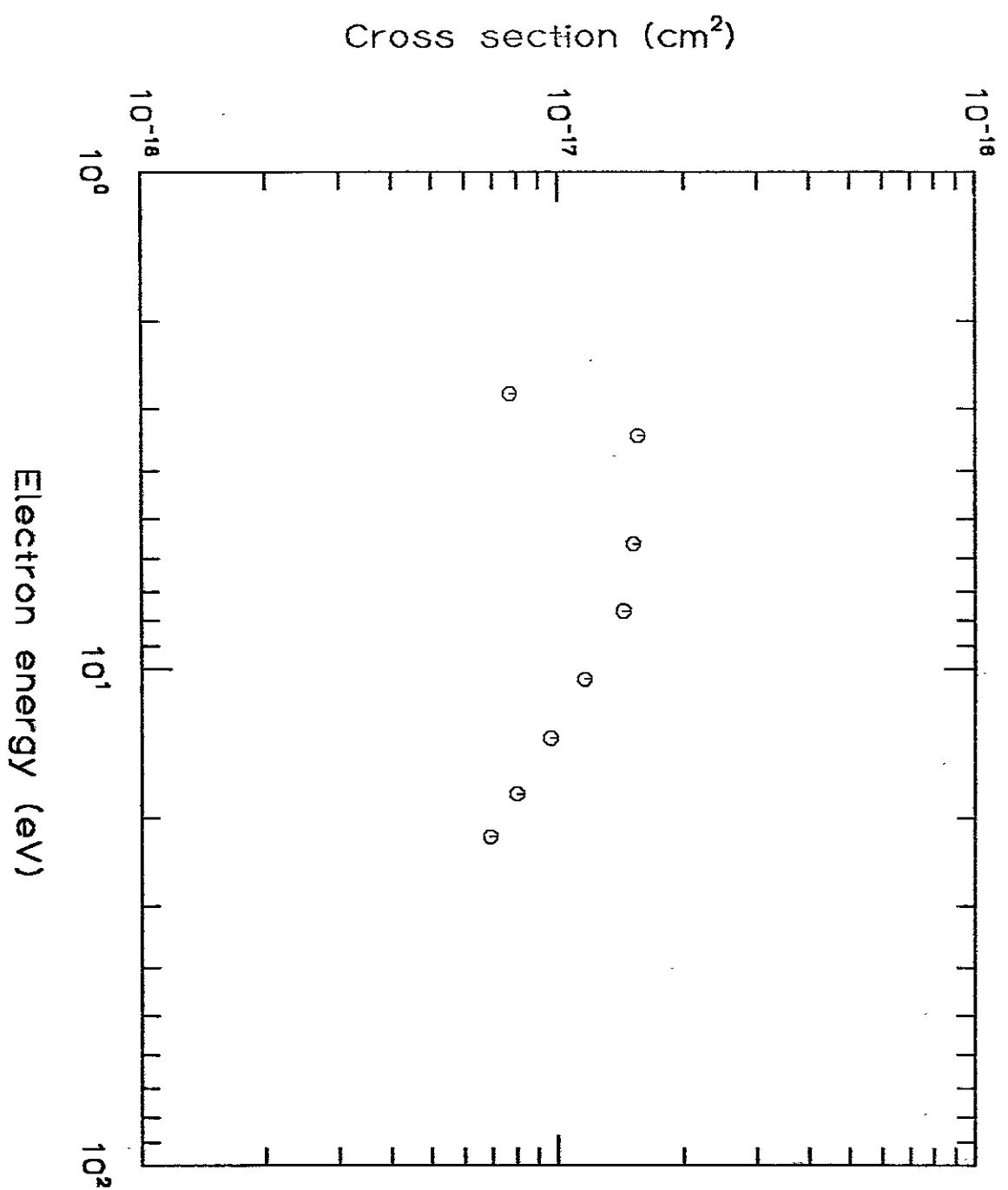
$N \{ 2s^2 2p^3 \ ^4S \rightarrow 2s 2p^4 \ ^2P \ }$



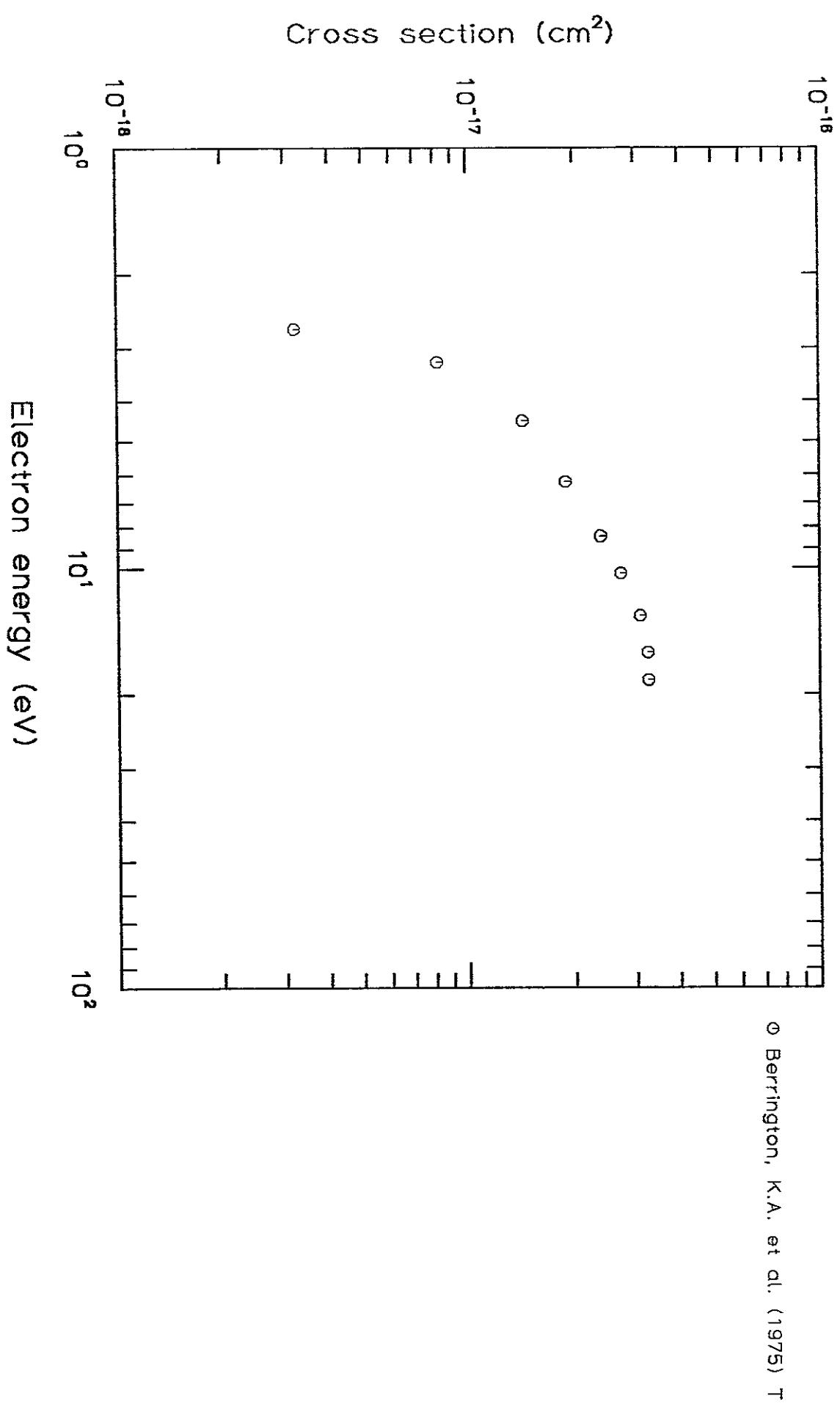




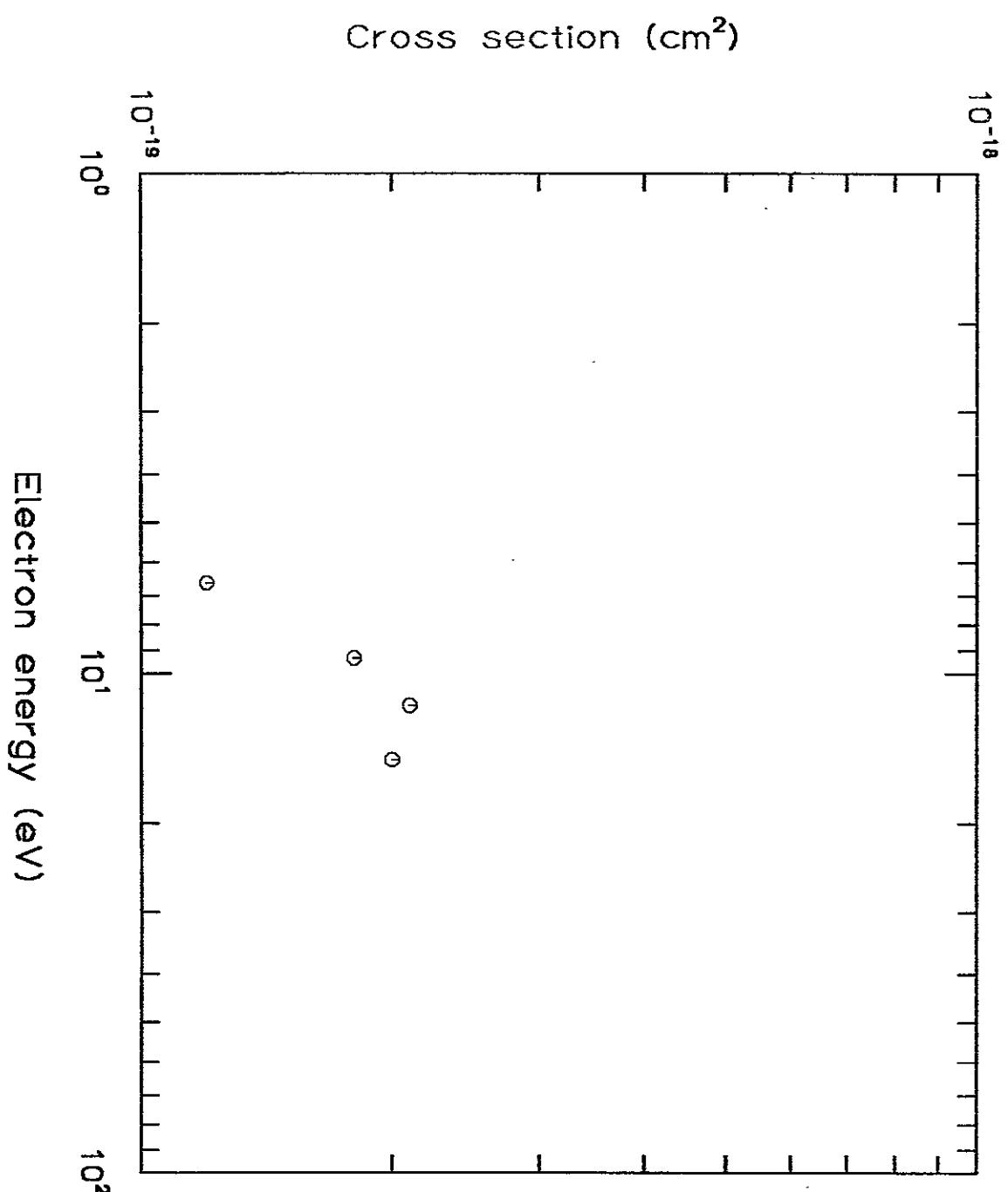
$N \{ 2s^2 2p^3 \ ^2D \rightarrow 2s2p^4 \ ^4P \ }$



$N \{ 2s^2 2p^3 \text{ } ^2D \rightarrow 2s2p^4 \text{ } ^2D \}$

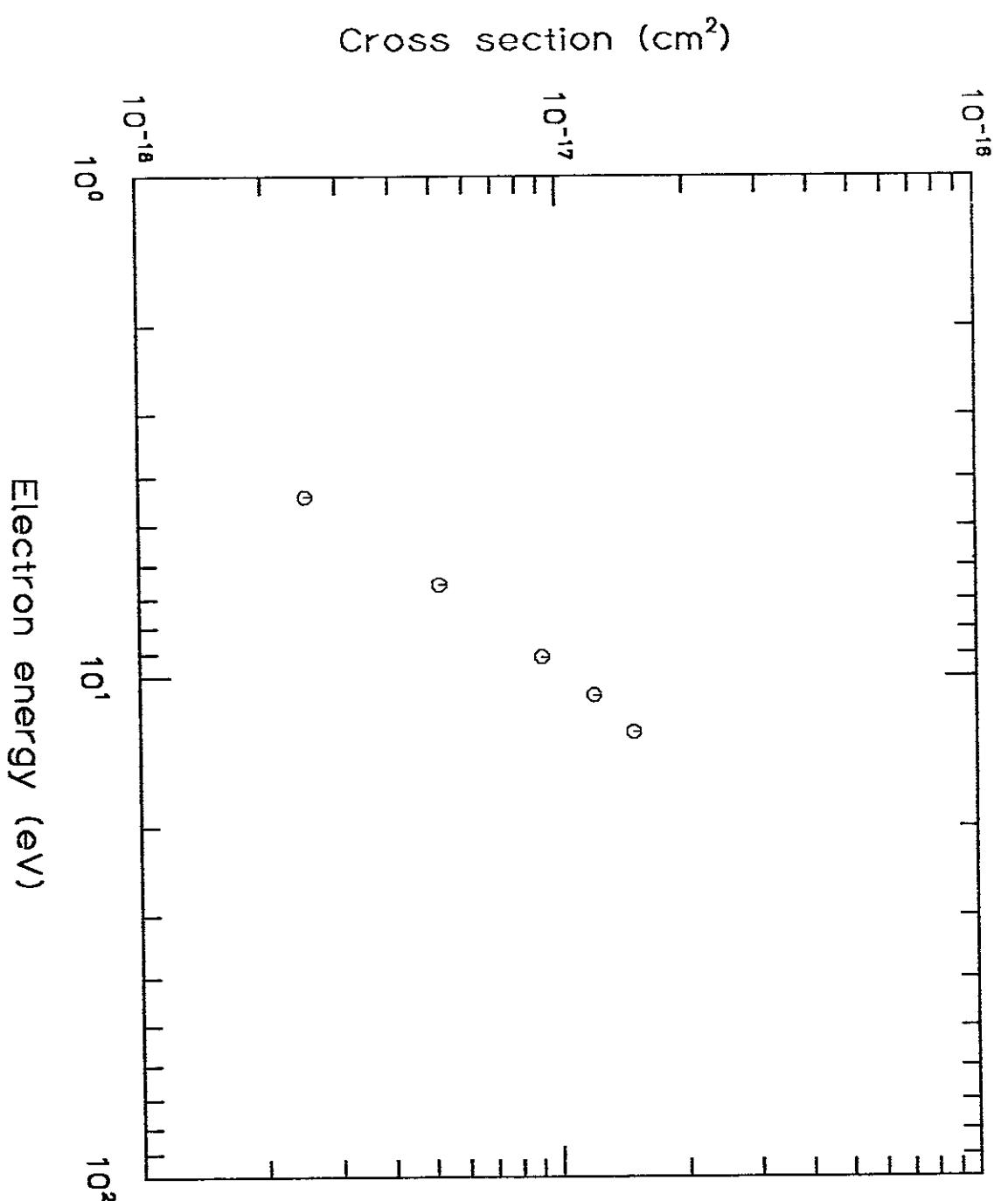


$N \{ 2s^2 2p^3 \text{ } ^2D \rightarrow 2s2p^4 \text{ } ^2S \}$

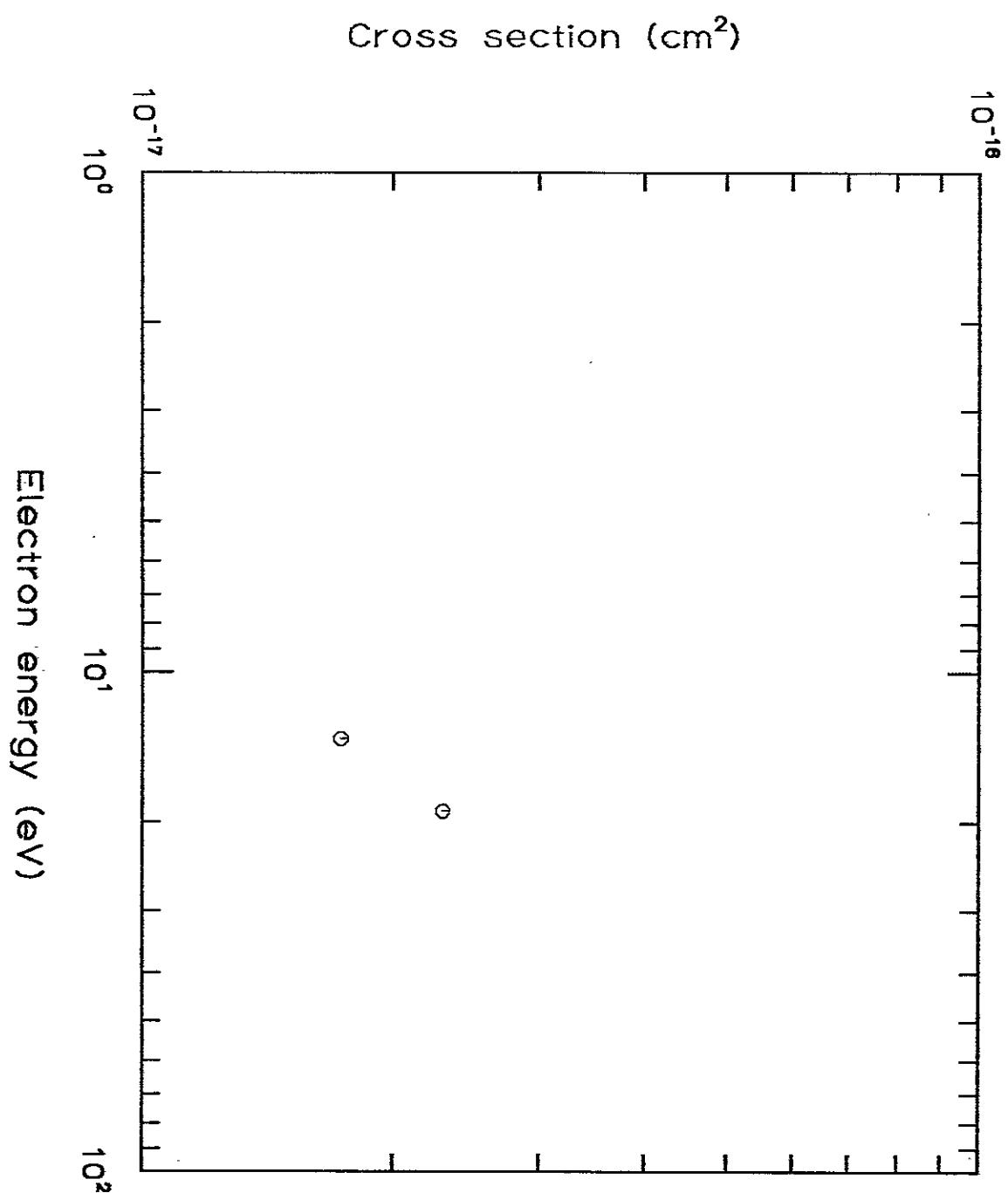


○ Barrington, K.A. et al. (1975) T

$N \{ 2s^2 2p^3 \text{ } ^2D \rightarrow 2s2p^4 \text{ } ^2P \}$

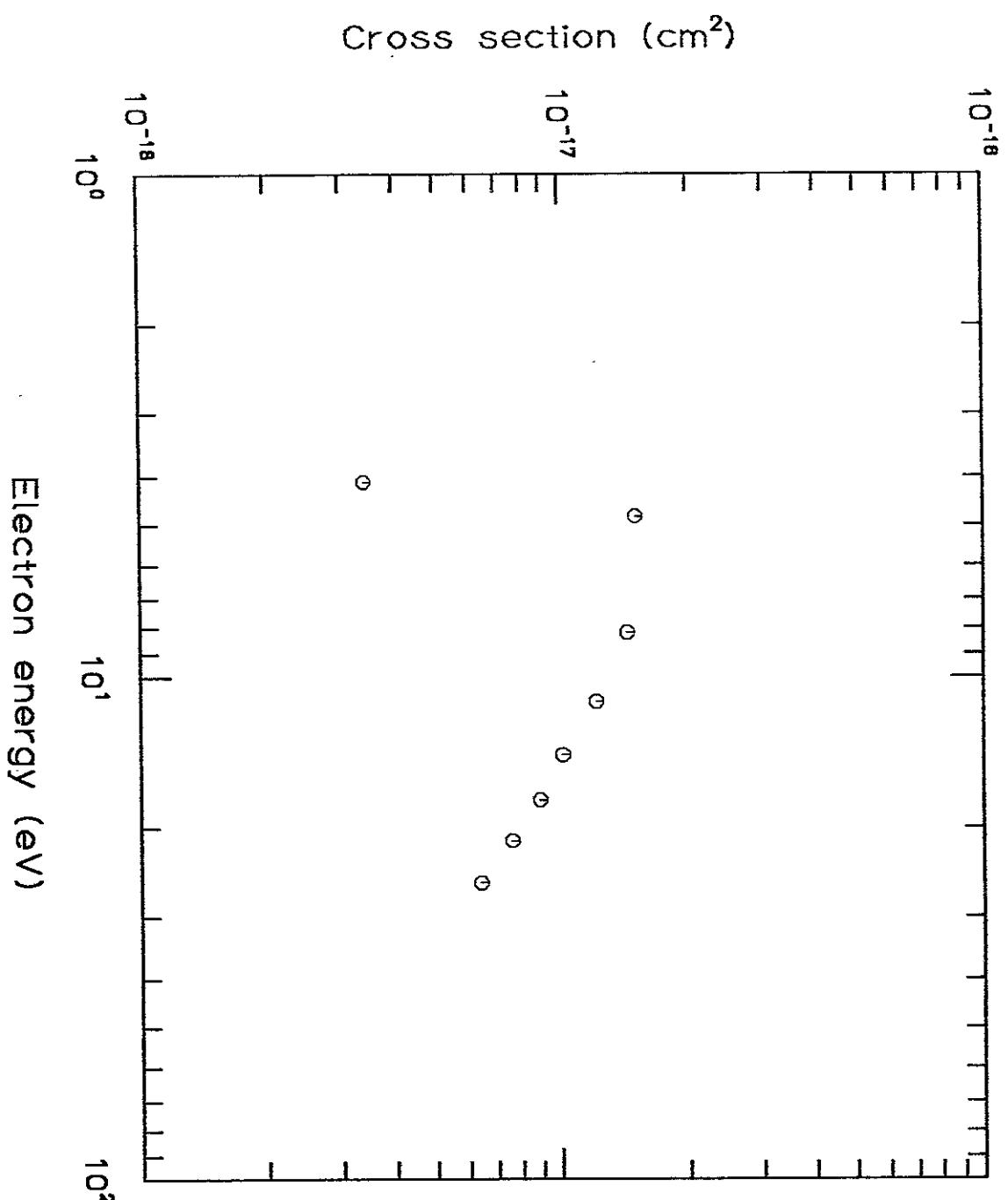


$N \{ 2s^2 2p^3 \text{ } ^2P \rightarrow 2s^2 2p^2 3s \text{ } ^4P \}$



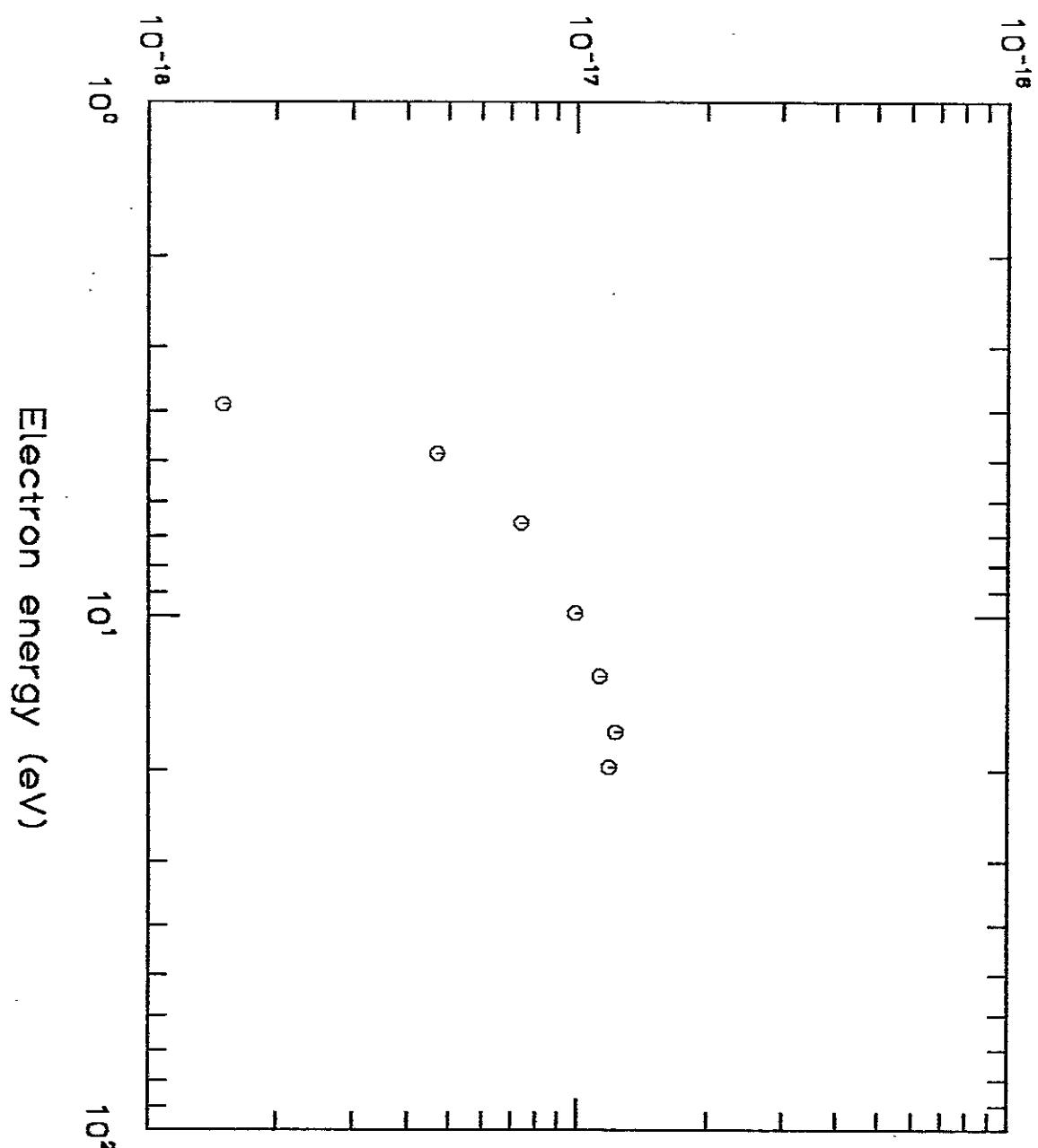
Ormonde, S. et al. (1973) T

$N \{ 2s^2 2p^3 \text{ } ^2P \rightarrow 2s2p^4 \text{ } ^4P \}$



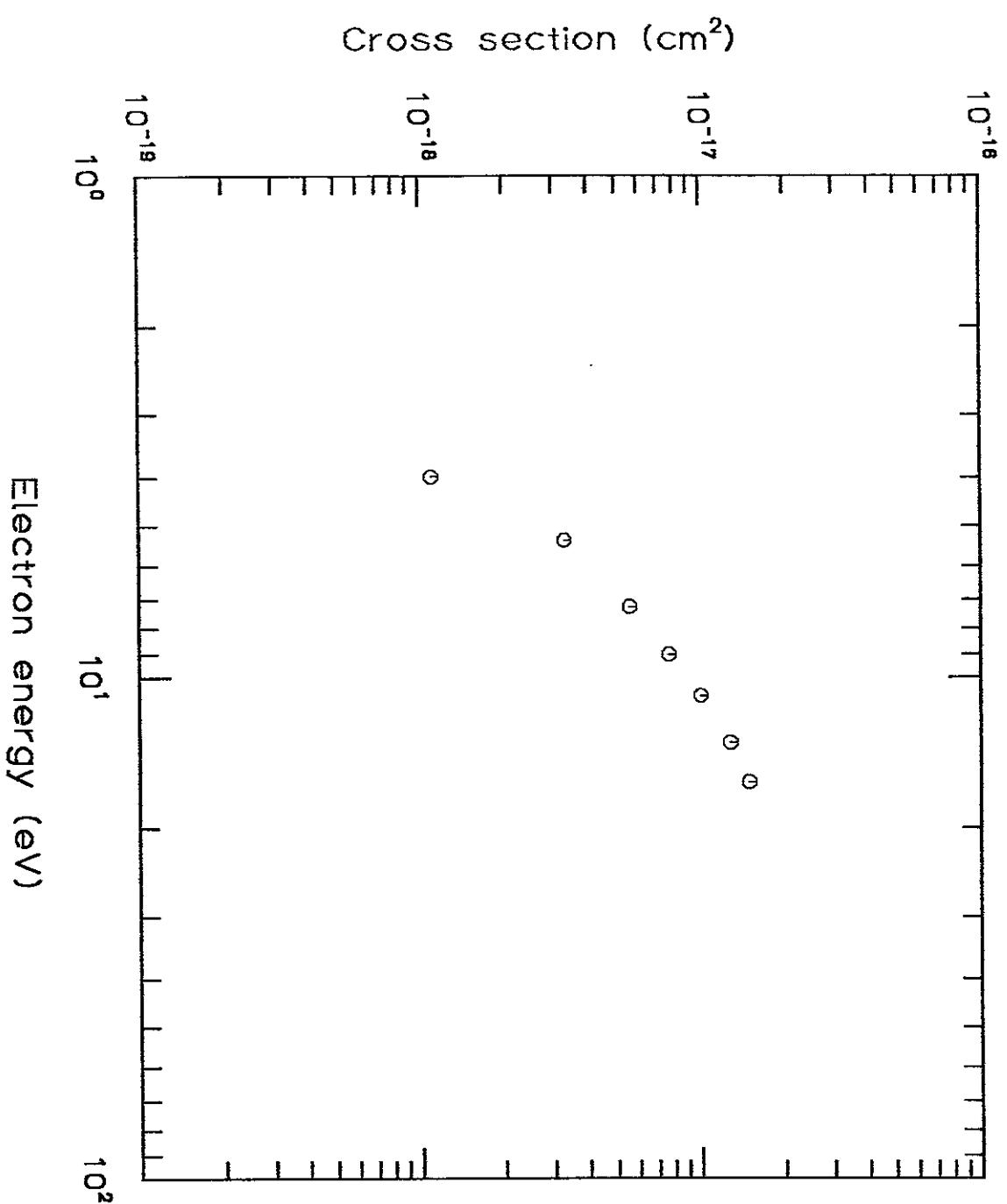
$N \{ 2s^2 2p^3 \text{ } ^2P \rightarrow 2s2p^4 \text{ } ^2D \}$

Cross section (cm^2)

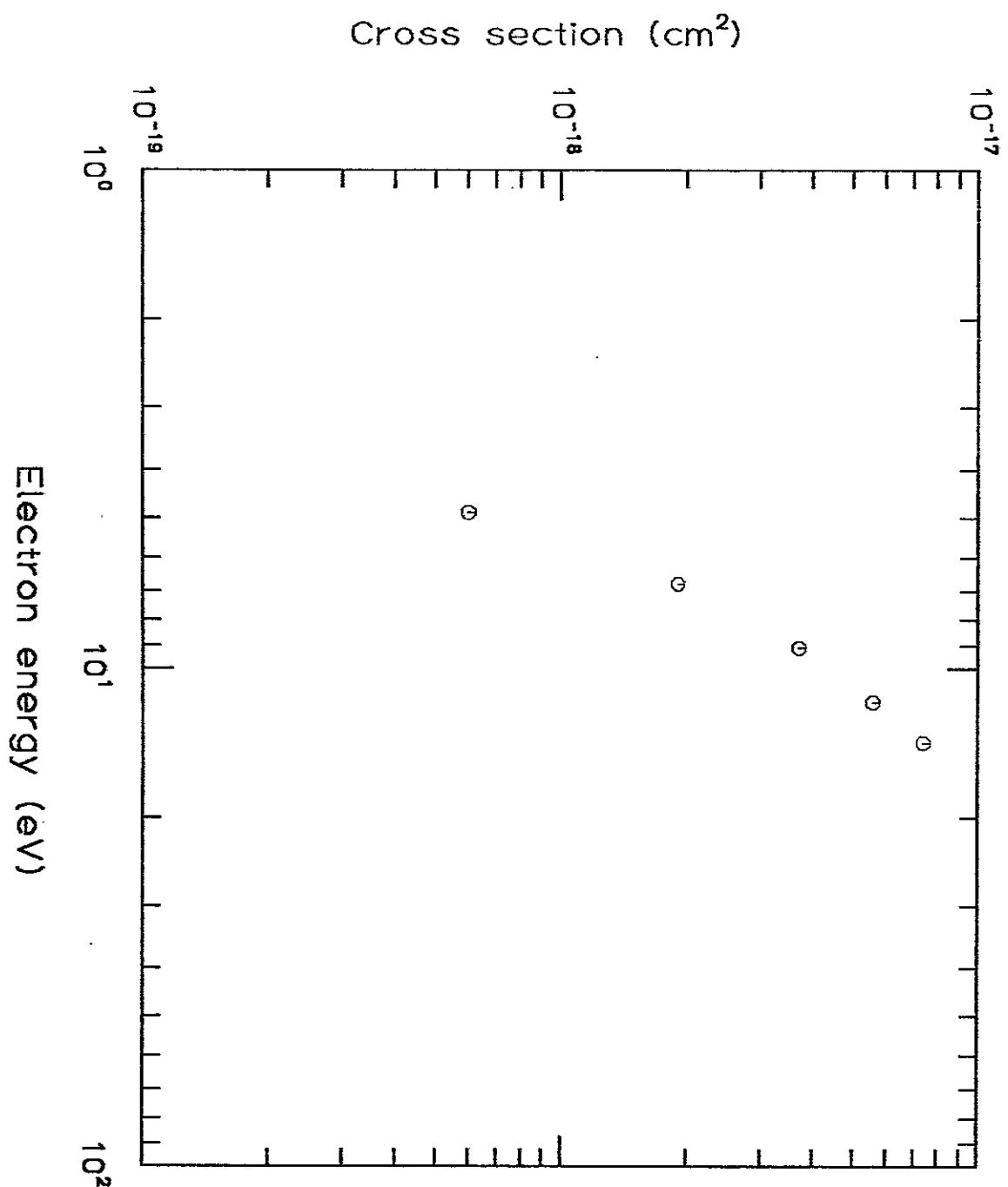


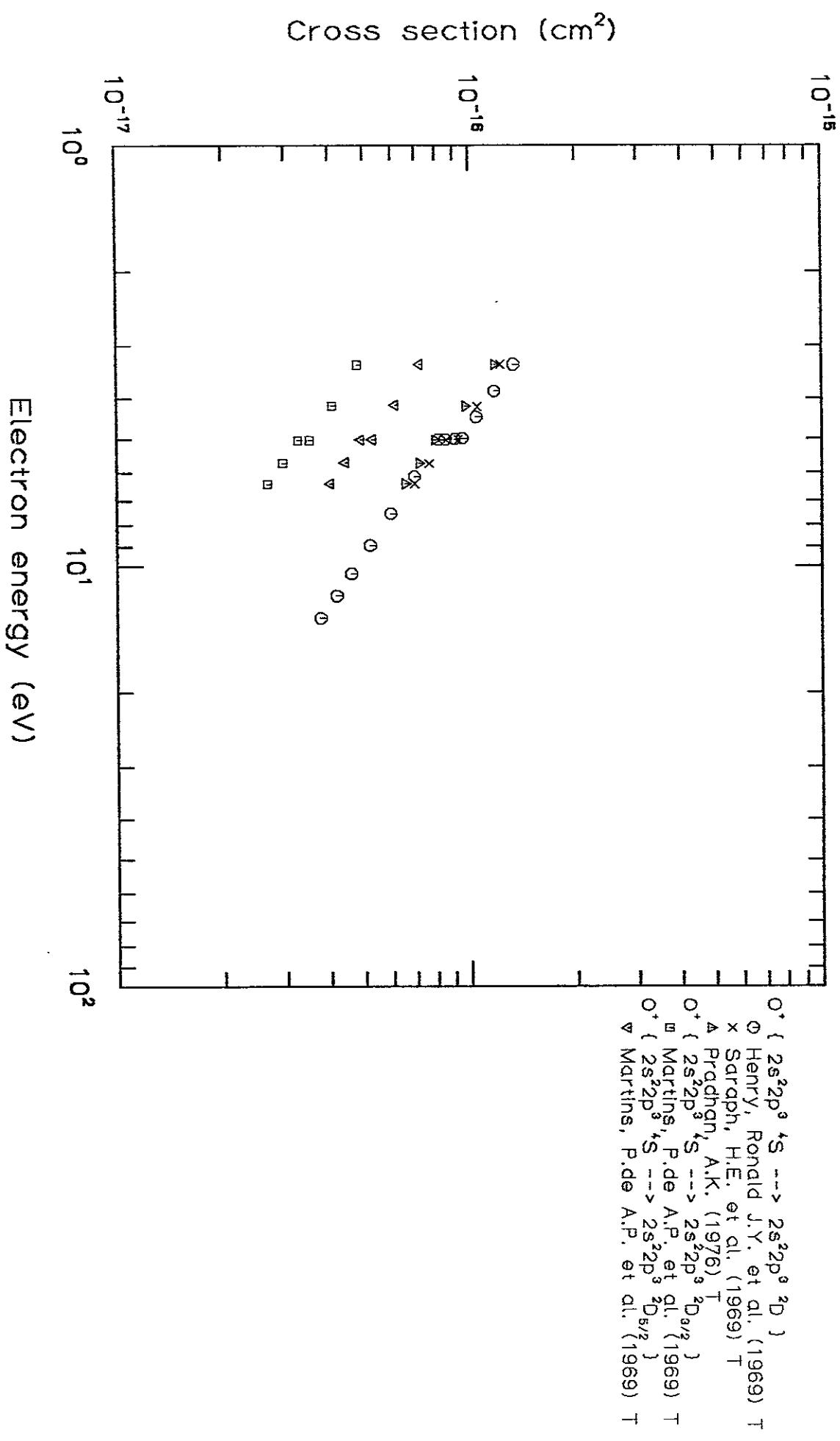
○ Berrington, K.A. et al. (1975) T

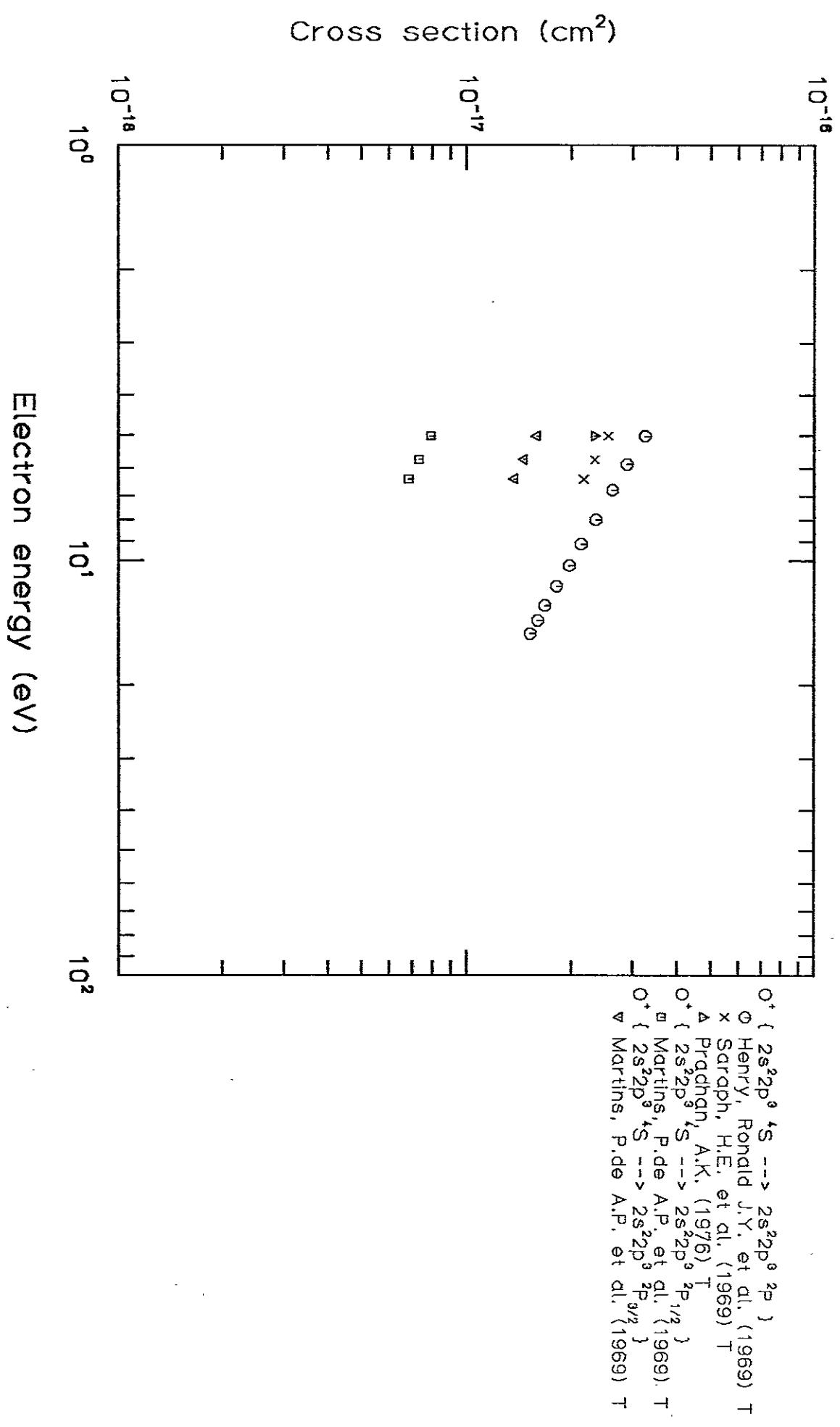
$N \{ 2s^2 2p^3 \text{ } ^2P \rightarrow 2s 2p^4 \text{ } ^2S \}$



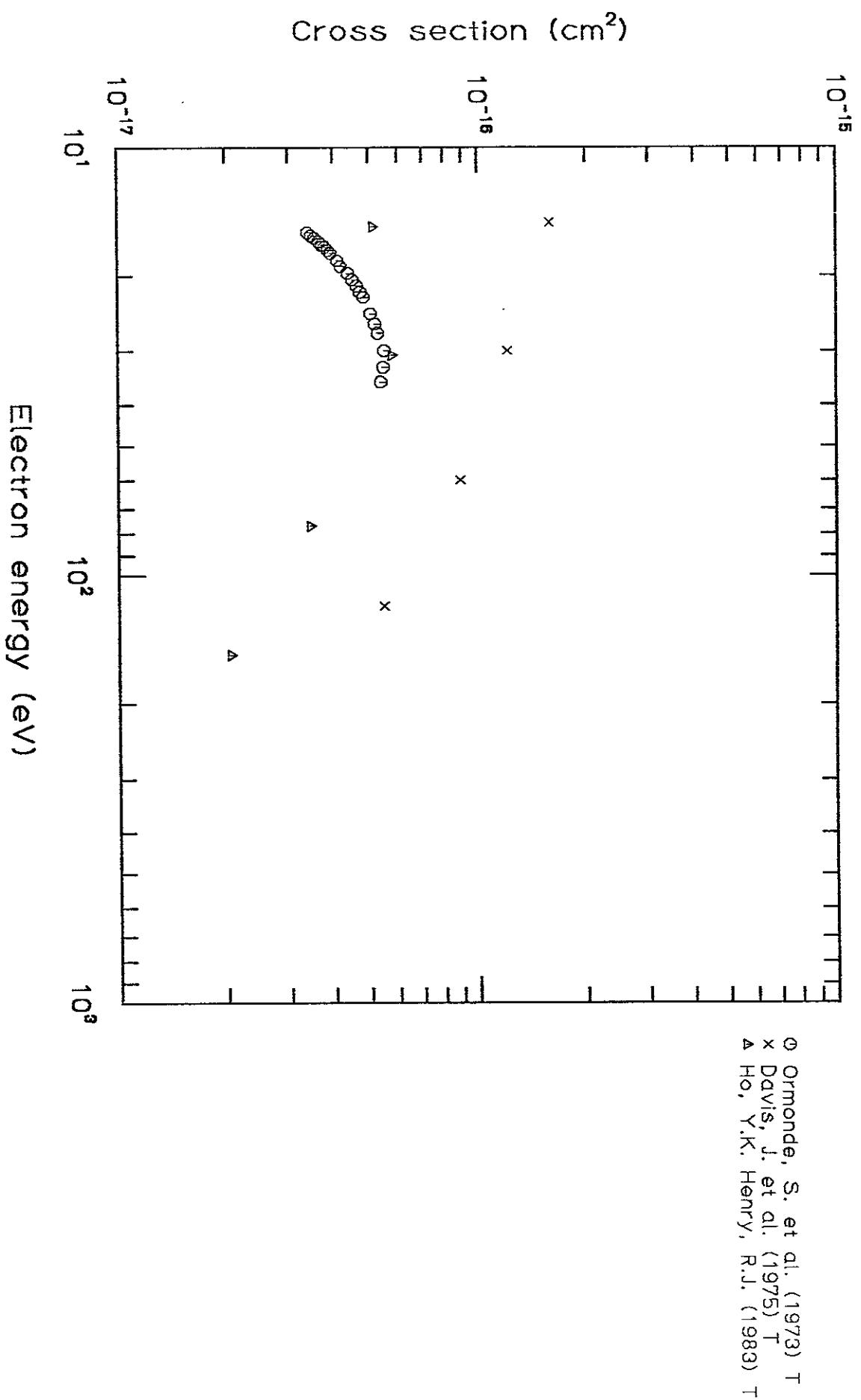
$N \{ 2s^2 2p^3 \text{ } ^2P \rightarrow 2s2p^4 \text{ } ^2P \}$



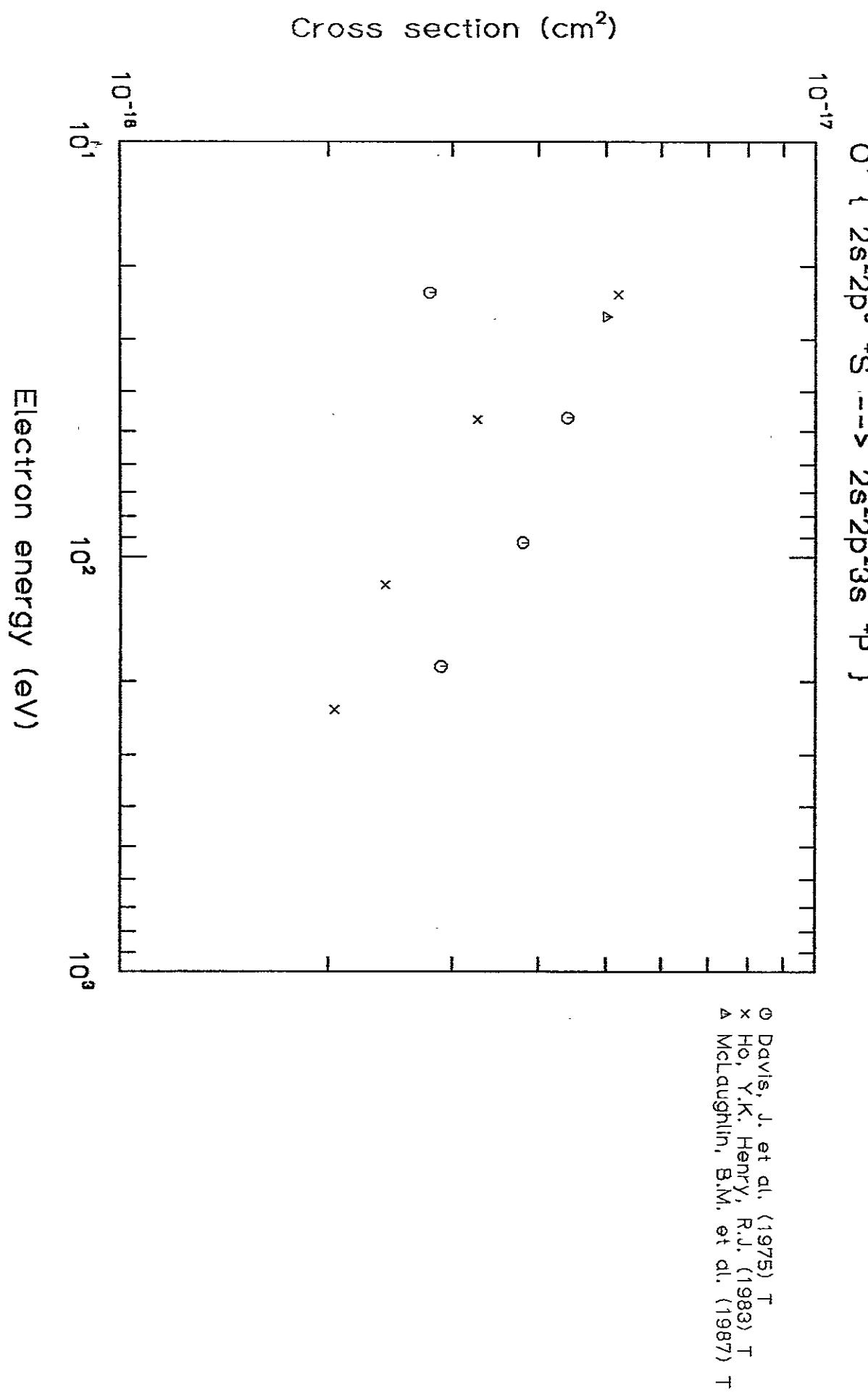


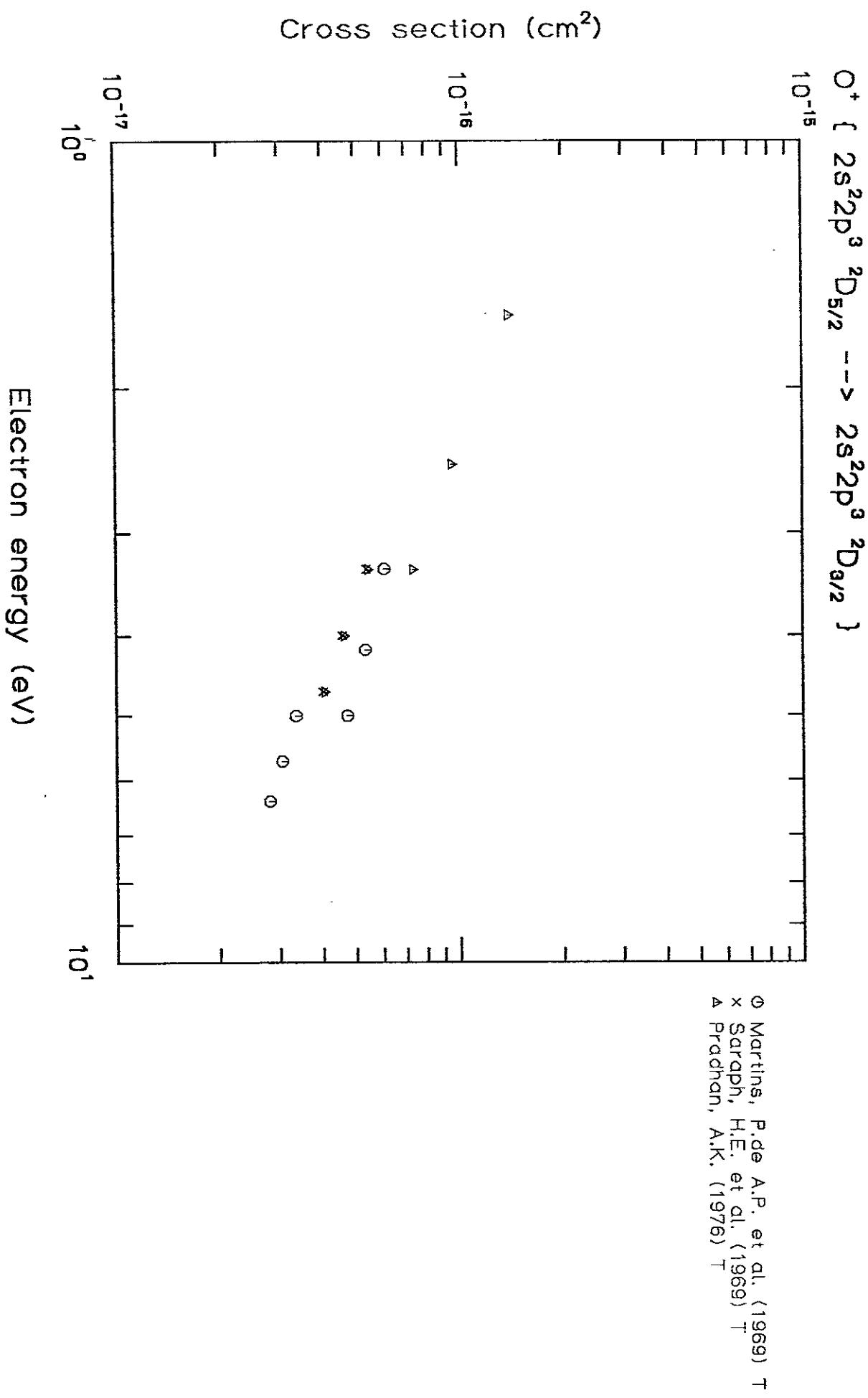


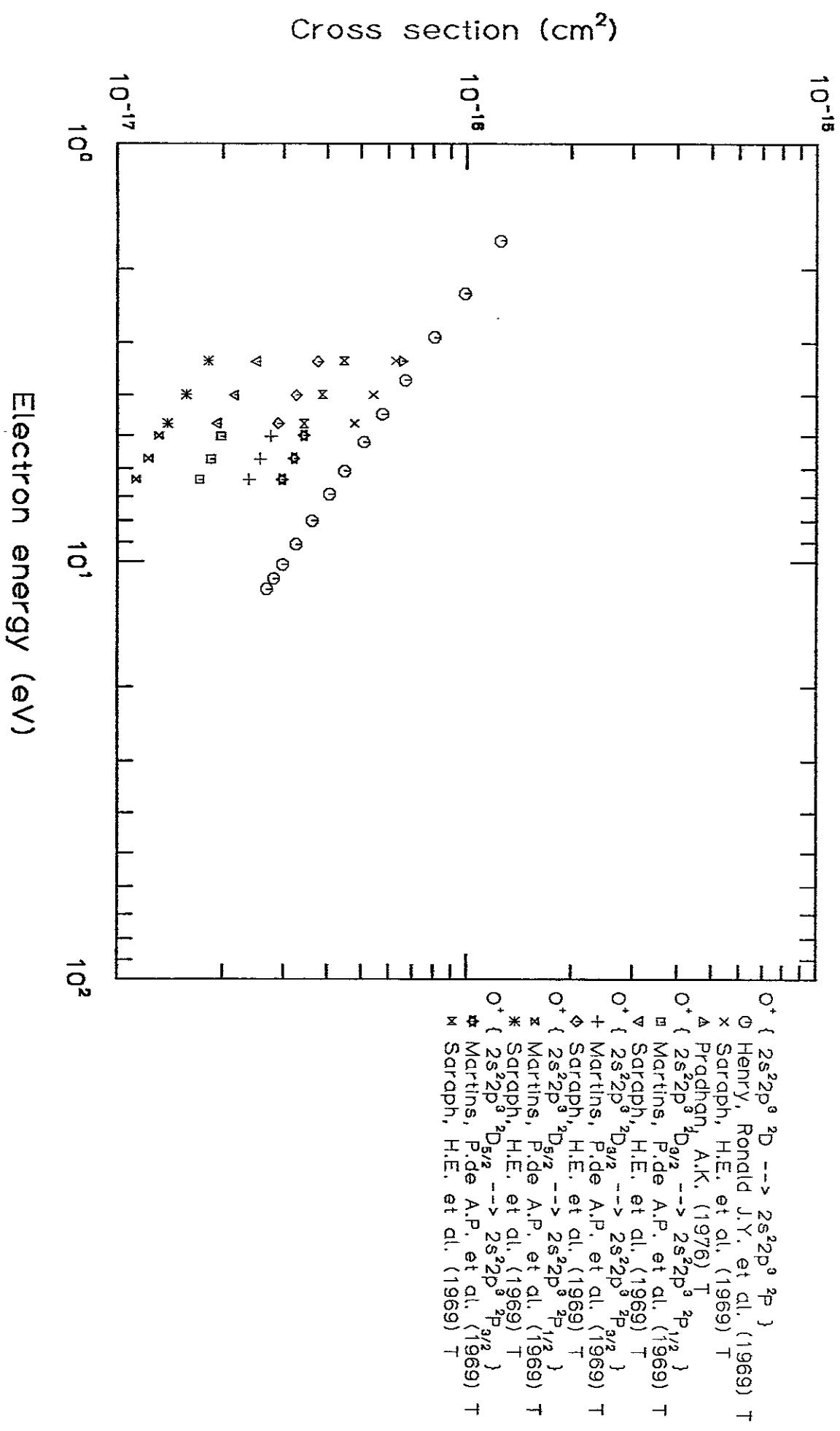
$O^+ \{ 2s^2 2p^3 \text{ } ^4S \rightarrow 2s 2p^4 \text{ } ^4P \}$

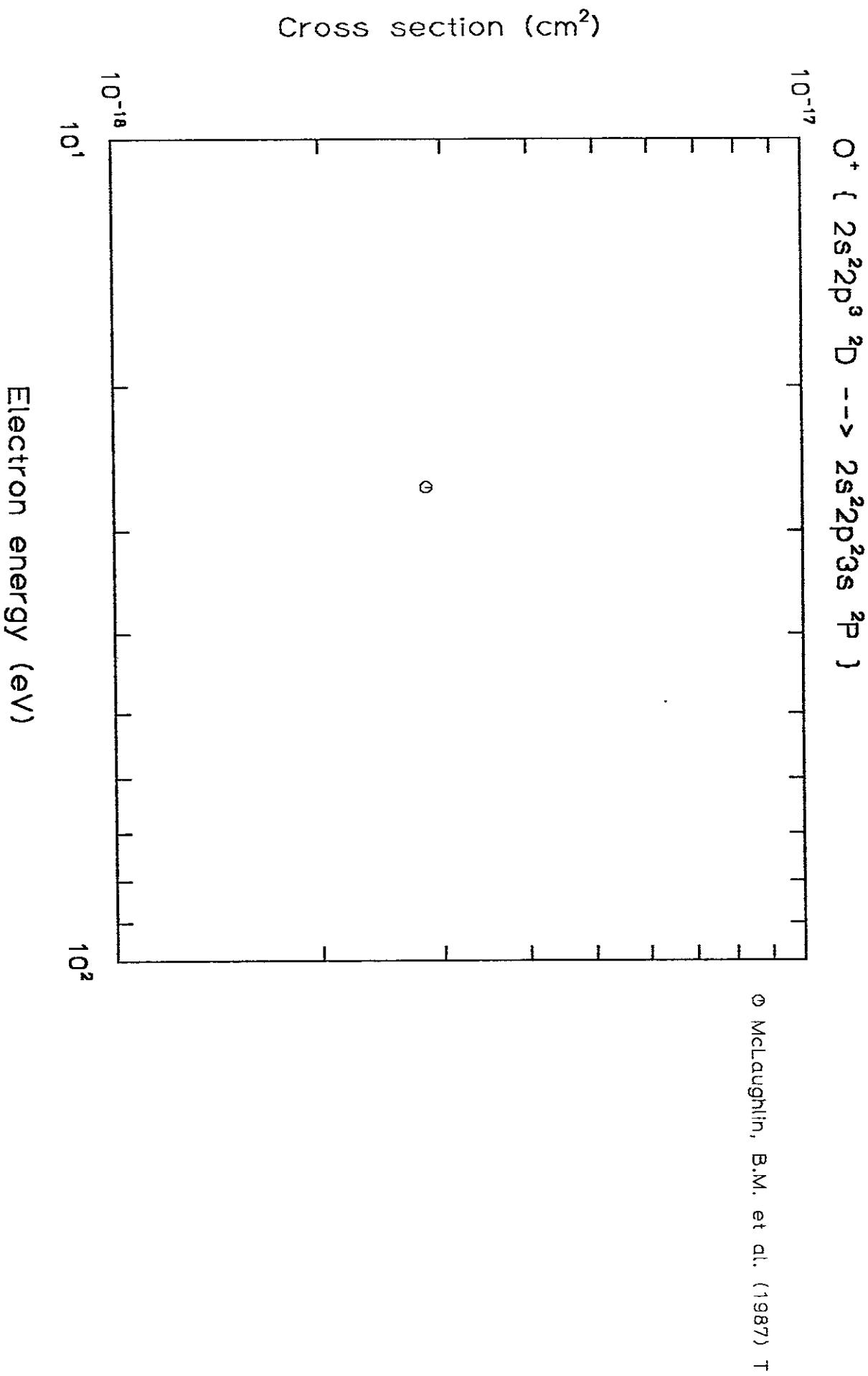


$O^+ \{ 2s^2 2p^3 \text{ } ^4S \rightarrow 2s^2 2p^2 3s \text{ } ^4P \}$



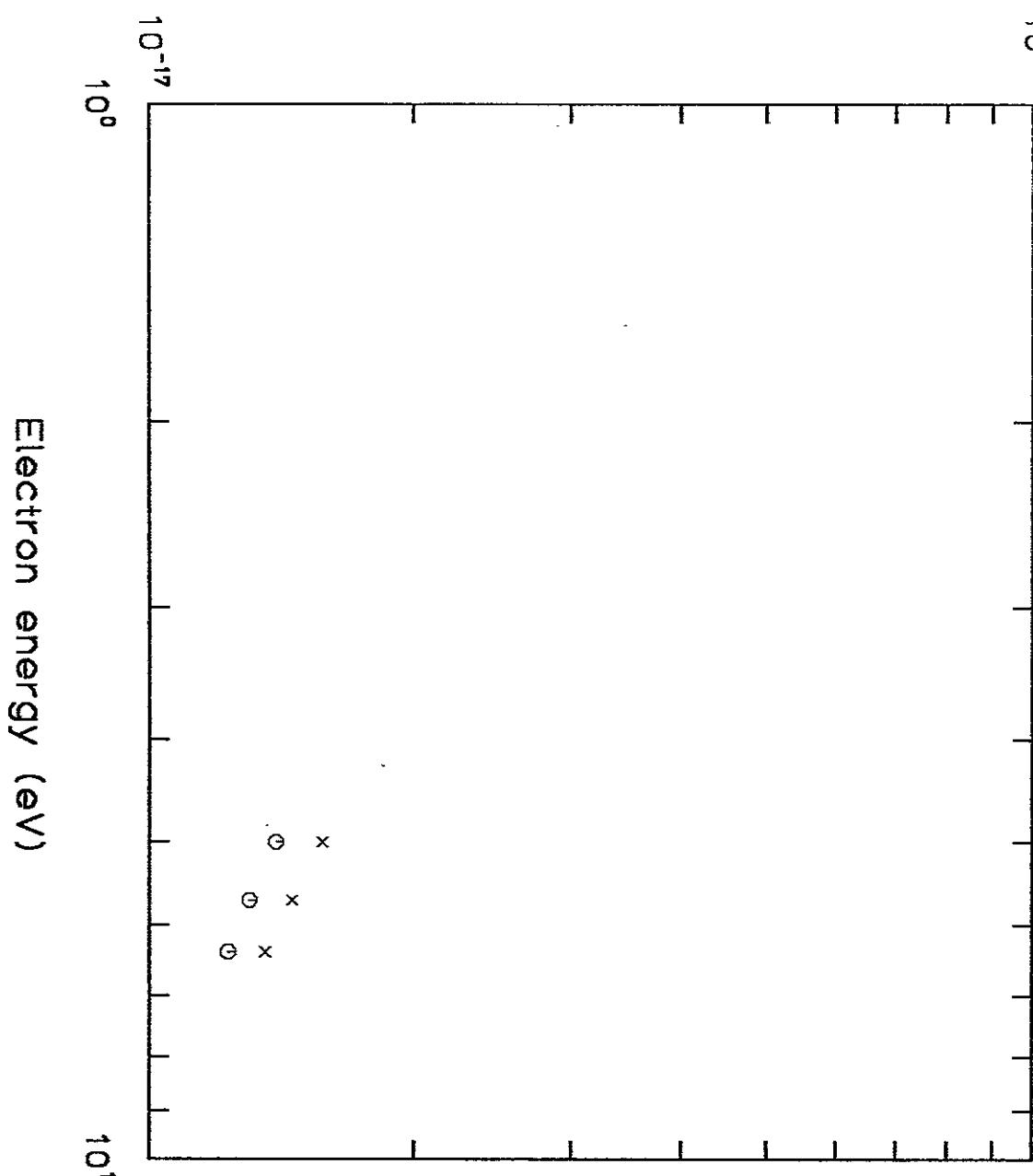




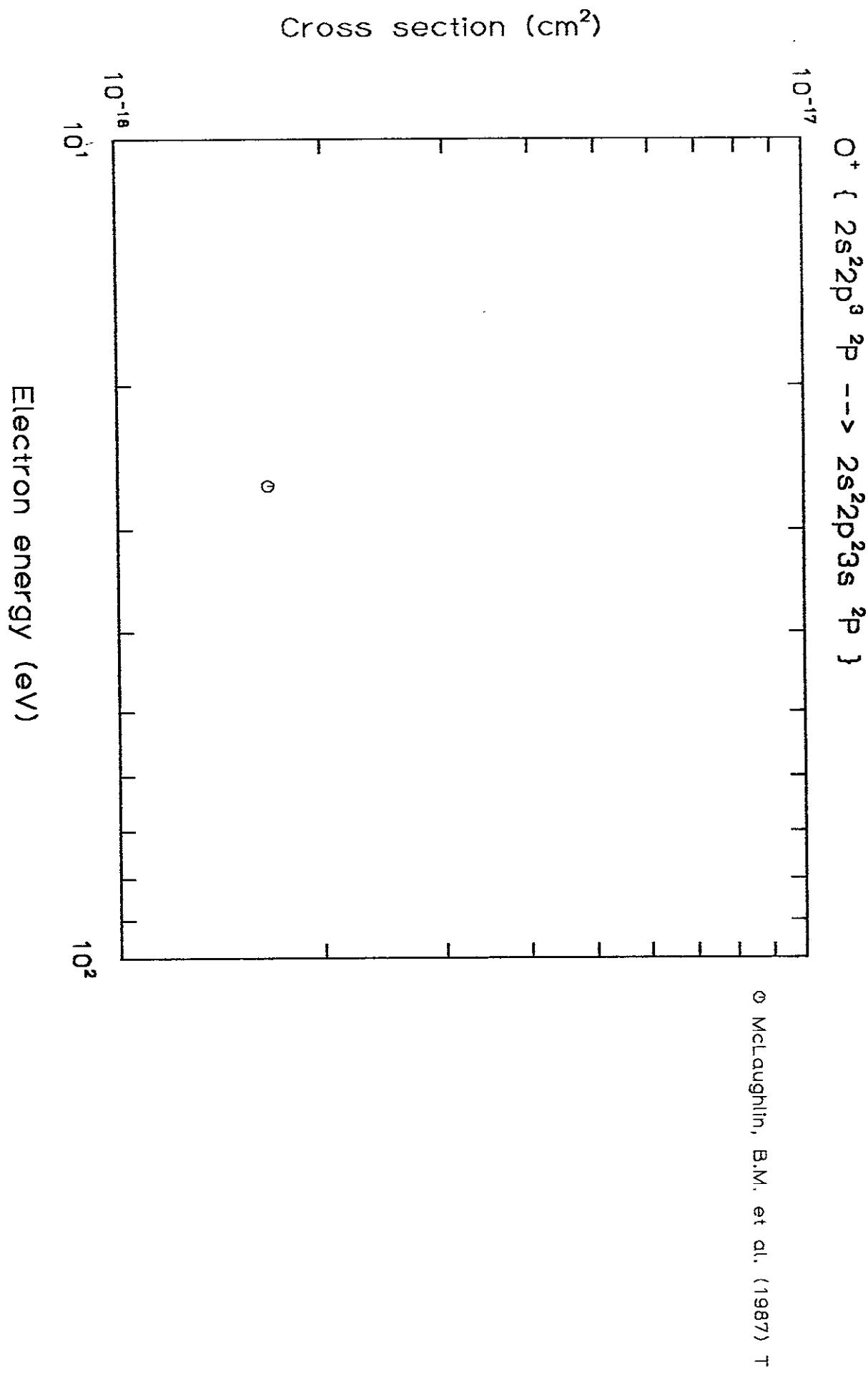


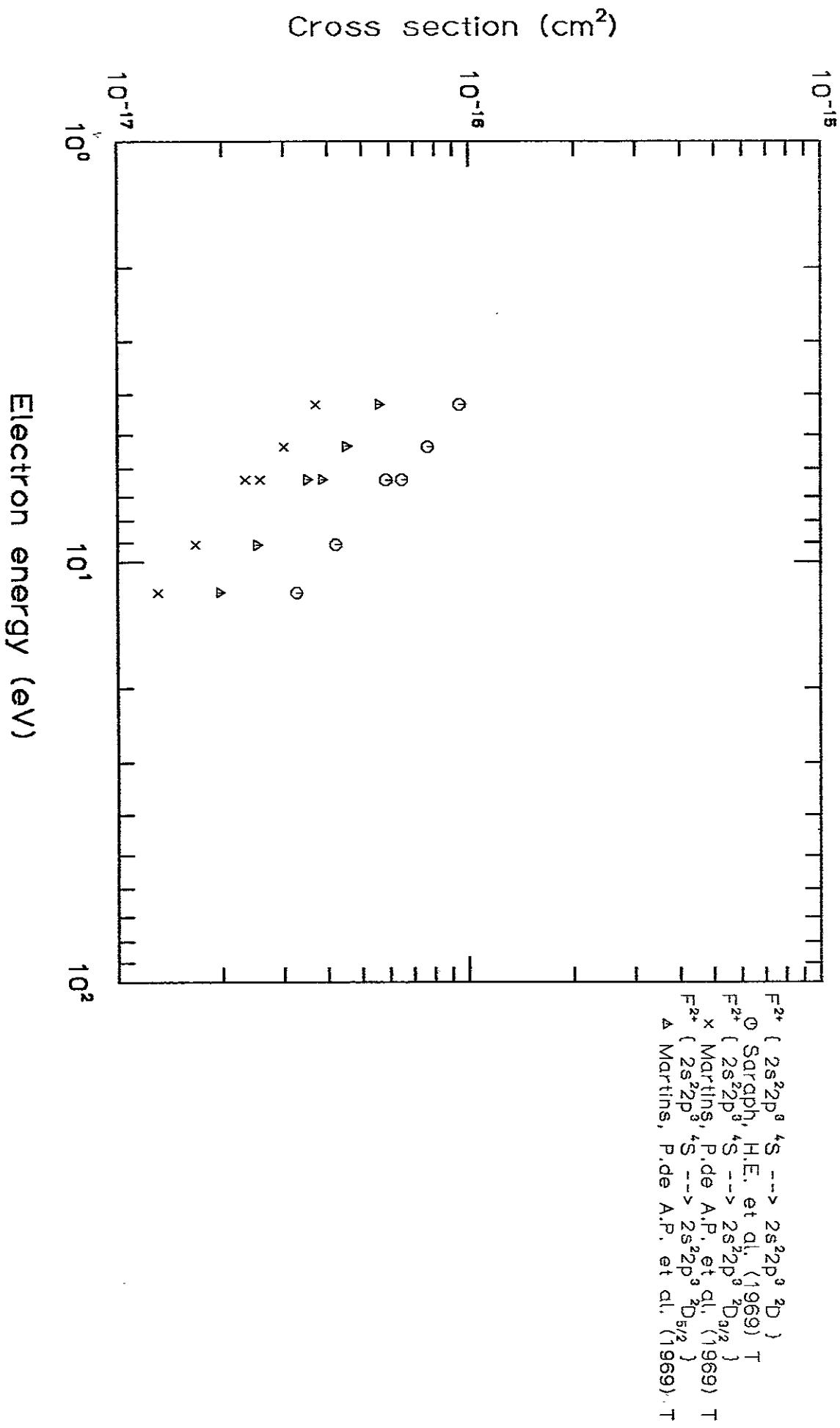
Cross section (cm^2)

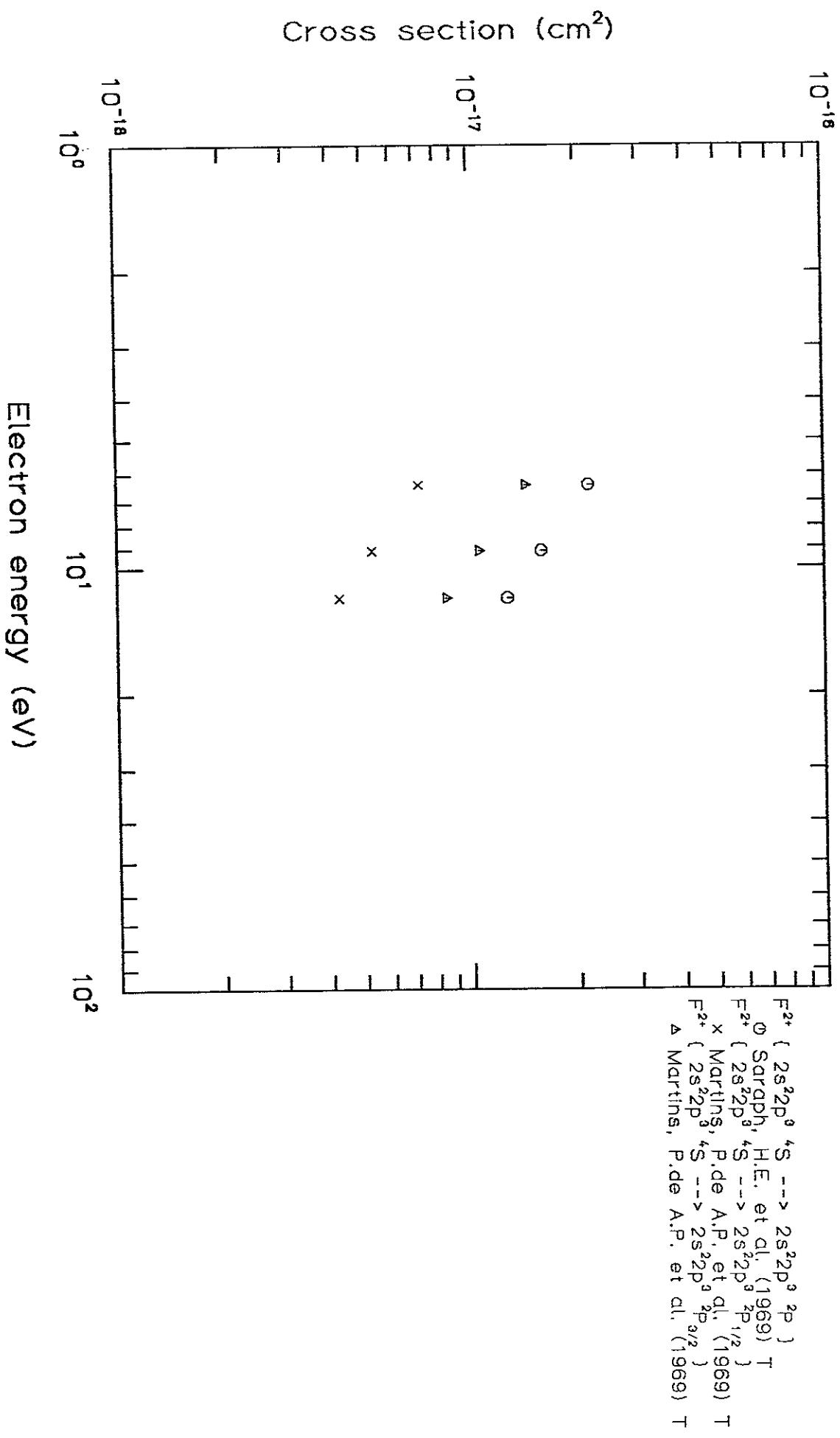
$O^+ \{ 2s^2 2p^3 \ 2P_{3/2} \rightarrow 2s^2 2p^3 \ 2P_{1/2} \ }$



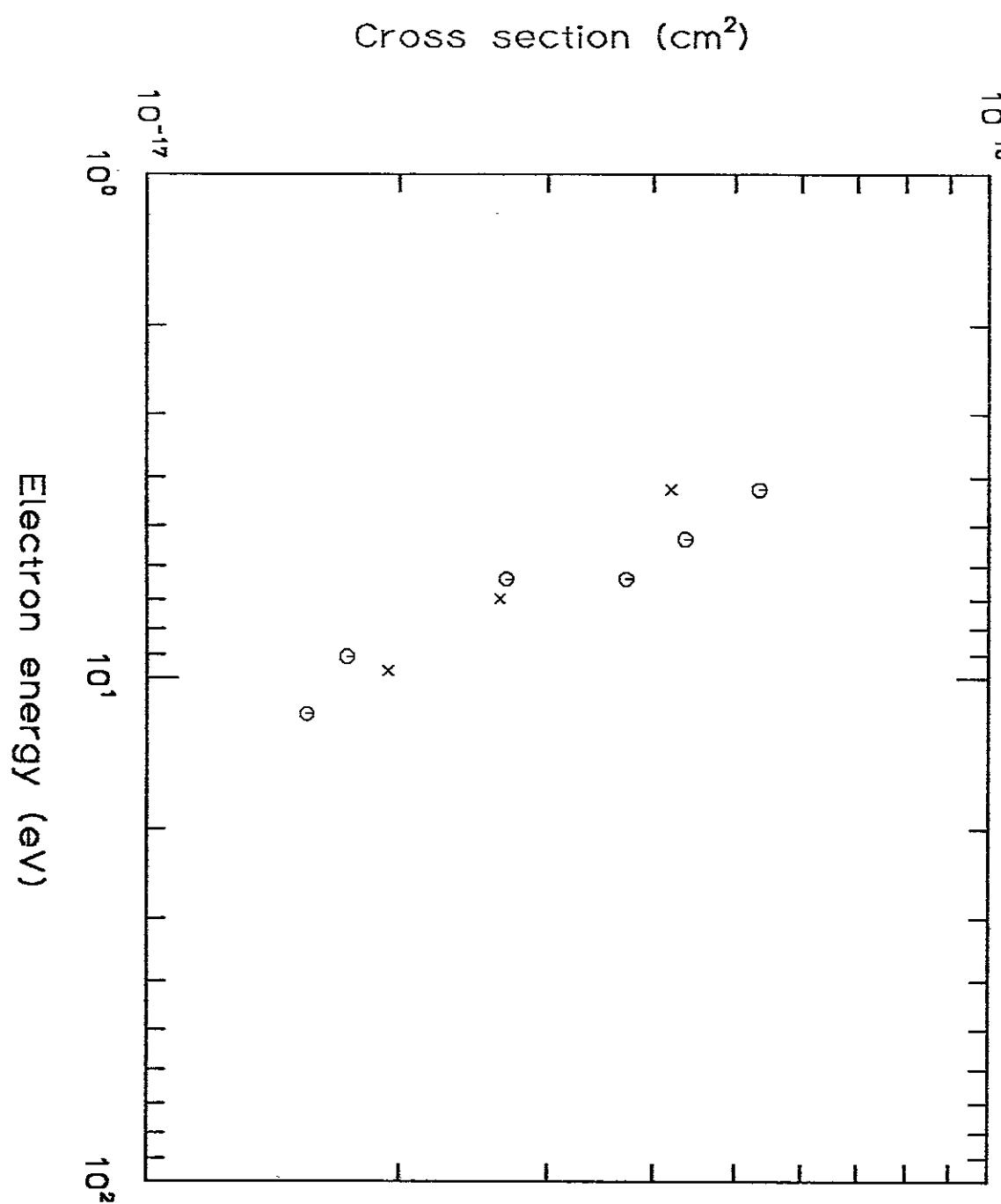
○ Martins, P.de A.P. et al. (1969) T
× Saraph, H.E. et al. (1969) T



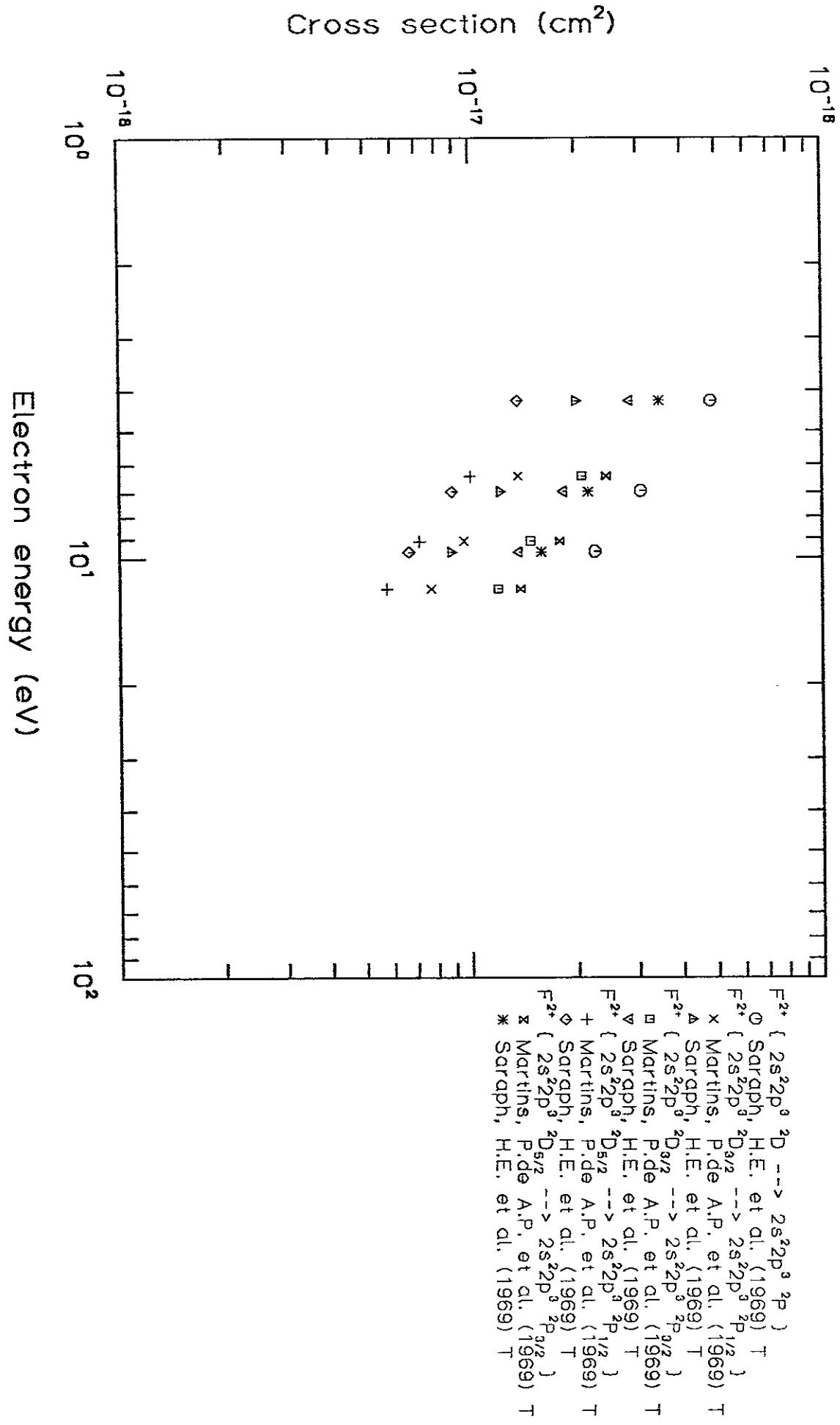


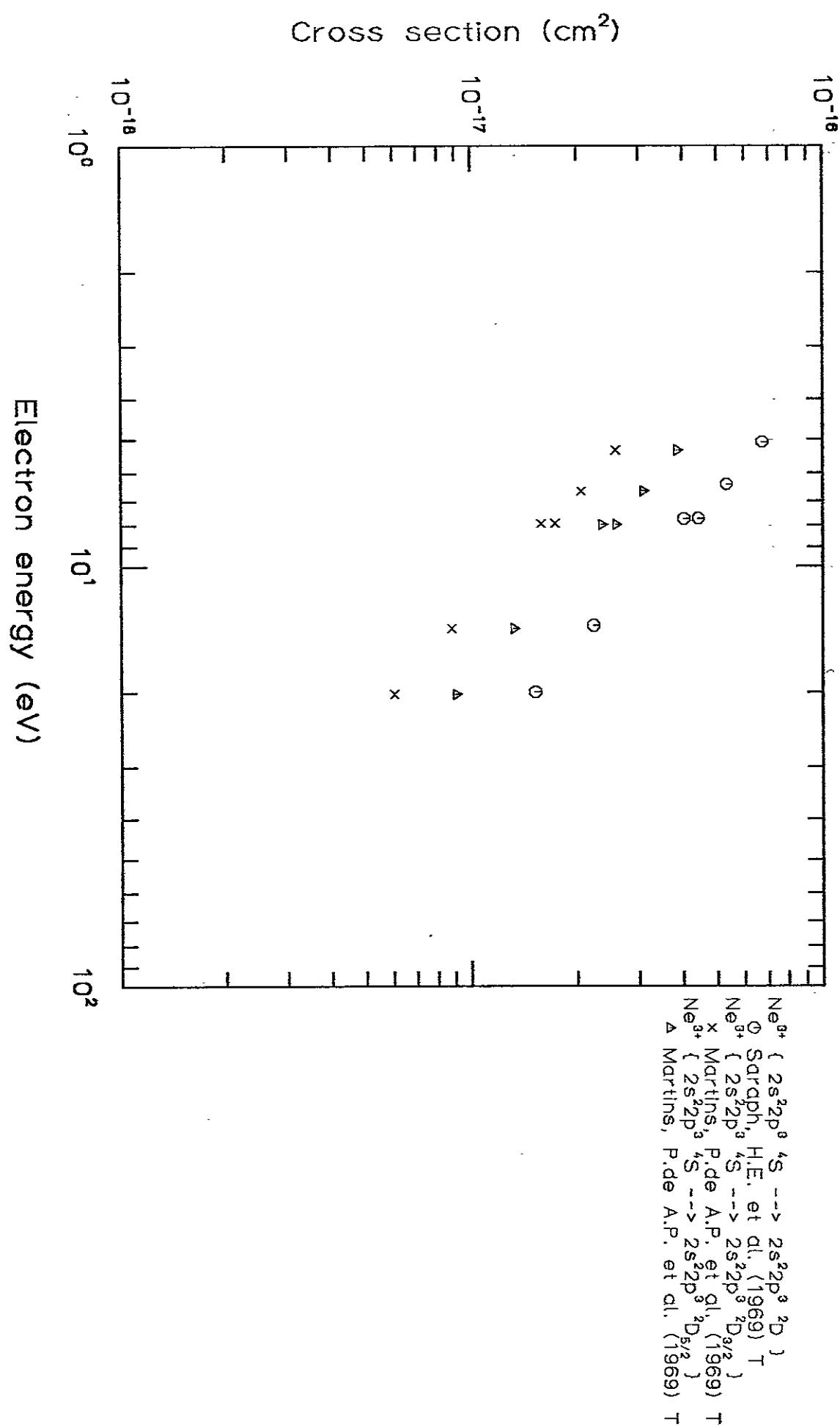


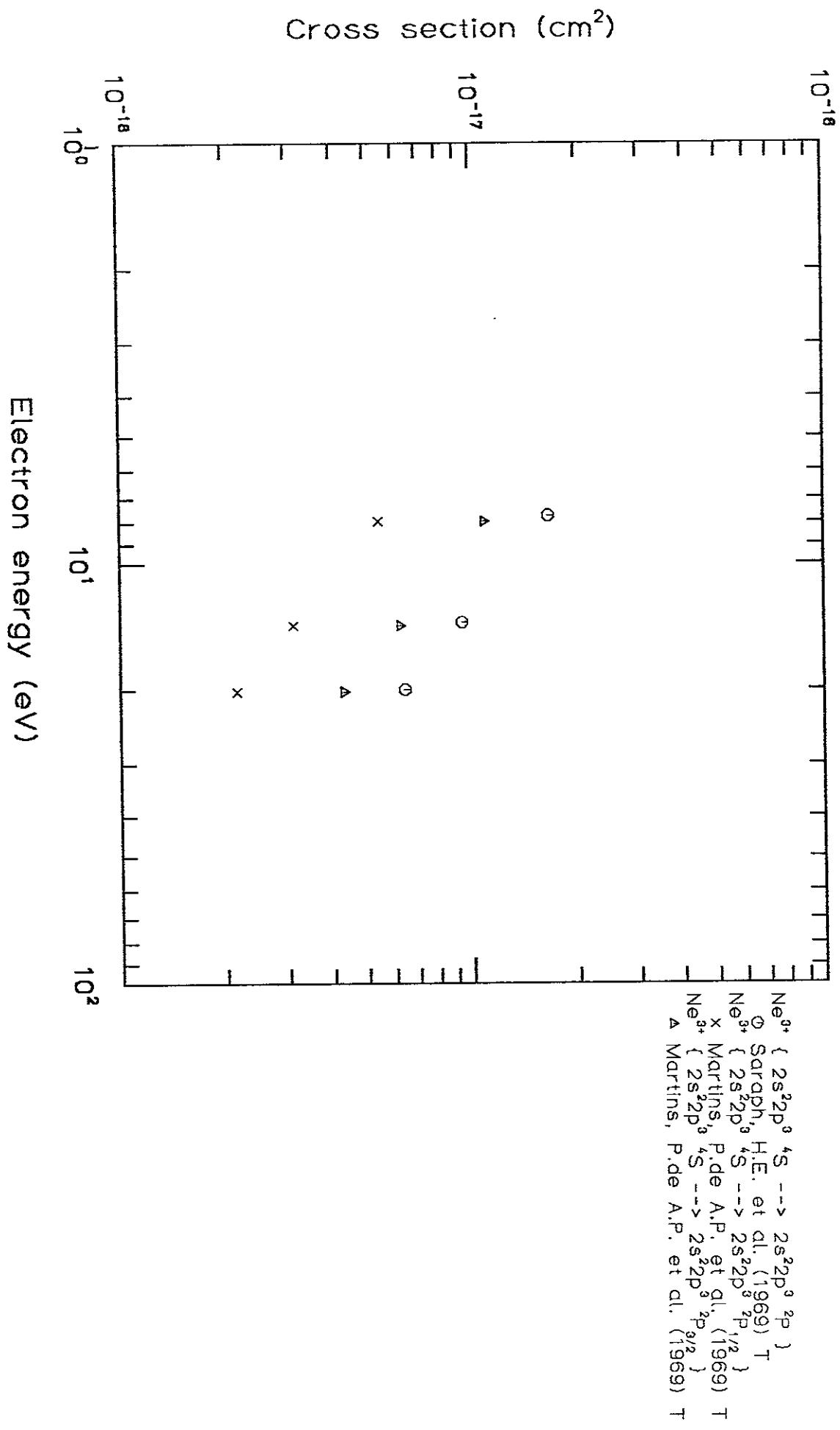
$F^{2+} \{ 2s^2 2p^3 \ ^2D_{5/2} \rightarrow 2s^2 2p^3 \ ^2D_{3/2} \ }$



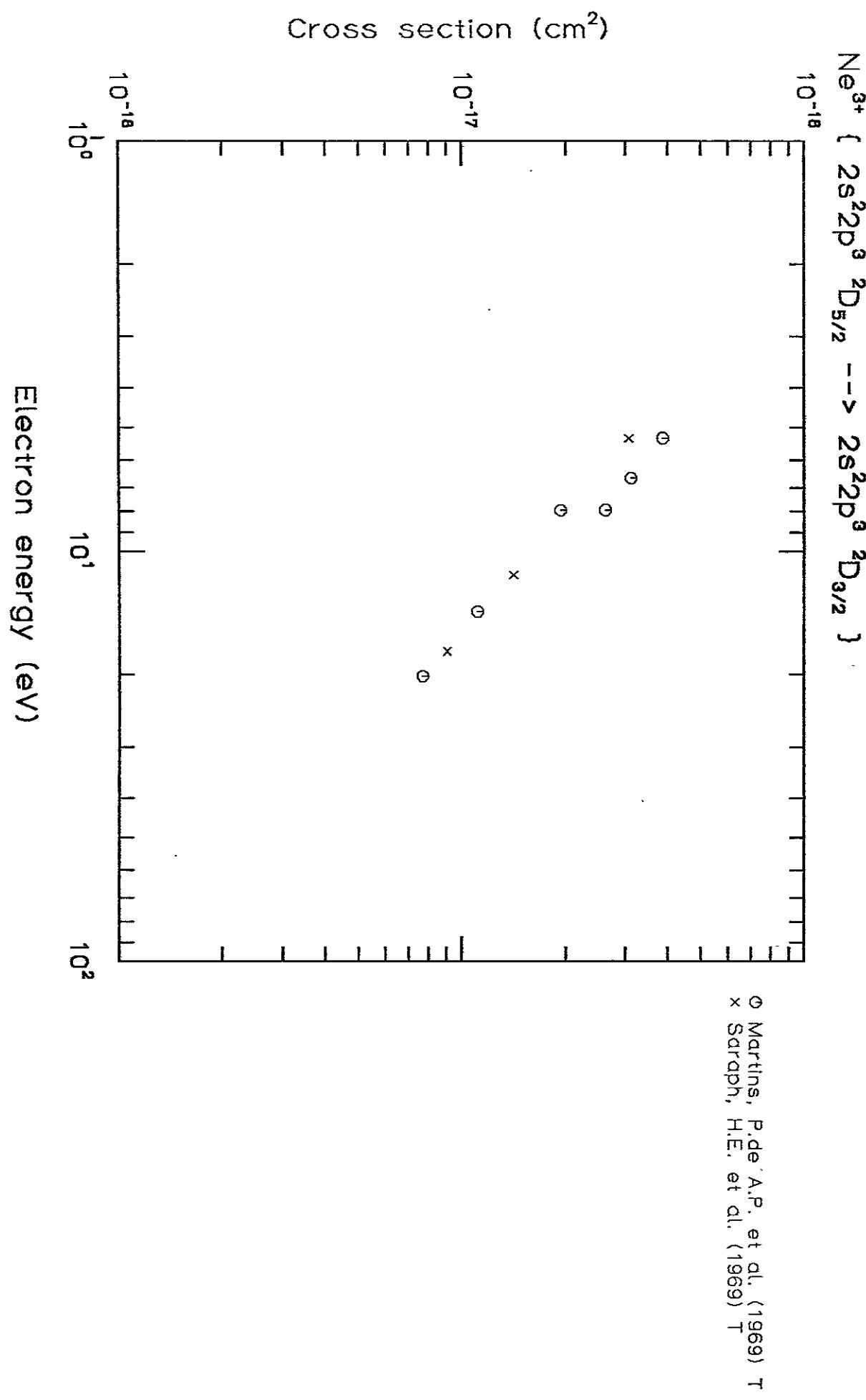
○ Martins, P.de A.P. et al. (1969)
× Saraph, H.E. et al. (1969)

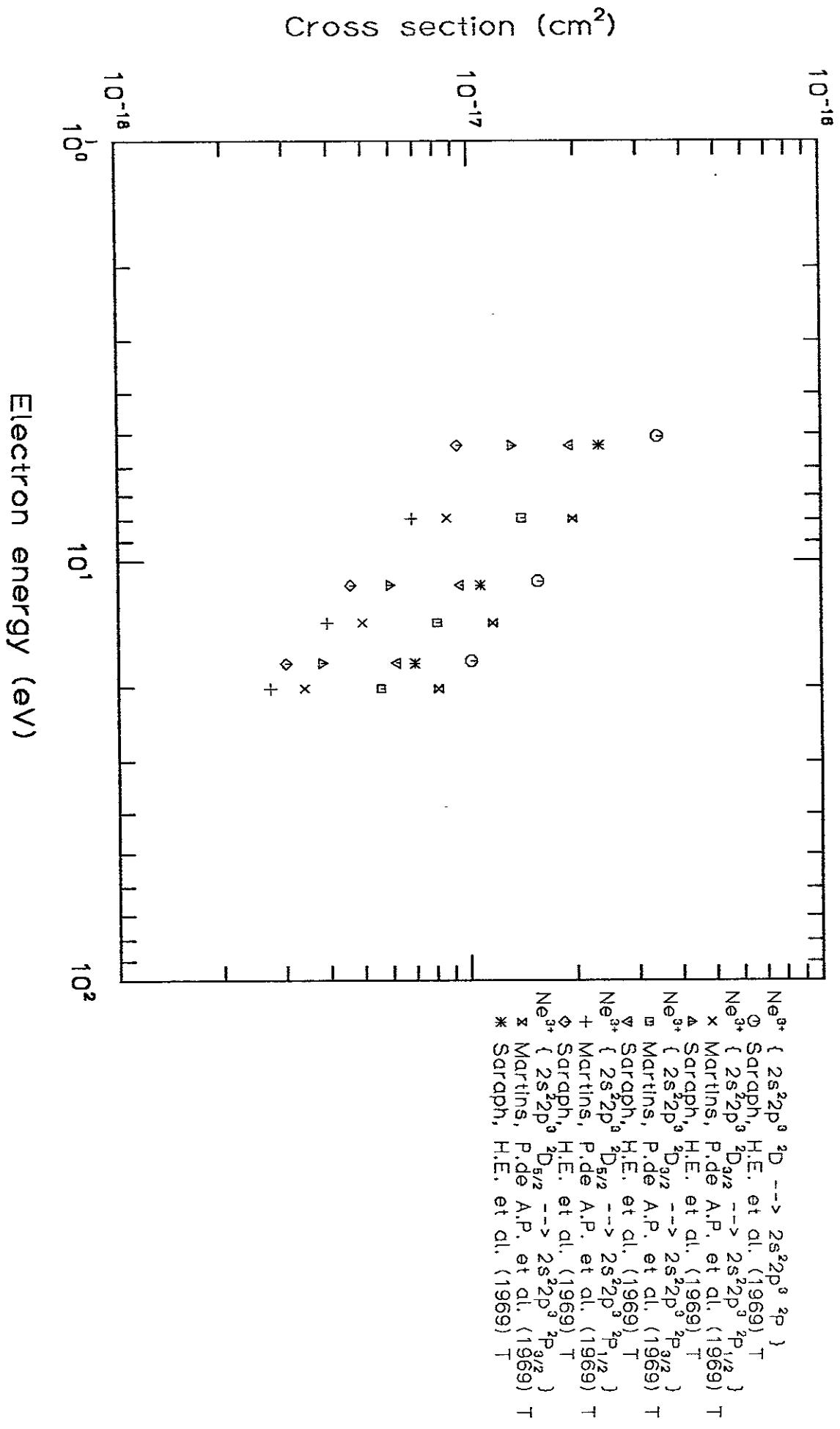




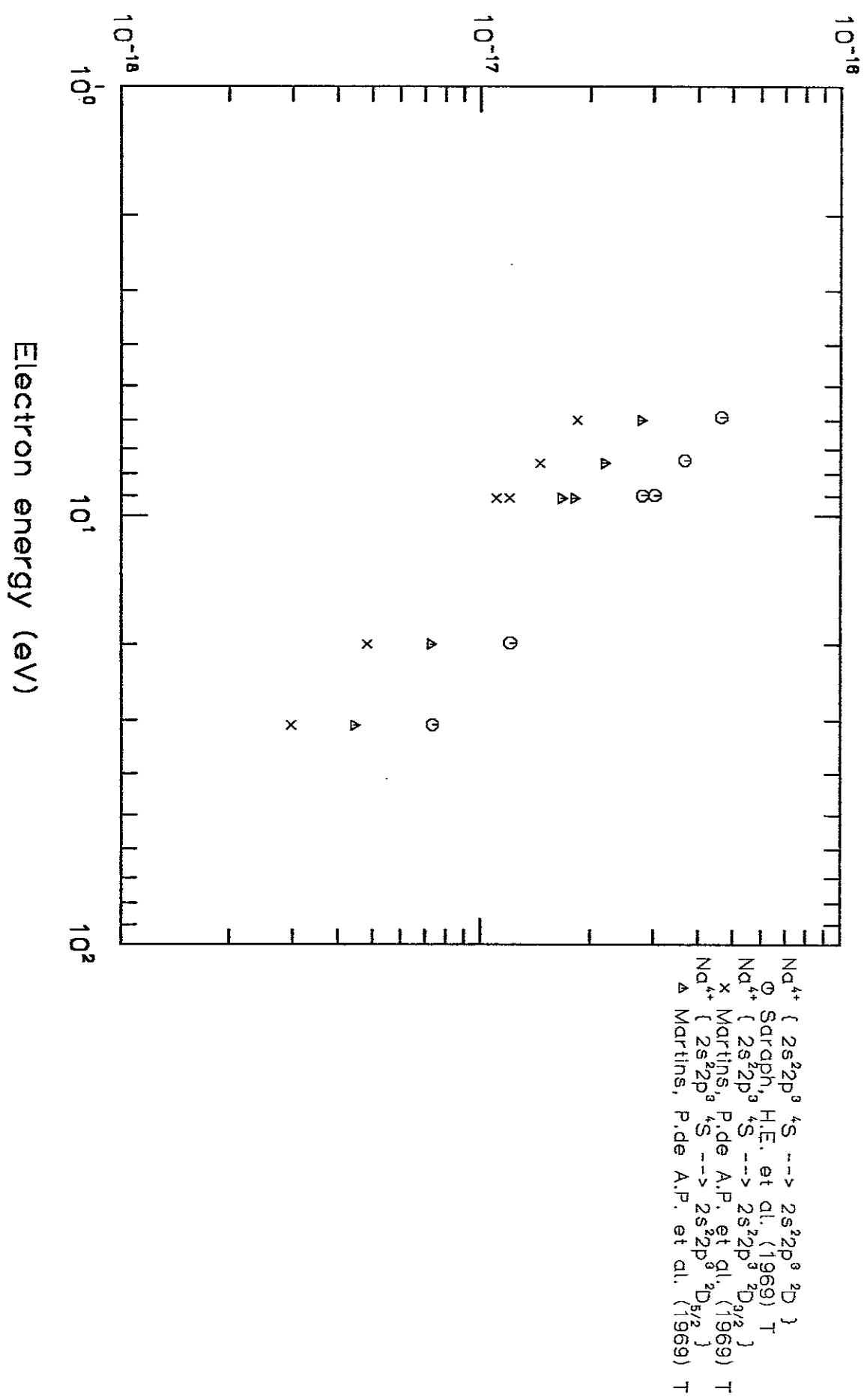


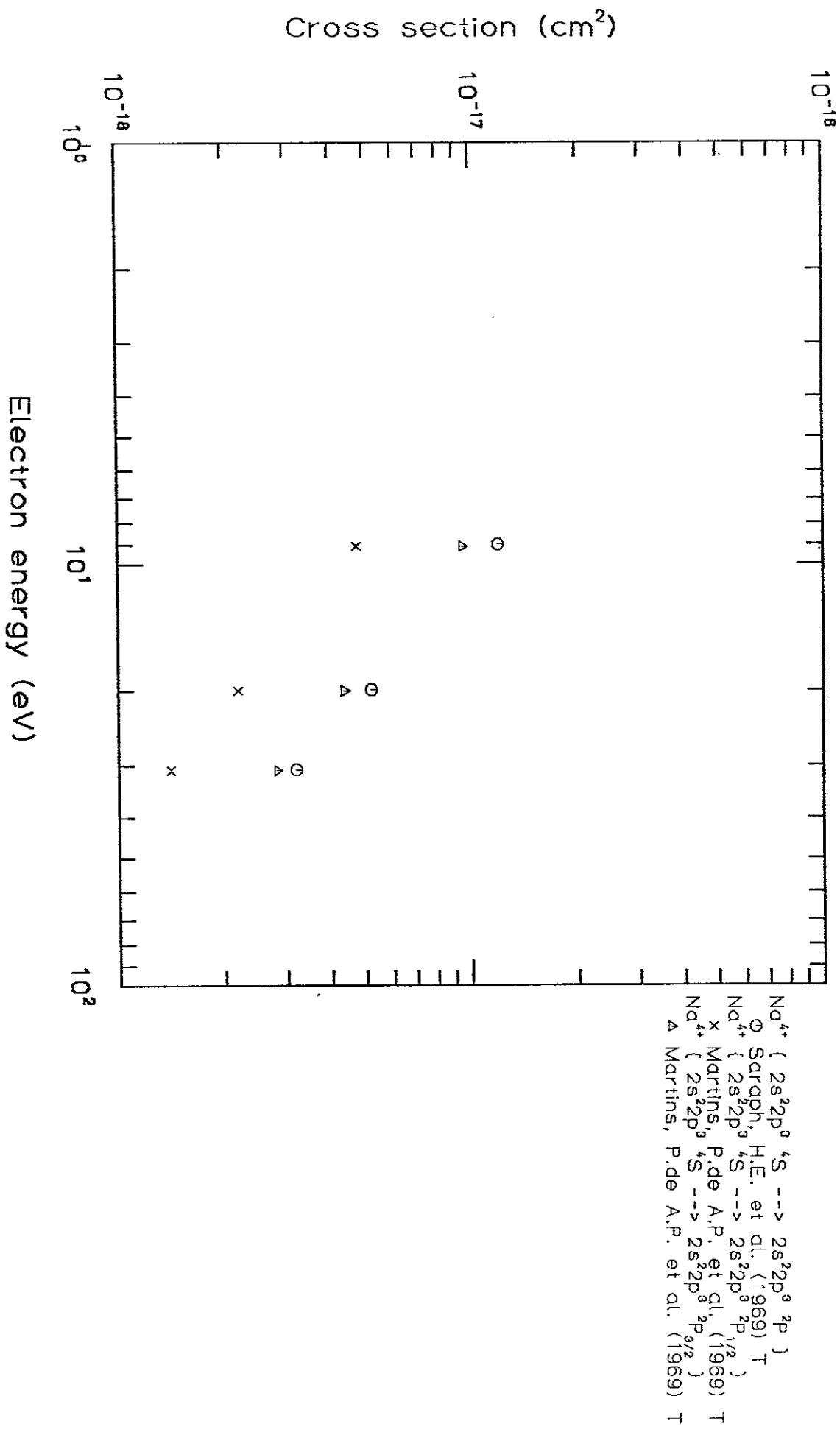
$N_{\Theta^{3+}} \{ 2s^2 2p^3 \ ^2D_{5/2} \rightarrow 2s^2 2p^3 \ ^2D_{3/2} \}$



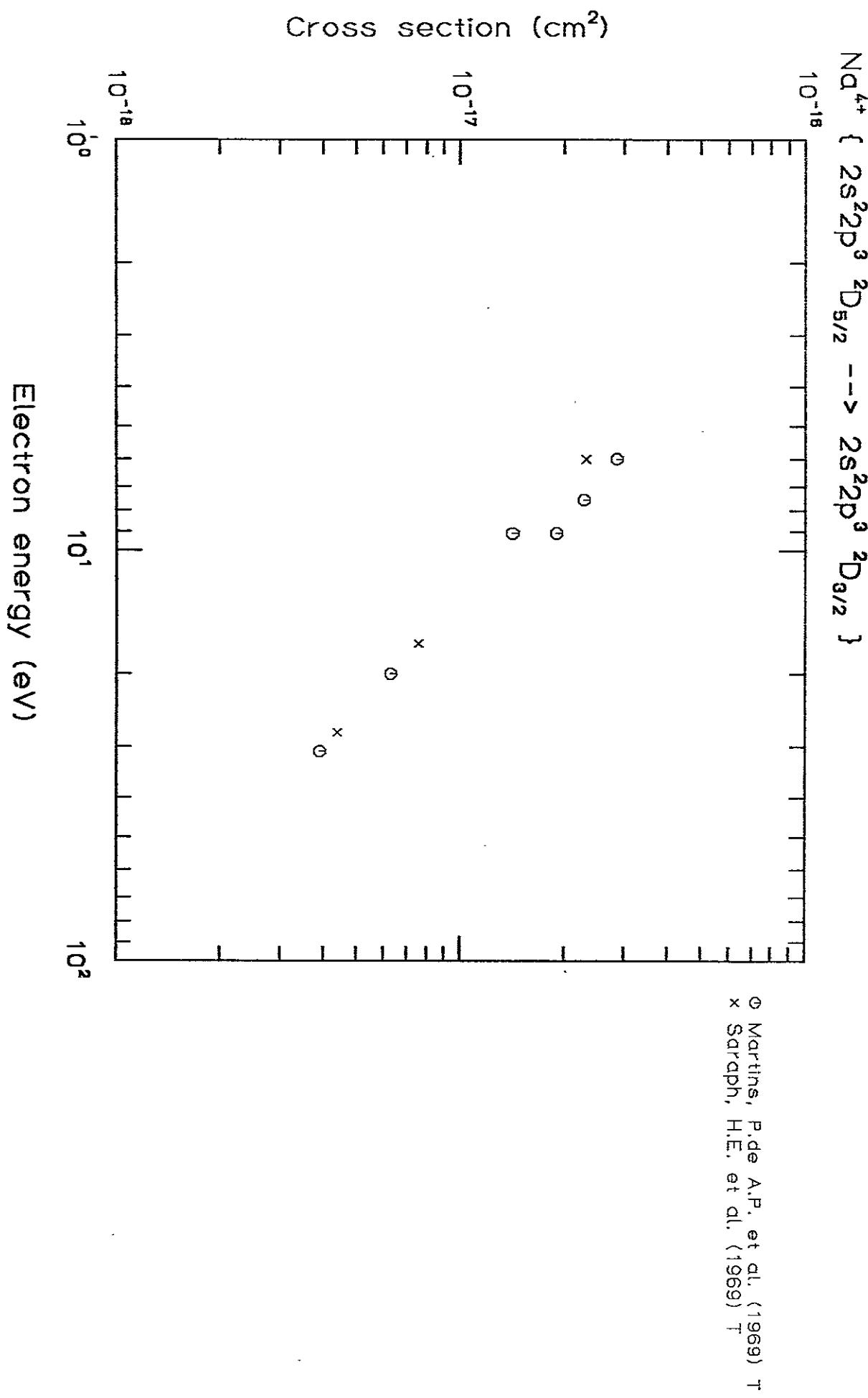


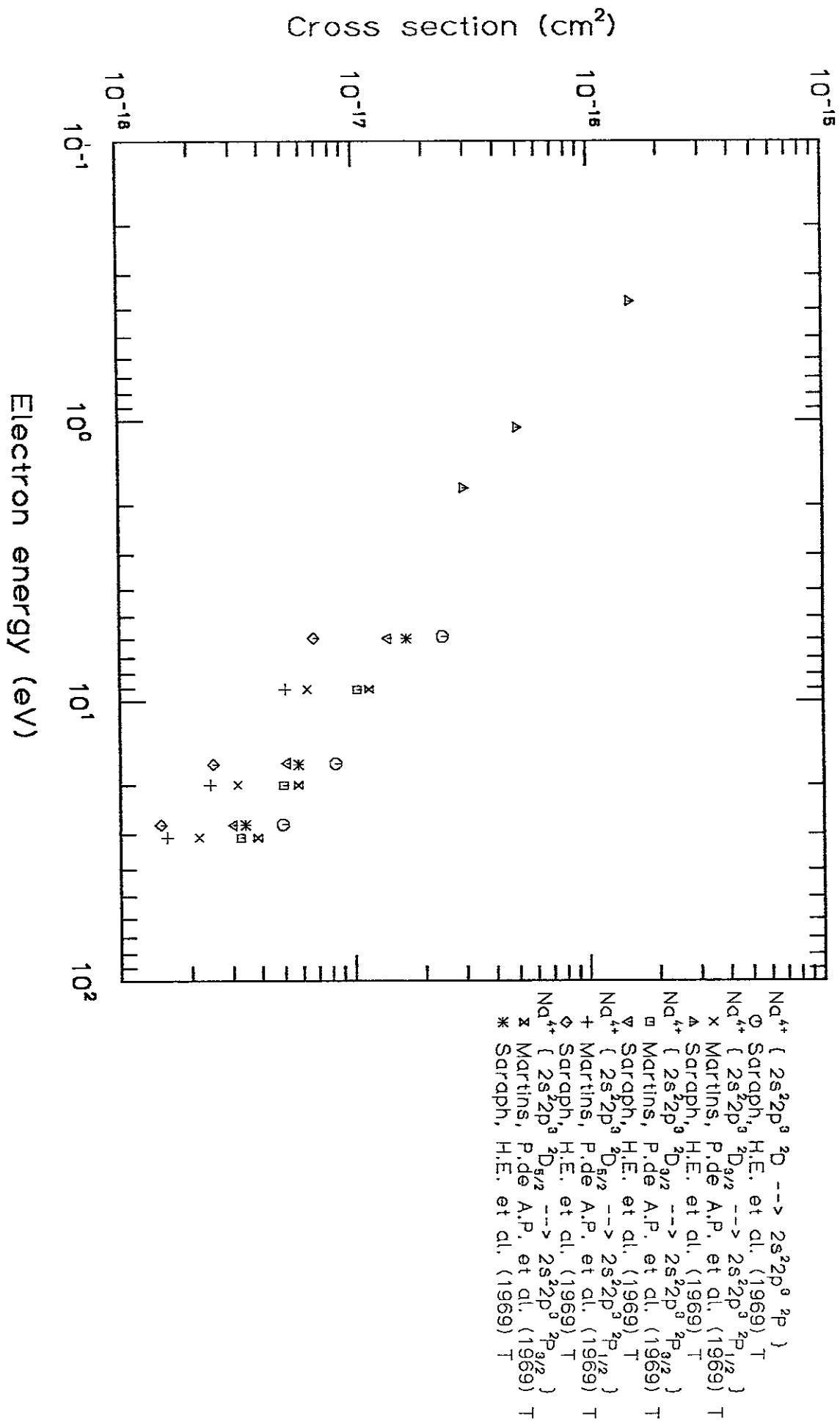
Cross section (cm^2)

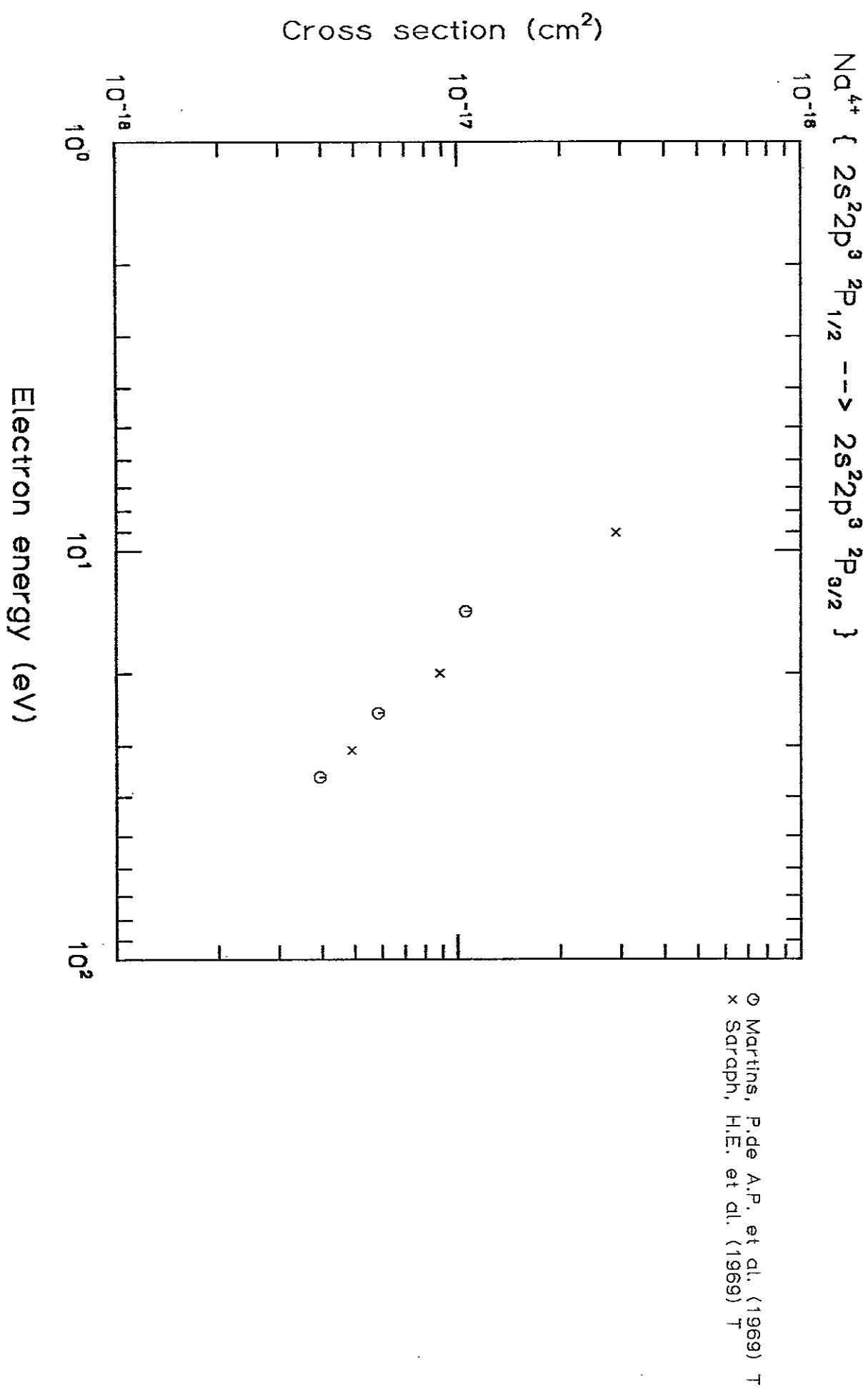




$\text{Na}^{4+} \{ 2s^2 2p^3 \ 2D_{5/2} \rightarrow 2s^2 2p^3 \ 2D_{3/2} \ }$







$Mg^{5+} \{ 2s^2 2p^3 \text{ } ^4S \rightarrow 2s^2 2p^3 \text{ } ^2D \}$

10^{-16}

Cross section (cm^2)

10^{-17}

10^0

10^1

10^2

Electron energy (eV)

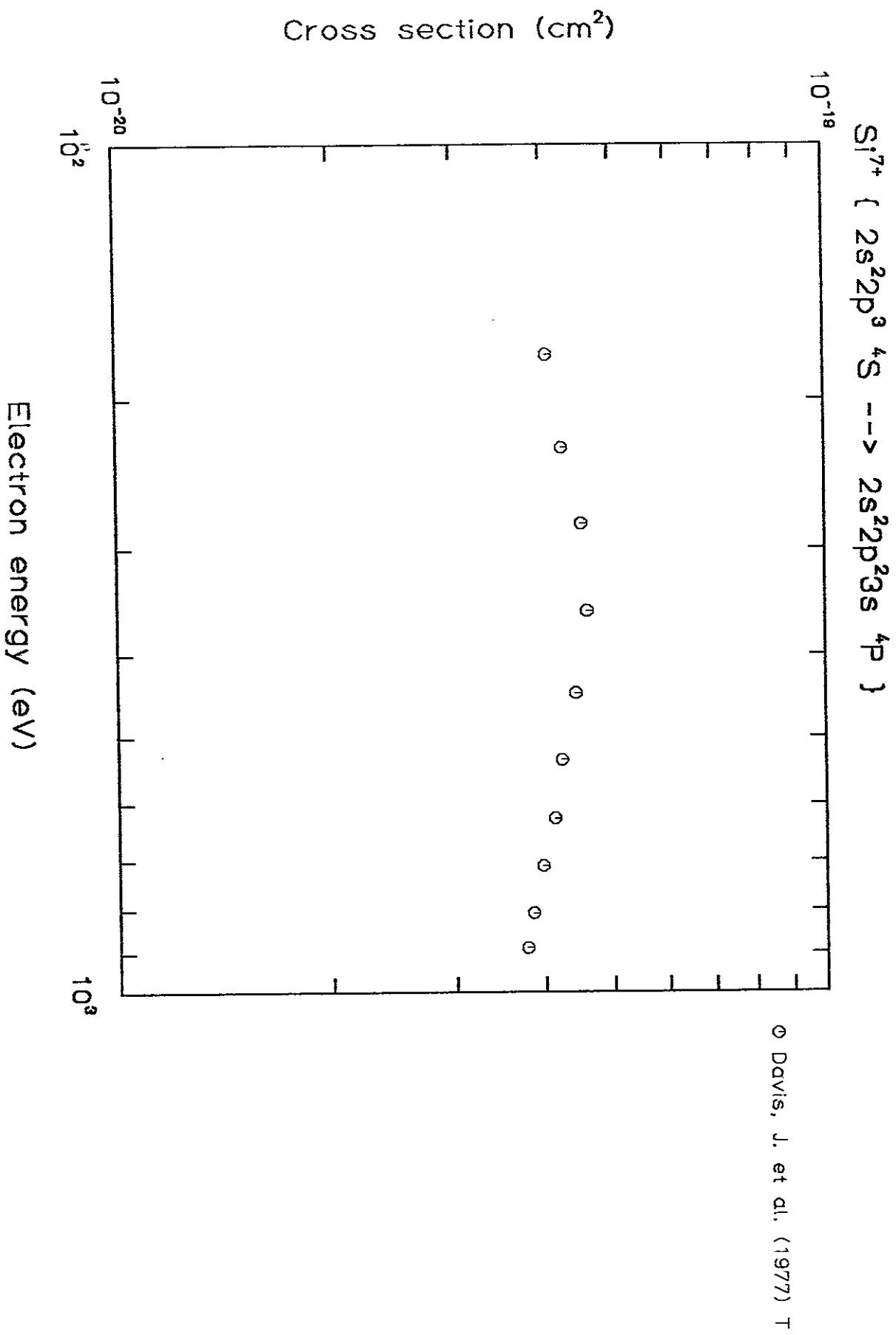
○ Saraph, H.E. et al. (1969) T

○
○
○
○
○

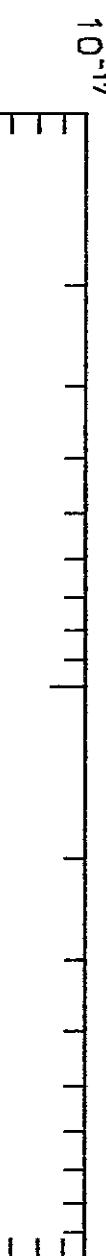
$Mg^{5+} \{ 2s^2 2p^3 \text{ } ^4S \rightarrow 2s^2 2p^3 \text{ } ^2P \}$



Θ Saraph, H.E. et al. (1969) T



$\text{Si}^{7+} \{ 2s^2 2p^3 \text{ } ^4S \rightarrow 2s^2 2p^2 3d \text{ } ^4P \}$



○ Davis, J. et al. (1977) T

$Mg^{5+} \{ 2s^2 2p^3 \text{ } ^2D \rightarrow 2s^2 2p^3 \text{ } ^2P \}$

10^{-18}

○ Saraph, H.E. et al. (1969) T

Cross section (cm^2)

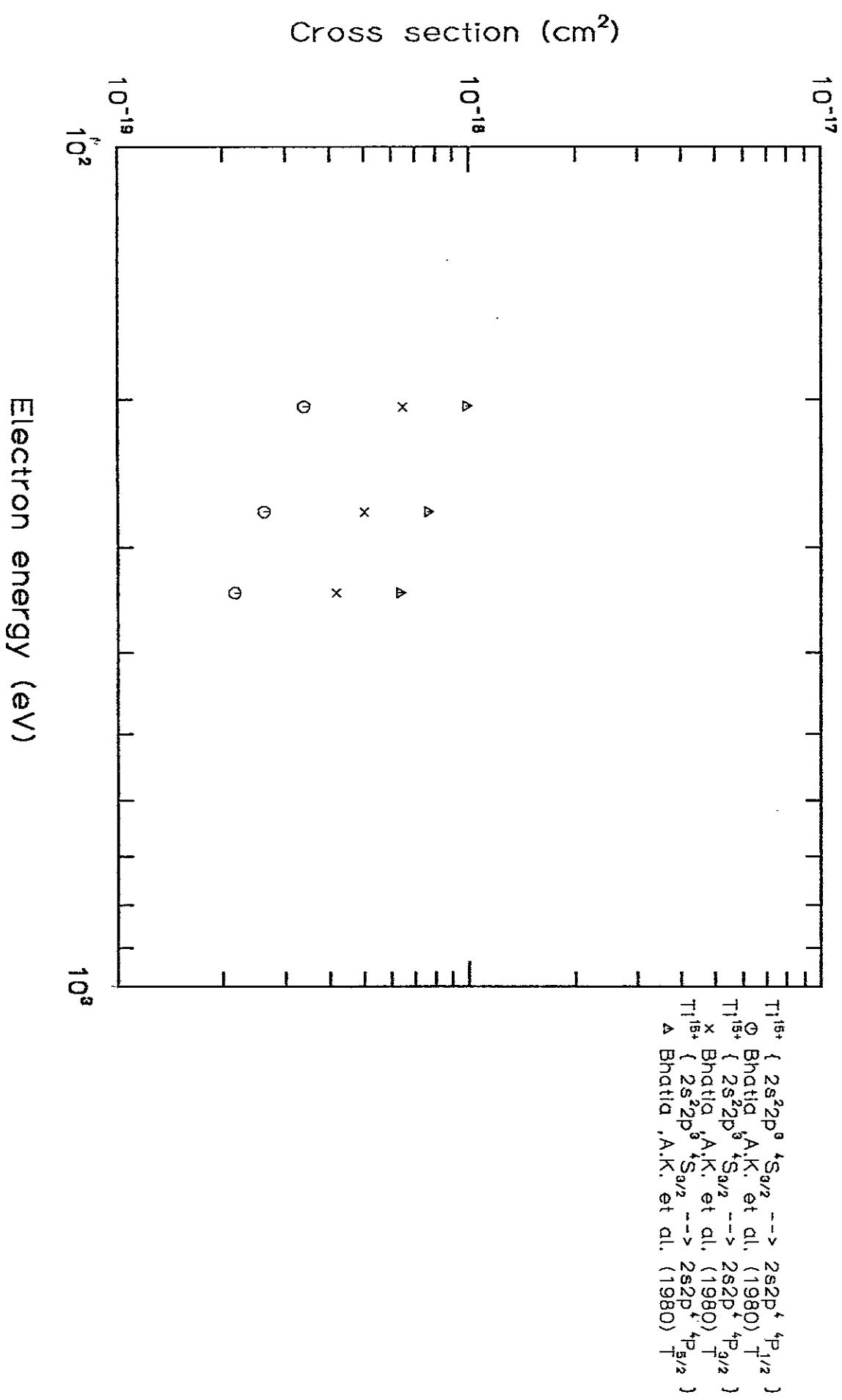
10^{-17}

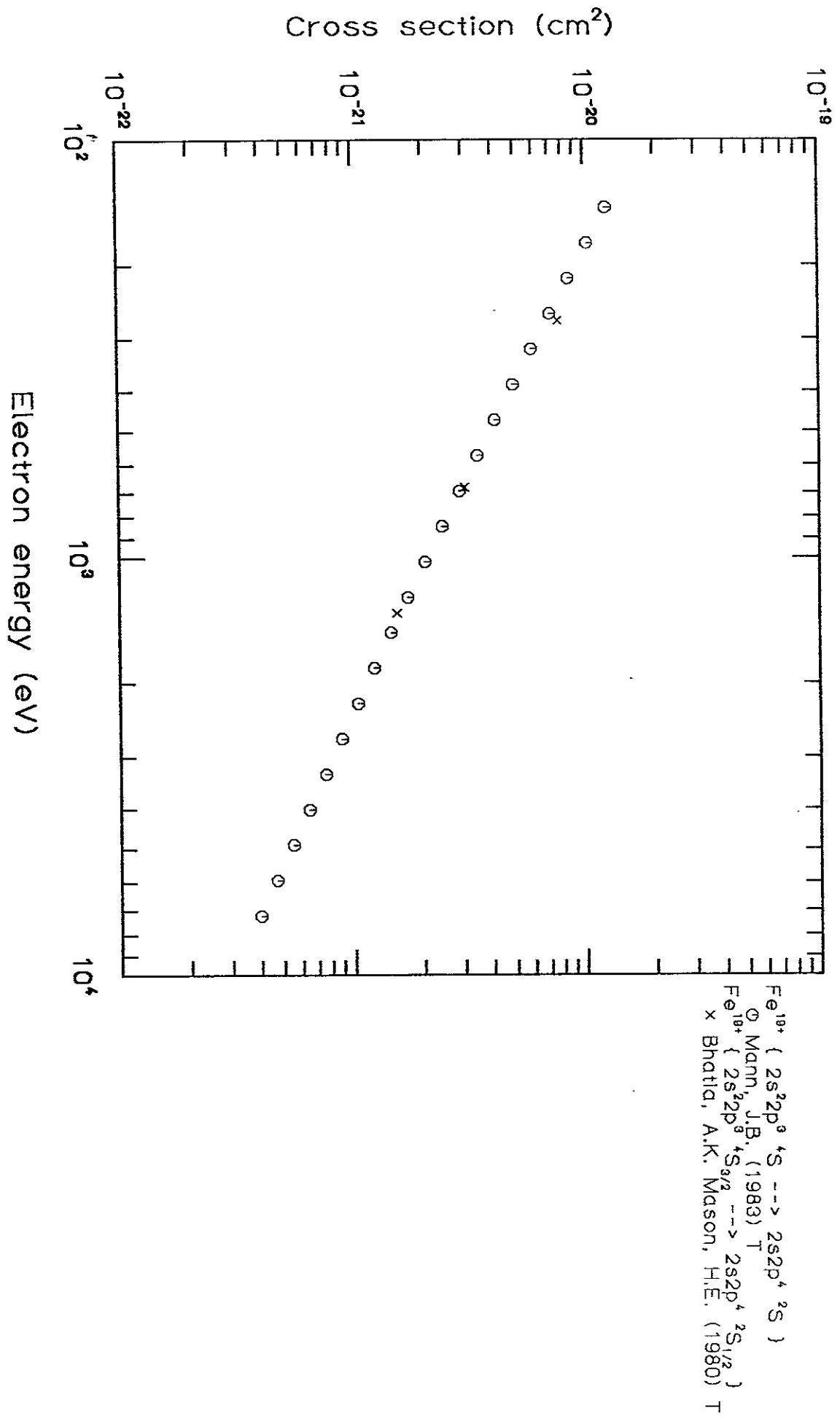
10^{-18}

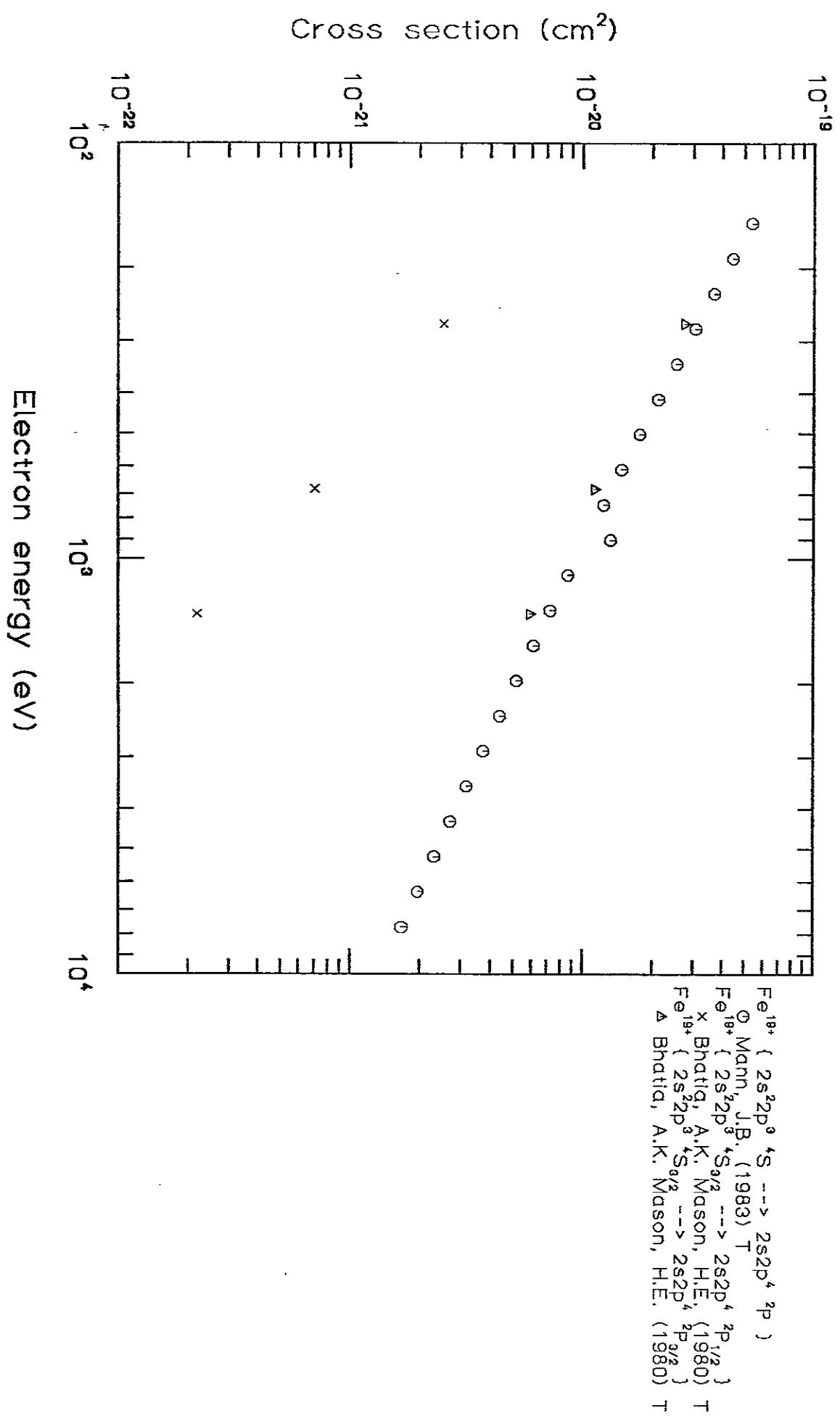
10^{-19}

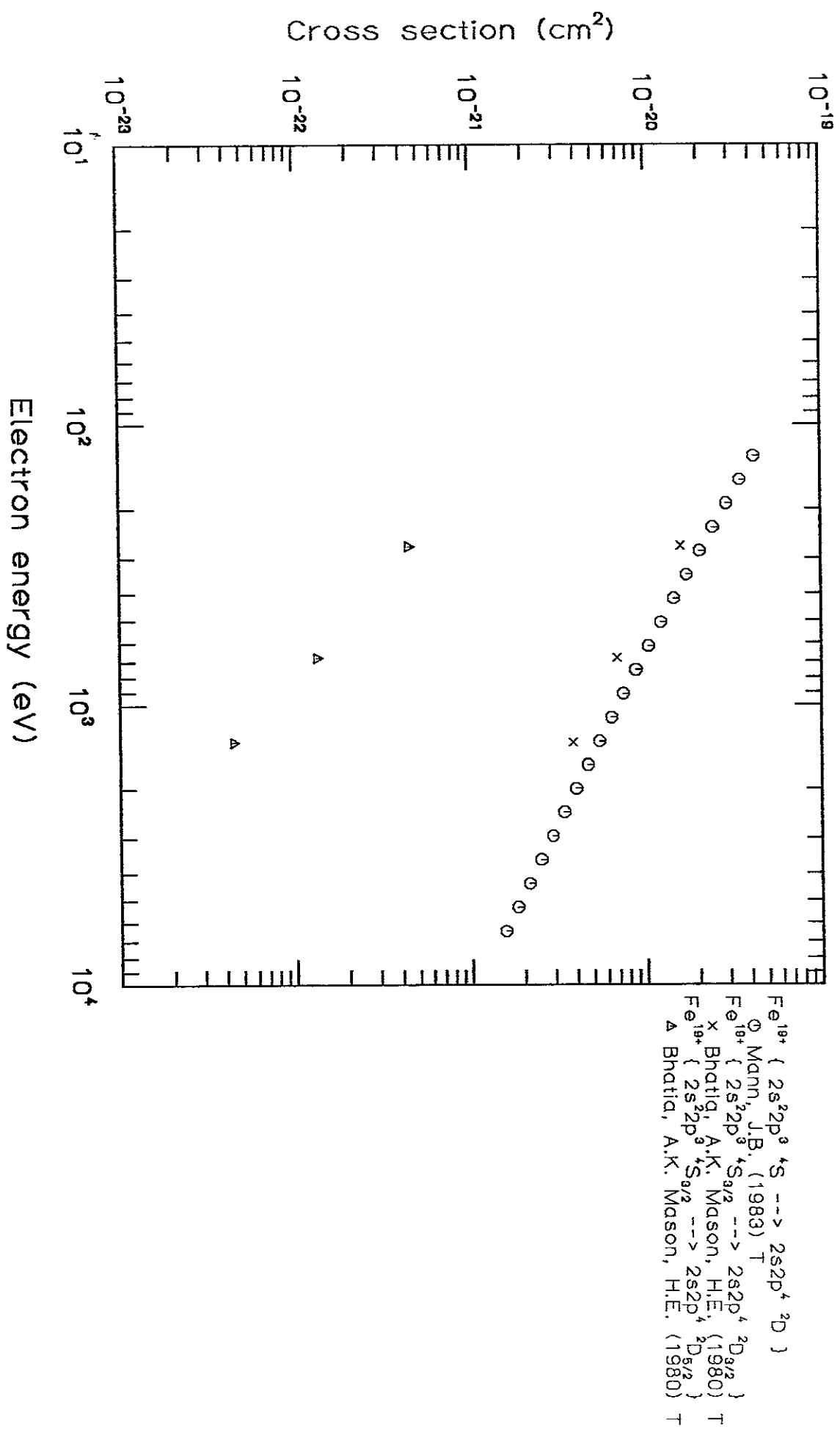
10^{-20}

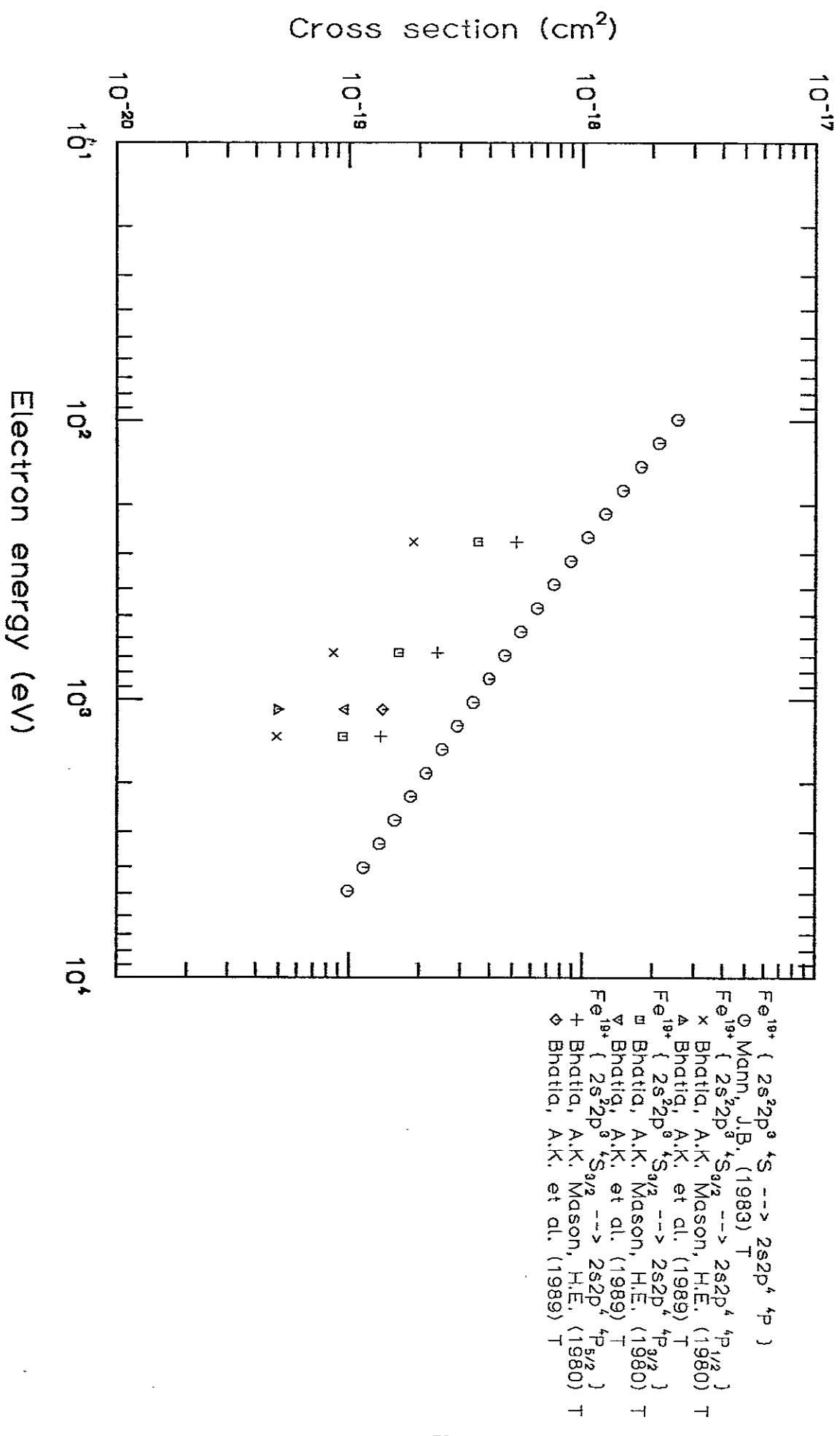
Electron energy (eV)



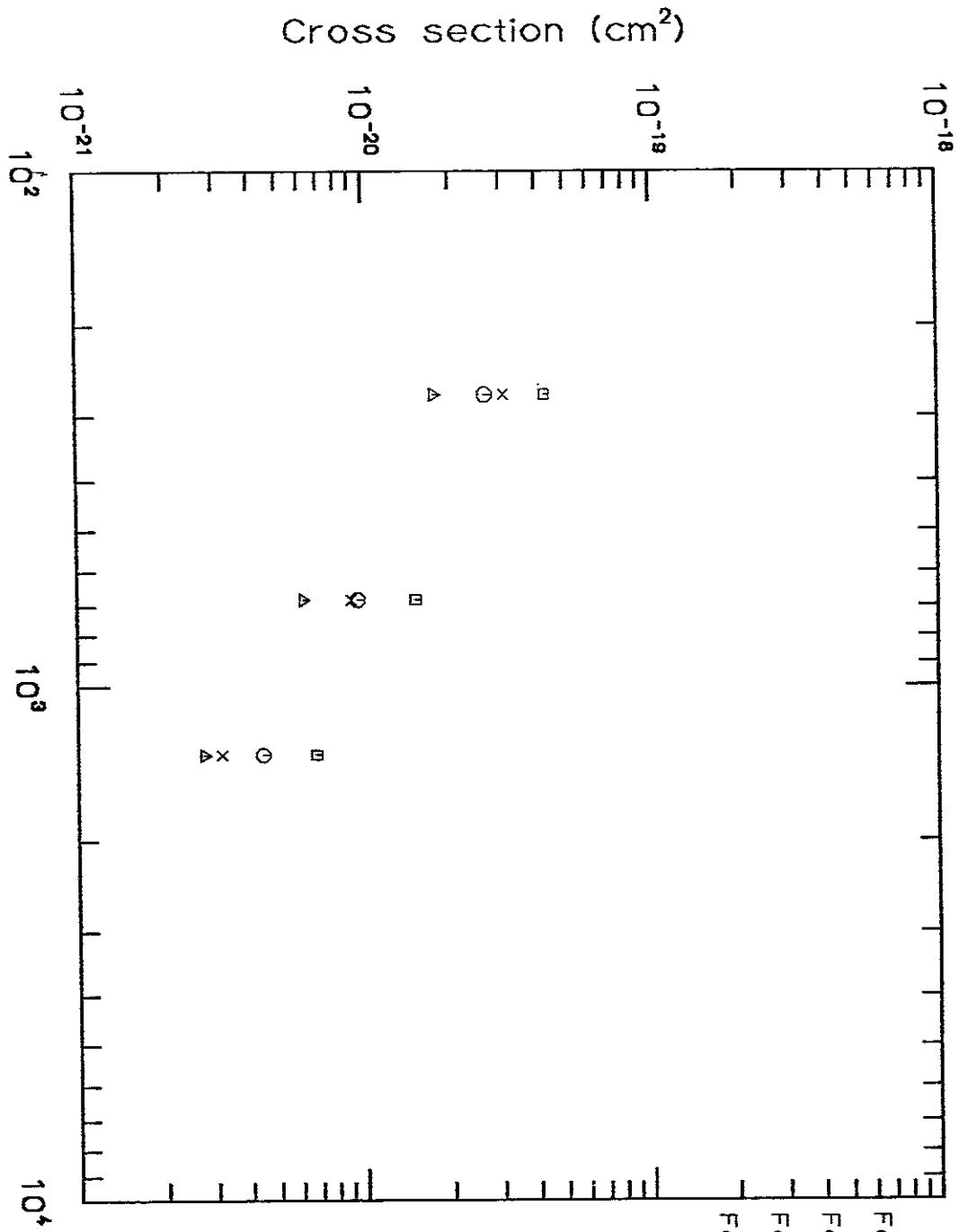








Electron energy (eV)

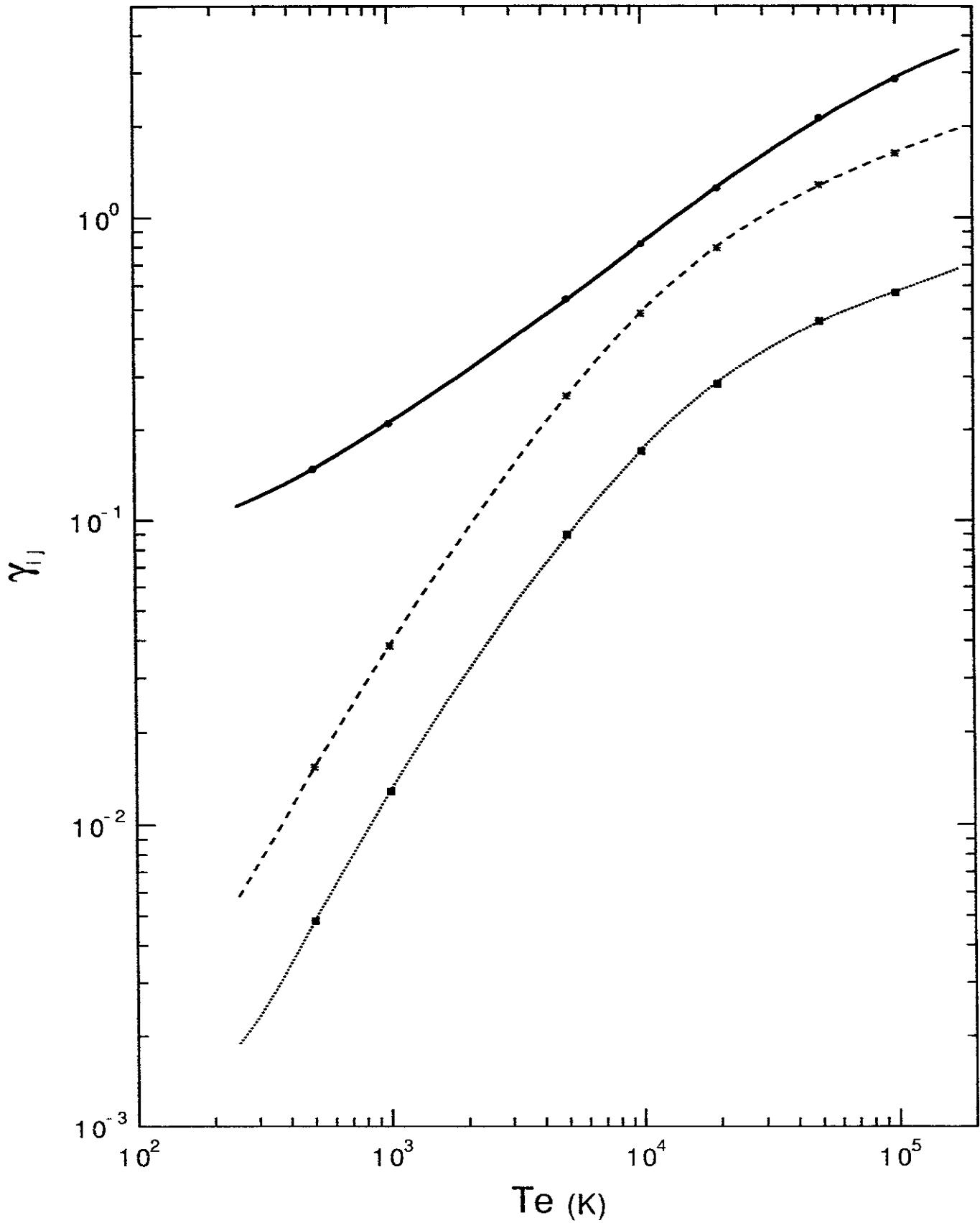


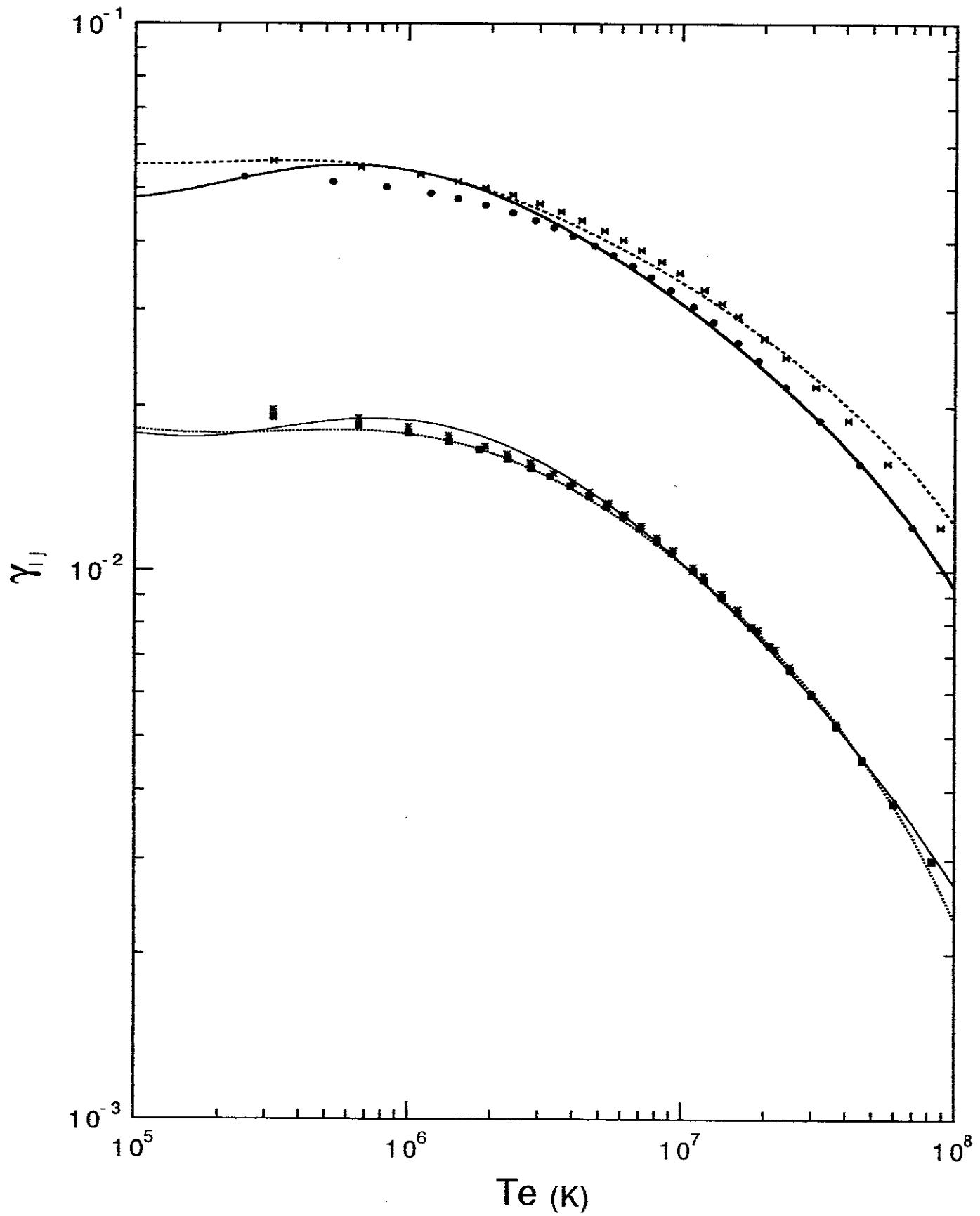
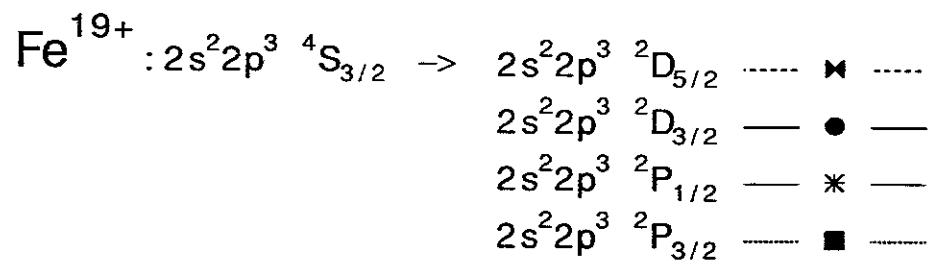
$\text{Fe}^{18+} \{ 2s^2 2p^3 \}^2D_{3/2} \rightarrow \{ 2s^2 2p^3 \}^2P_{1/2}$
 \ominus Bratia, A.K., Mason, H.E., (1980) T
 $\text{Fe}^{18+} \{ 2s^2 2p^3 \}^2D_{5/2} \rightarrow \{ 2s^2 2p^3 \}^2P_{3/2}$
 \times Bratia, A.K., Mason, H.E., (1980) T
 $\text{Fe}^{18+} \{ 2s^2 2p^3 \}^2D_{5/2} \rightarrow \{ 2s^2 2p^3 \}^2P_{1/2}$
 Δ Bratia, A.K., Mason, H.E., (1980) T
 $\text{Fe}^{18+} \{ 2s^2 2p^3 \}^2D_{5/2} \rightarrow \{ 2s^2 2p^3 \}^2P_{3/2}$
 \square Bratia, A.K., Mason, H.E., (1980) T

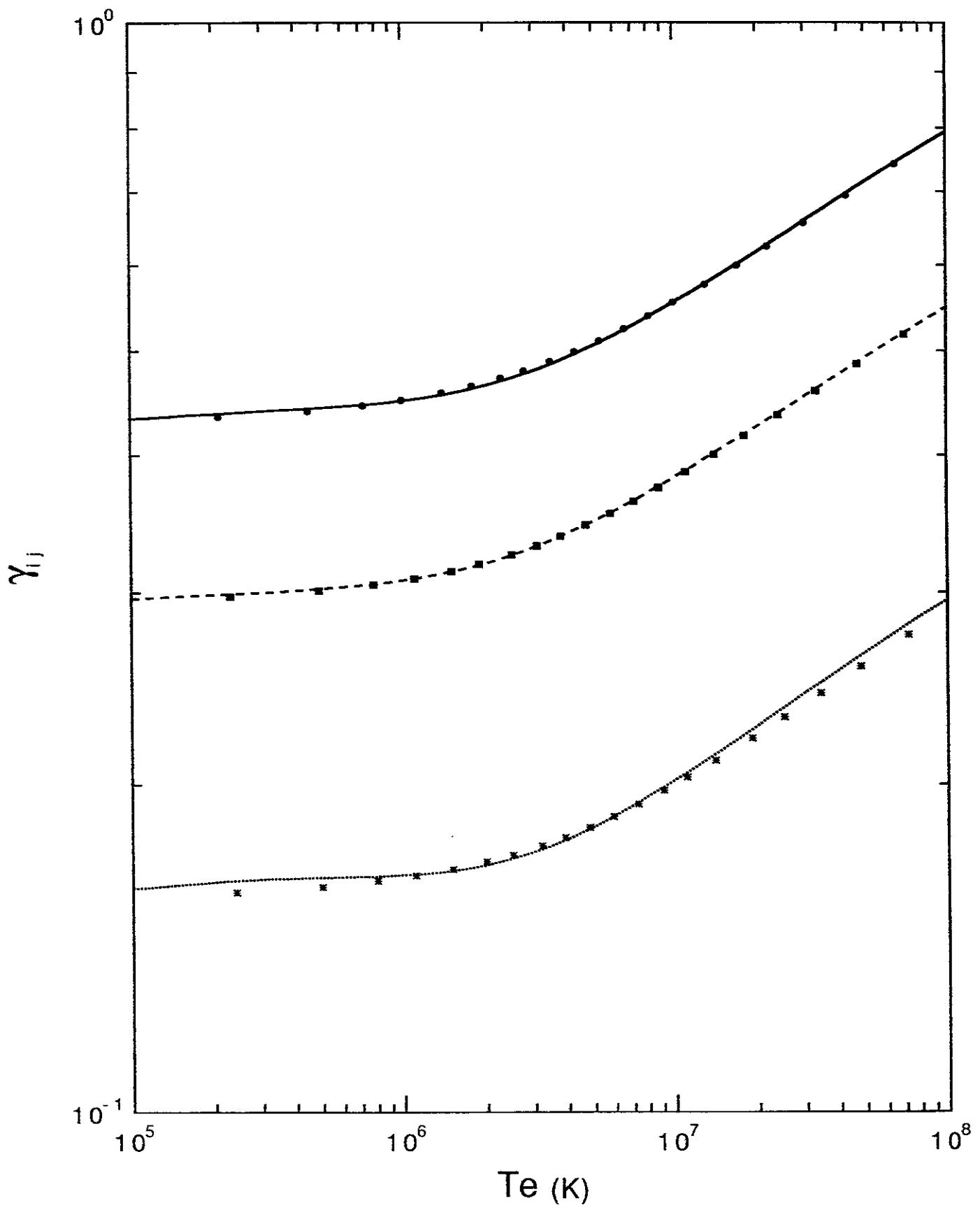
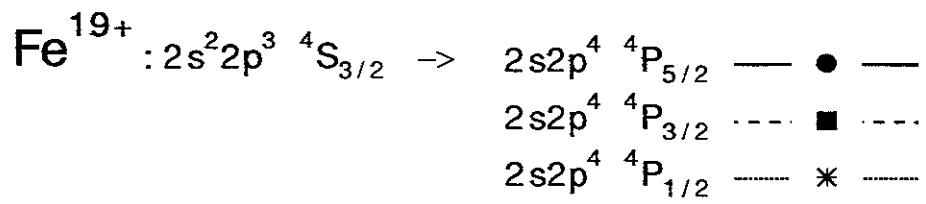
Graphs II.

Recommended effective collision strengths

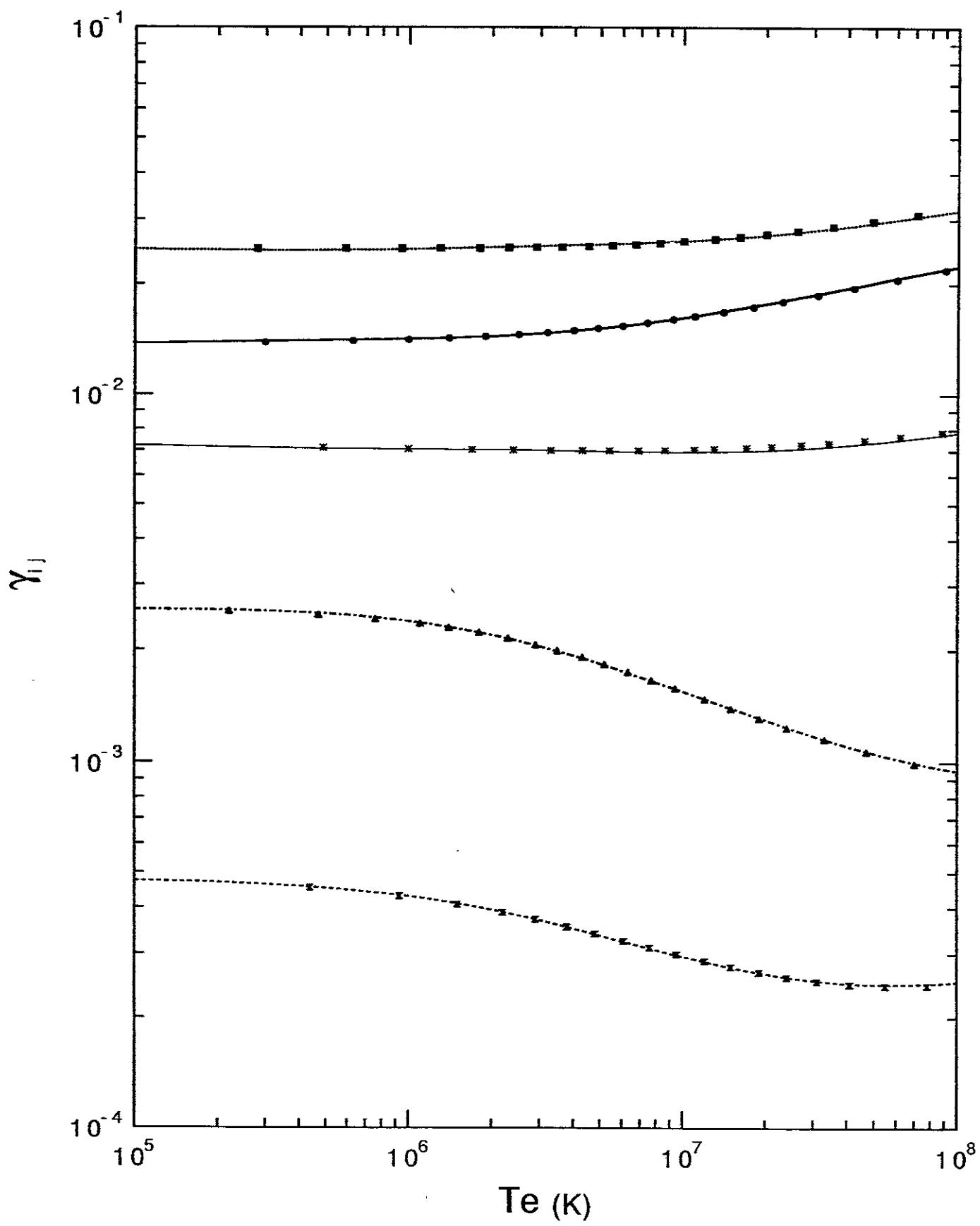
$N : 2s^2 2p^3 \ ^2D^o \rightarrow 2s^2 2p^3 \ ^2P^o$ —●—
 $2s^2 2p^3 \ ^4S^o \rightarrow 2s^2 2p^3 \ ^2D^o$ -·-*—
 $2s^2 2p^3 \ ^4S^o \rightarrow 2s^2 2p^3 \ ^2P^o$ ·——■—

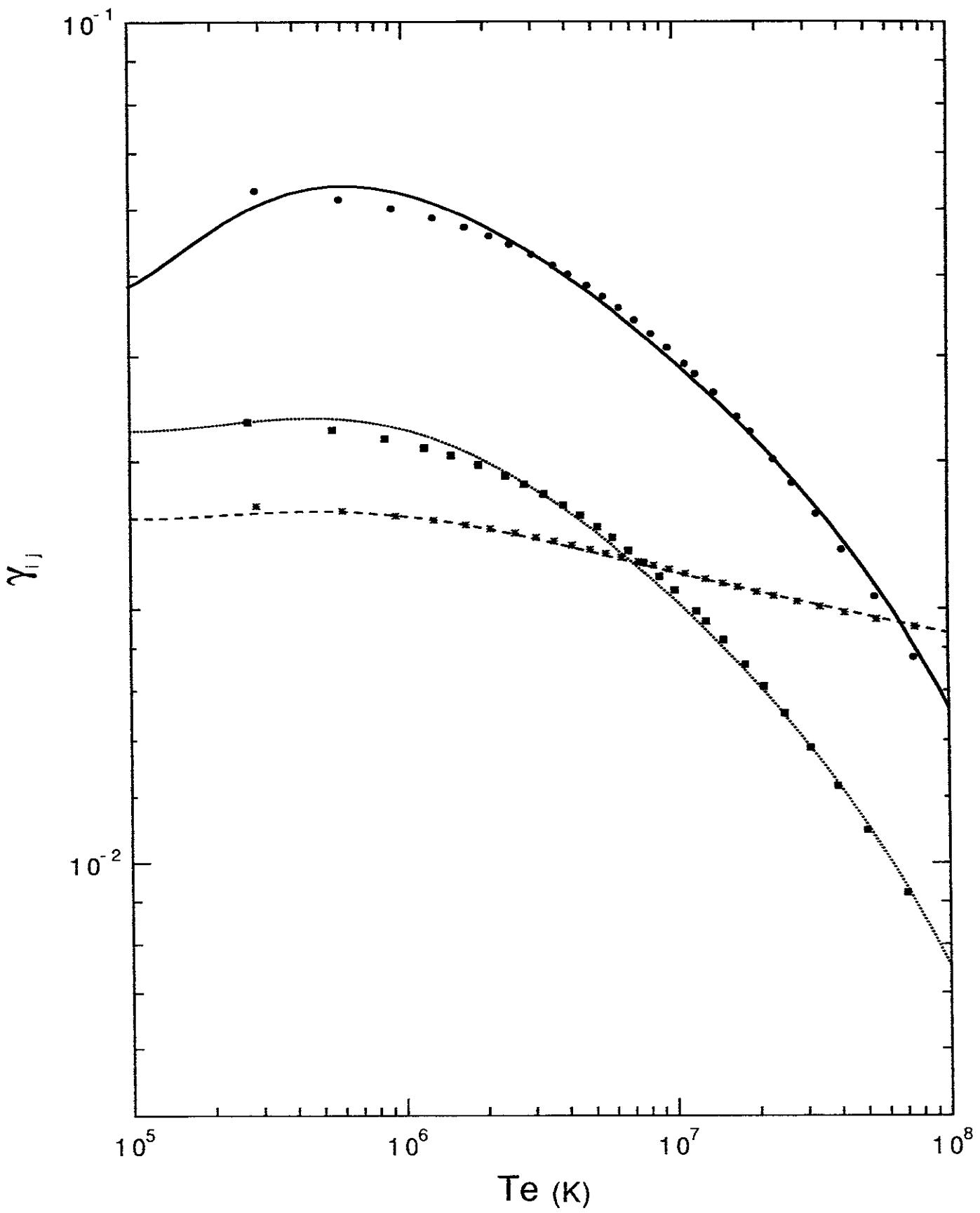
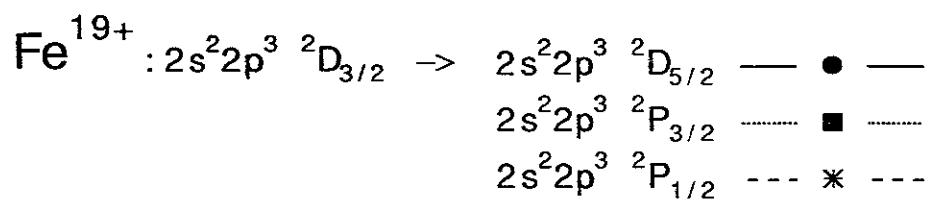


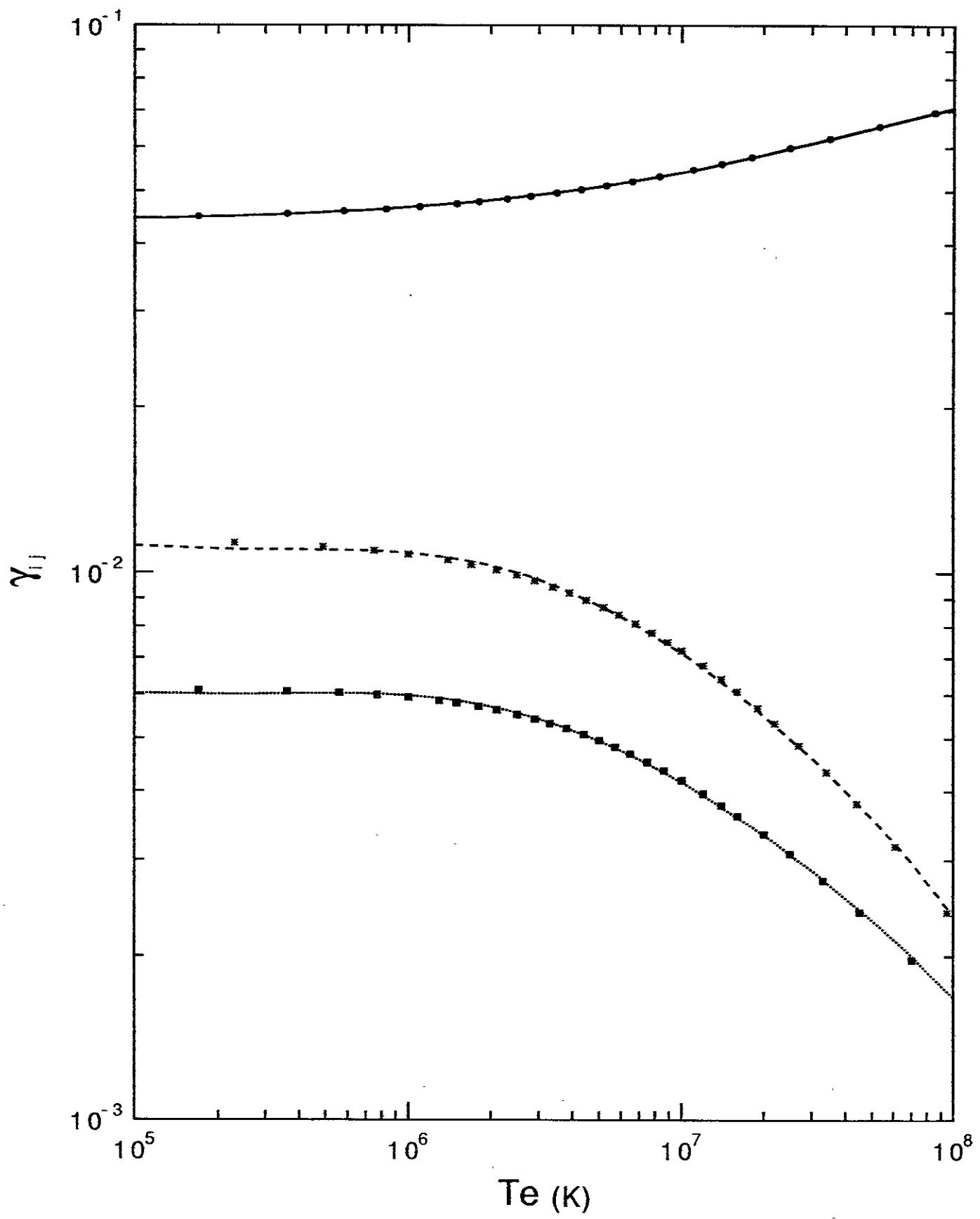
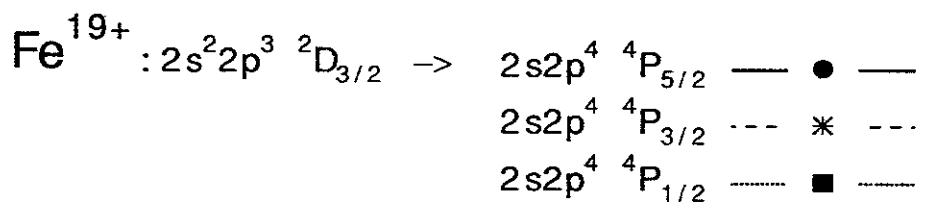




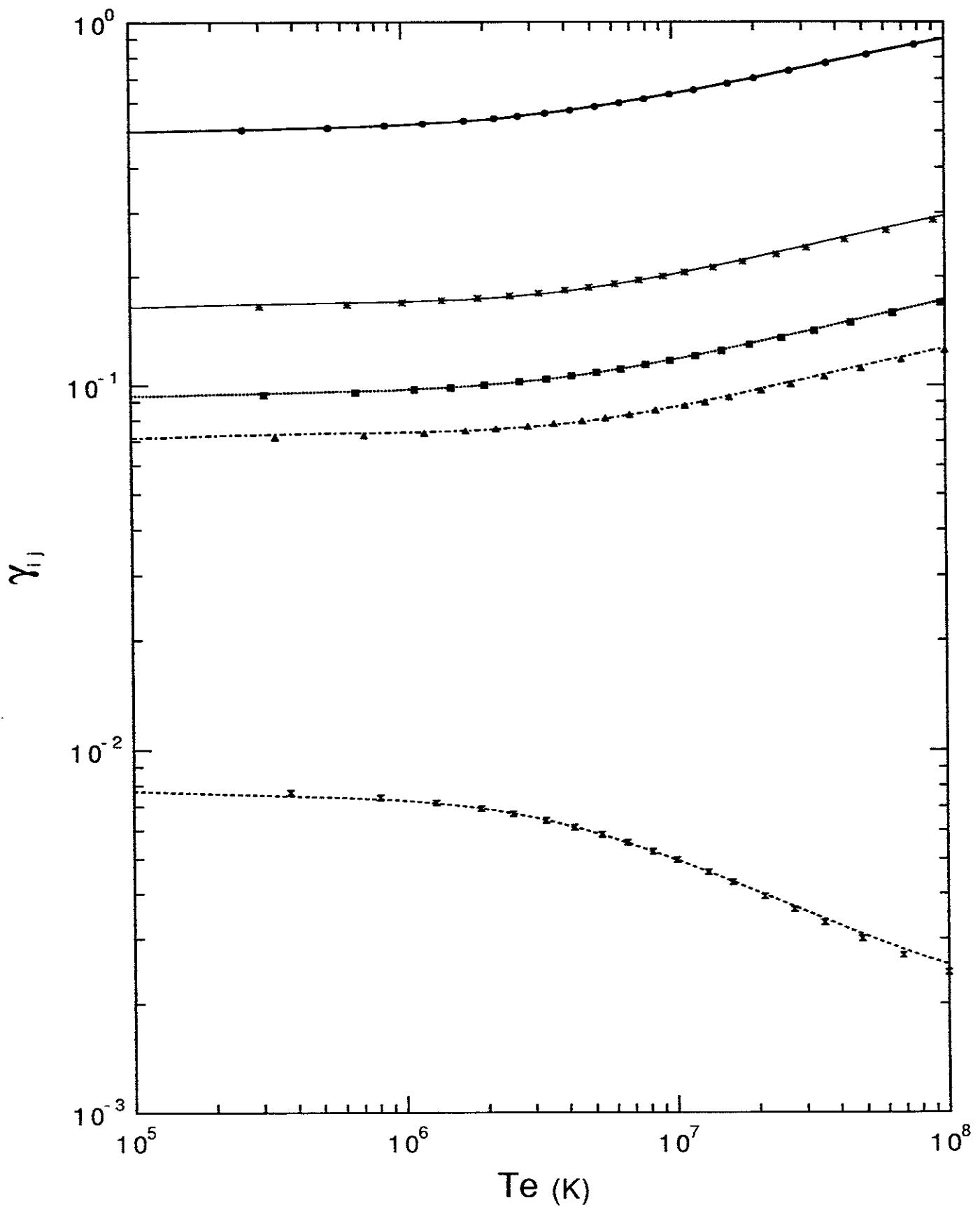
Fe^{19+} : $2s^2 2p^3$ $^4S_{3/2} \rightarrow$ $2s2p^4$ $^2P_{3/2}$ \square
 $2s2p^4$ $^2D_{3/2}$ \bullet
 $2s2p^4$ $^2S_{1/2}$ $*$
 $2s2p^4$ $^2P_{1/2}$ \blacktriangle
 $2s2p^4$ $^2D_{5/2}$ \times

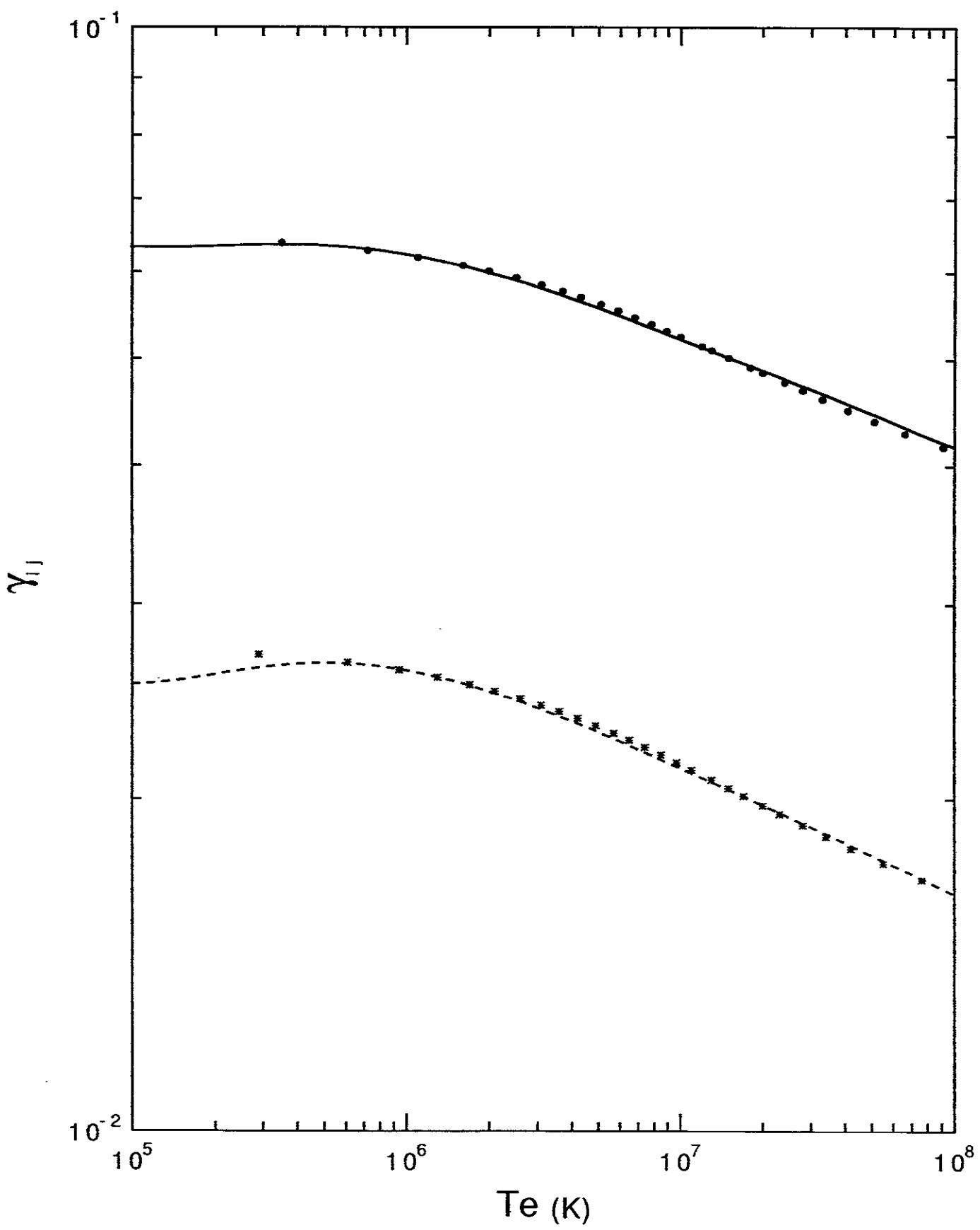
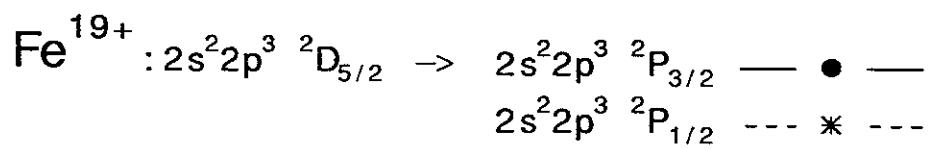


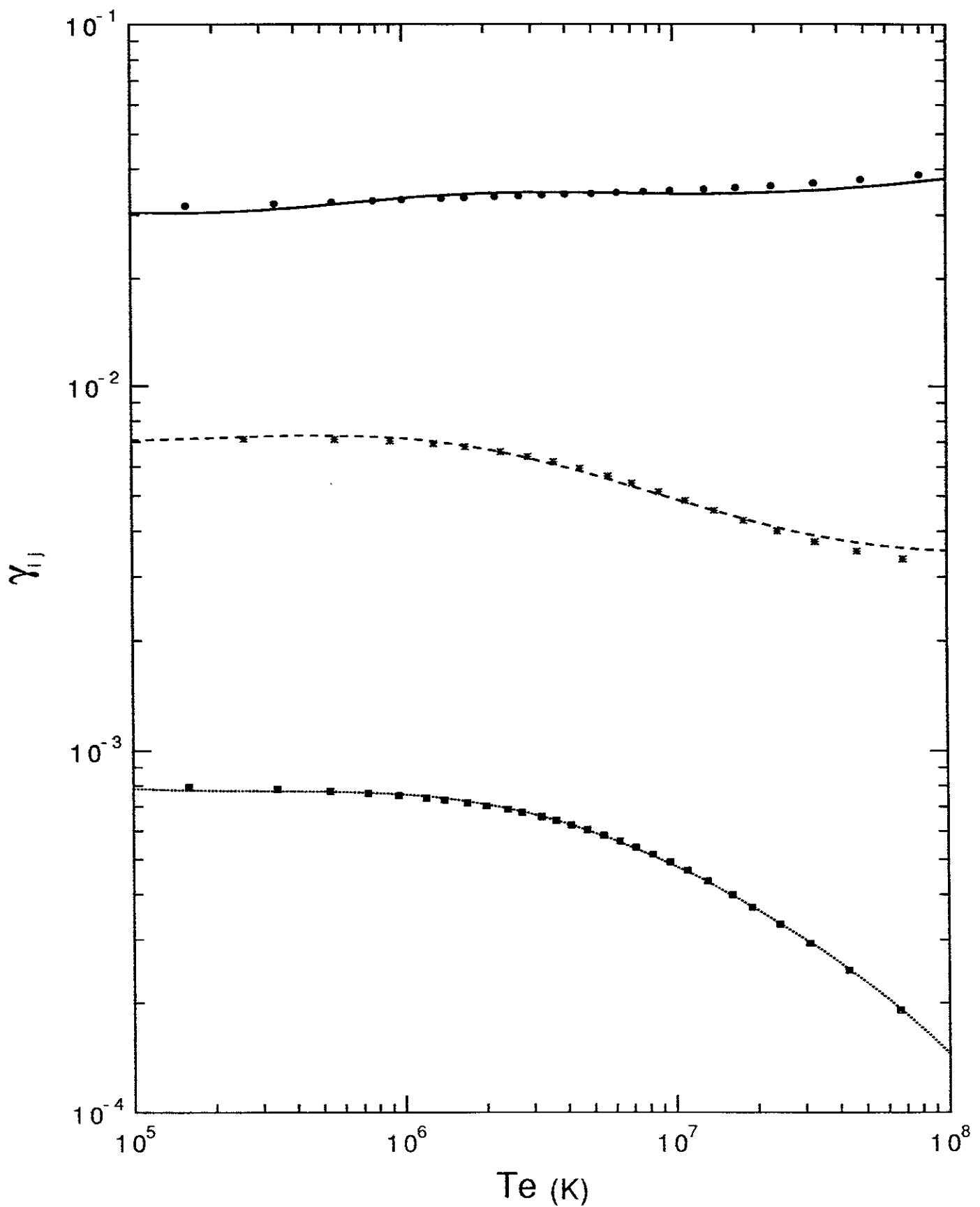
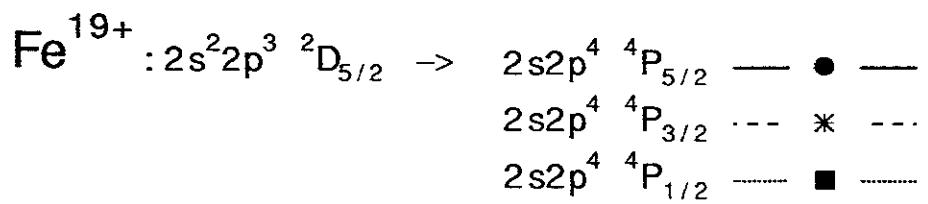




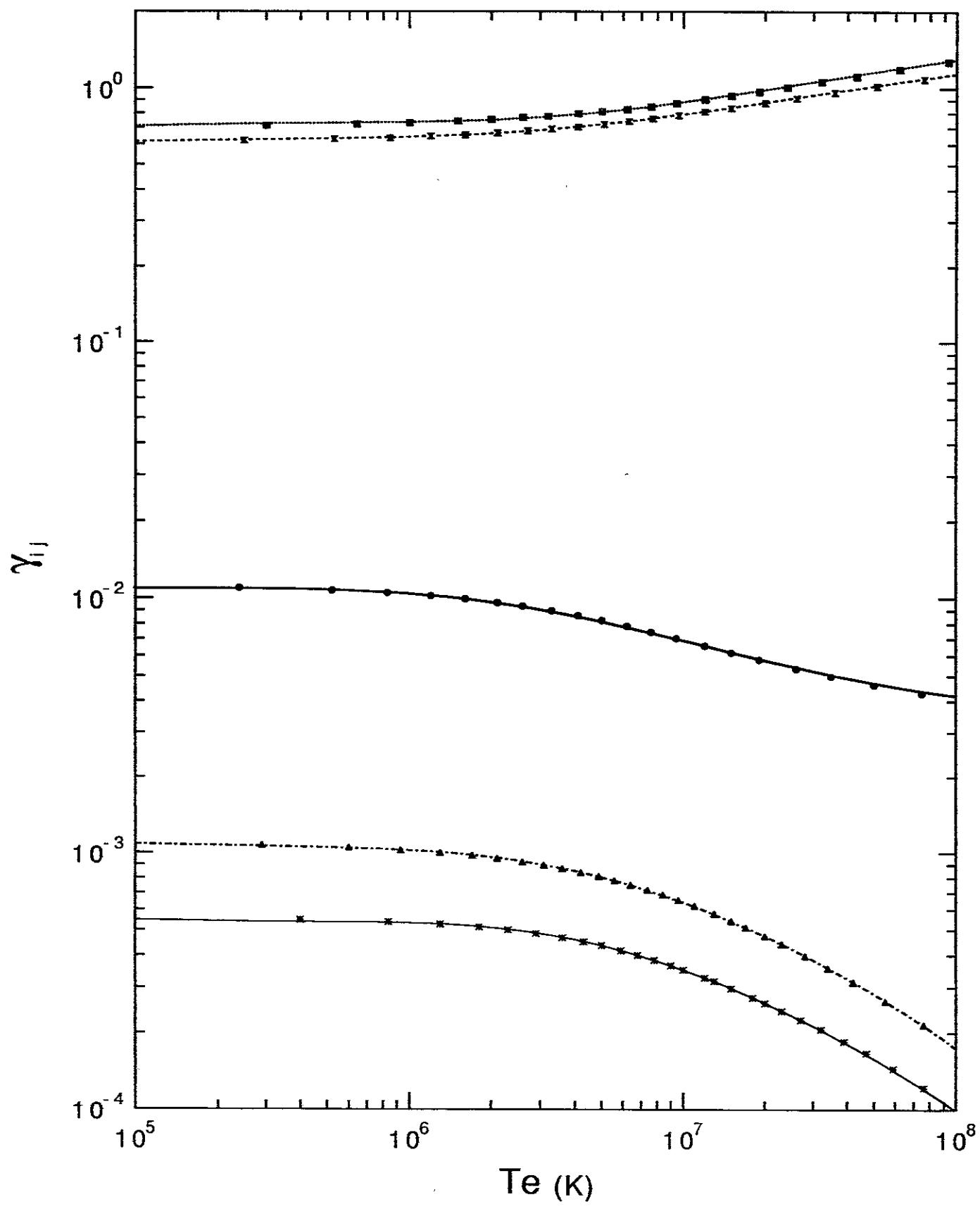
$\text{Fe}^{19+} : 2s^2 2p^3 \ ^2D_{3/2} \rightarrow$ $2s2p^4 \ ^2D_{3/2}$ —●—
 $2s2p^4 \ ^2S_{1/2}$ —*—
 $2s2p^4 \ ^2P_{3/2}$ ■.....
 $2s2p^4 \ ^2P_{1/2}$ ▲.....
 $2s2p^4 \ ^2D_{5/2}$ ×.....

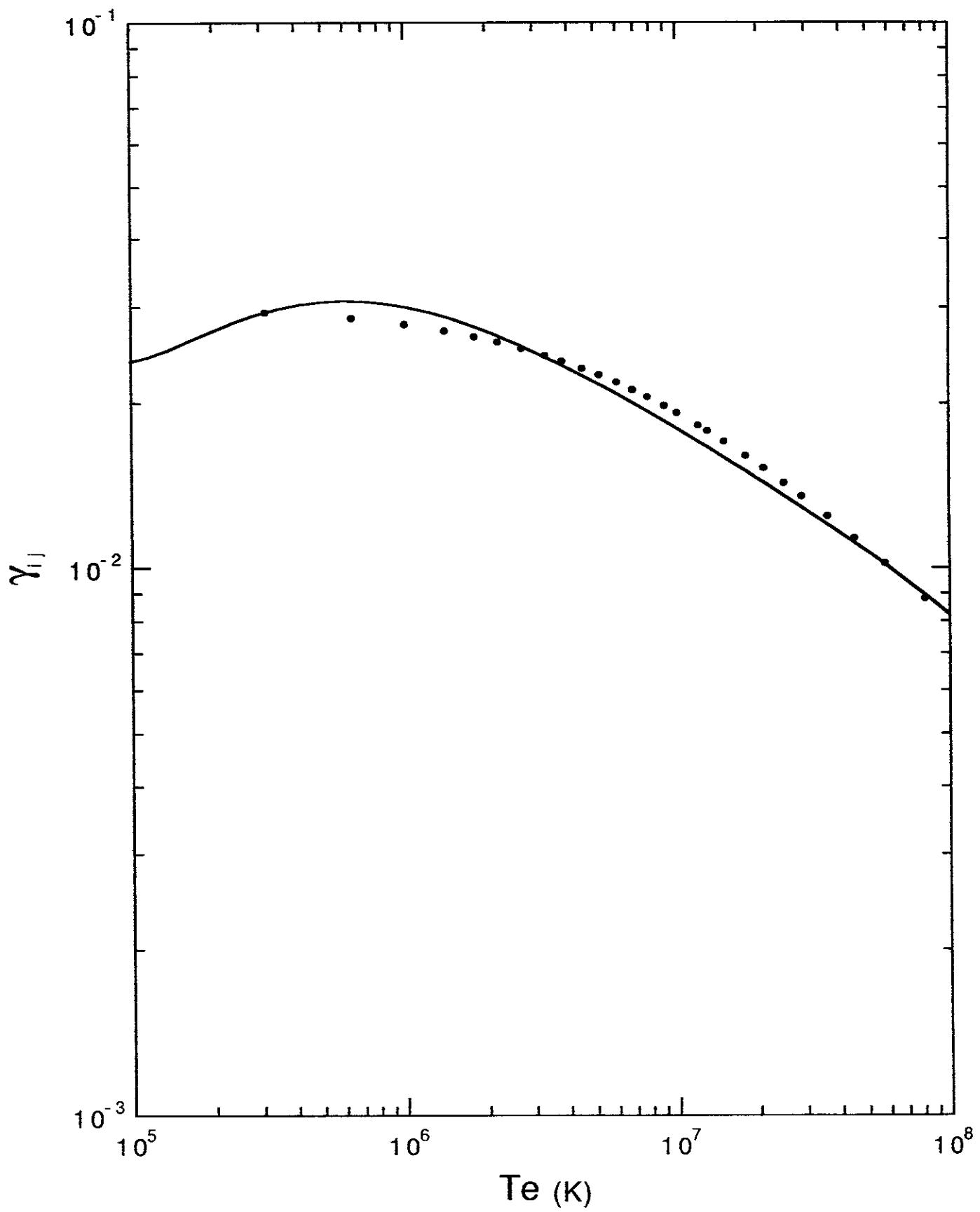
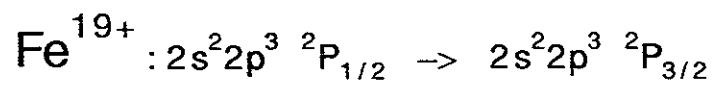


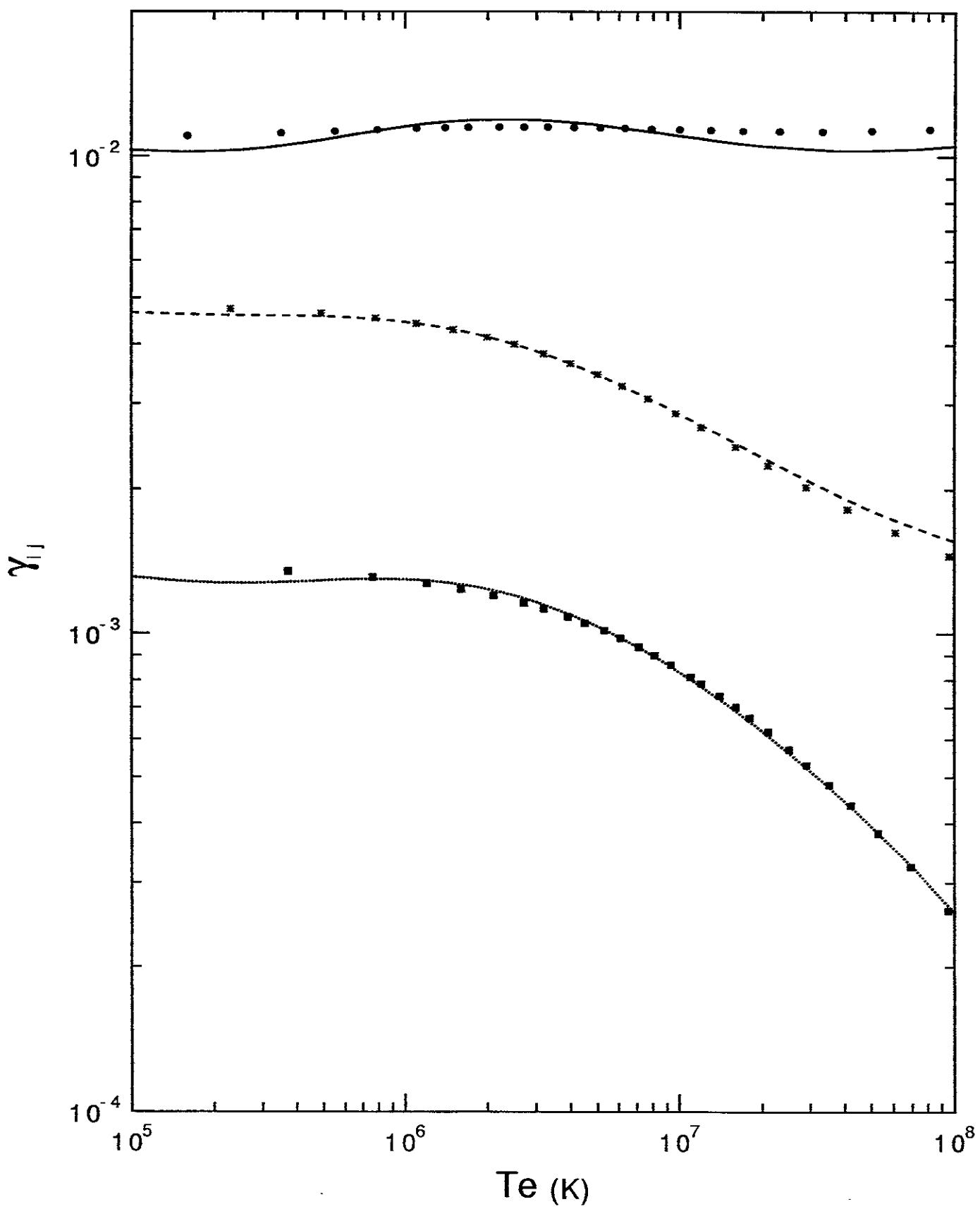
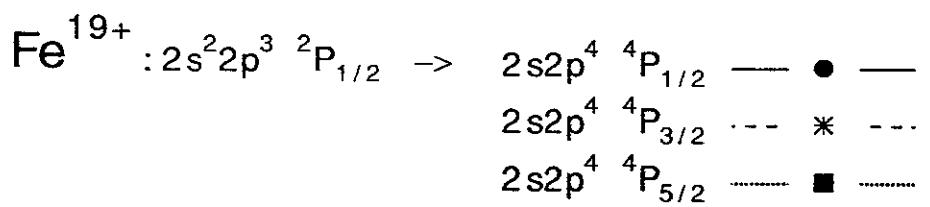




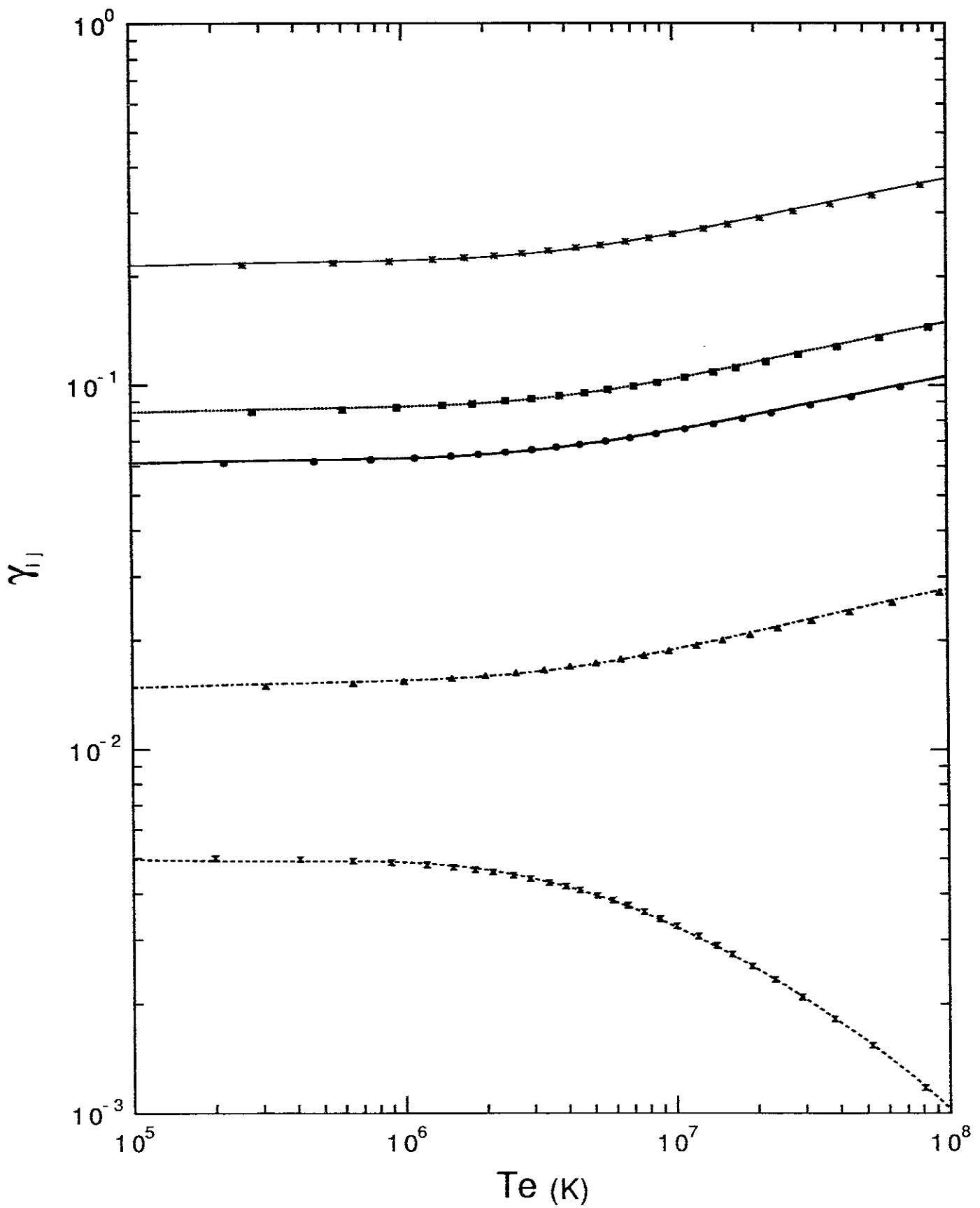
Fe^{19+} : $2s^2 2p^3$ $^2D_{5/2} \rightarrow$ $2s2p^4$ $^2P_{3/2}$ —■—
 $2s2p^4$ $^2D_{5/2}$ ×
 $2s2p^4$ $^2D_{3/2}$ —●—
 $2s2p^4$ $^2P_{1/2}$ ▲
 $2s2p^4$ $^2S_{1/2}$ —*—

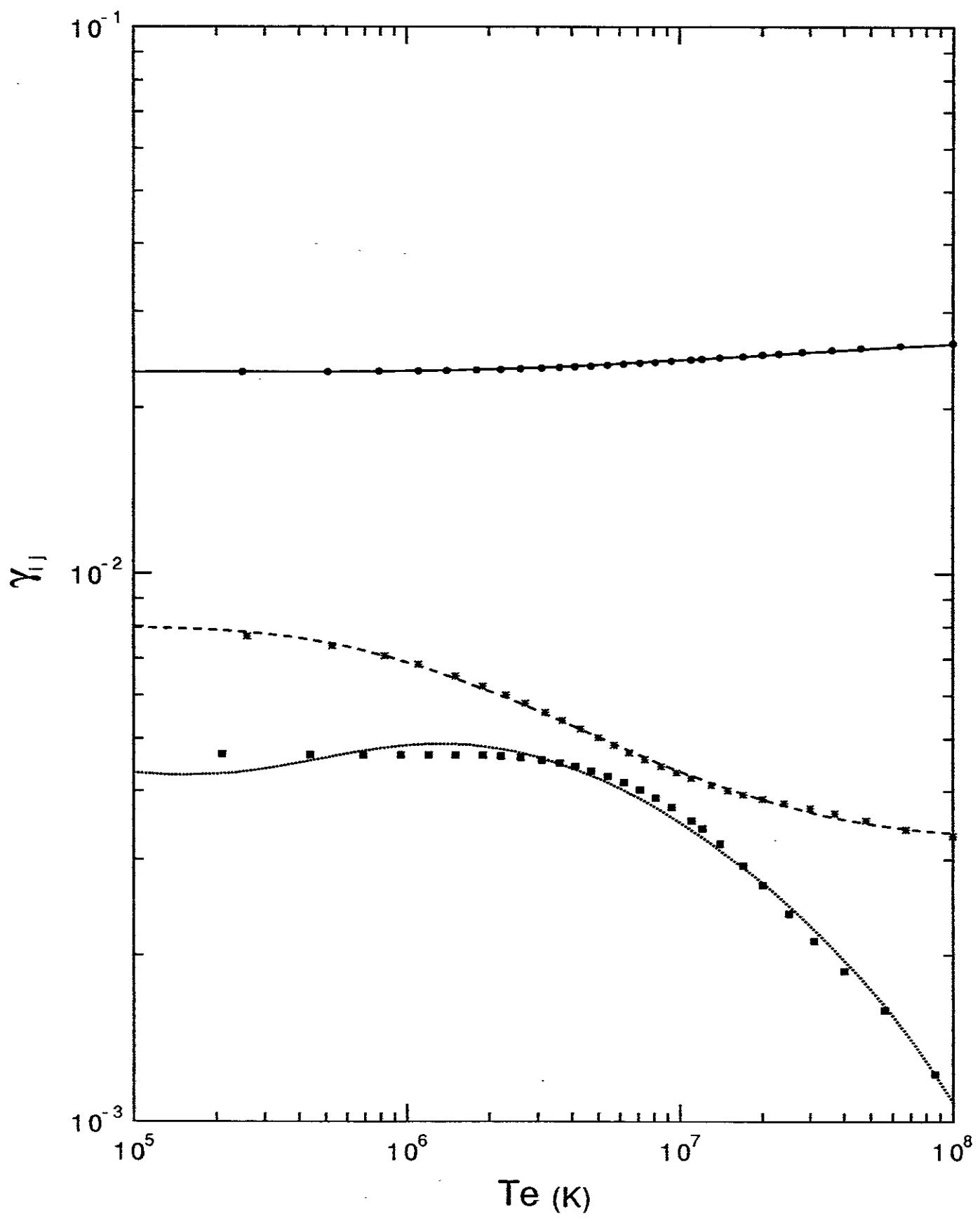
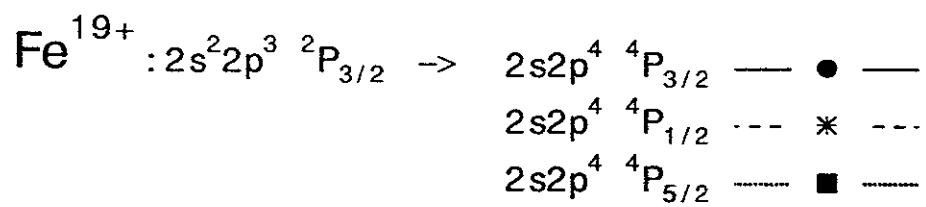




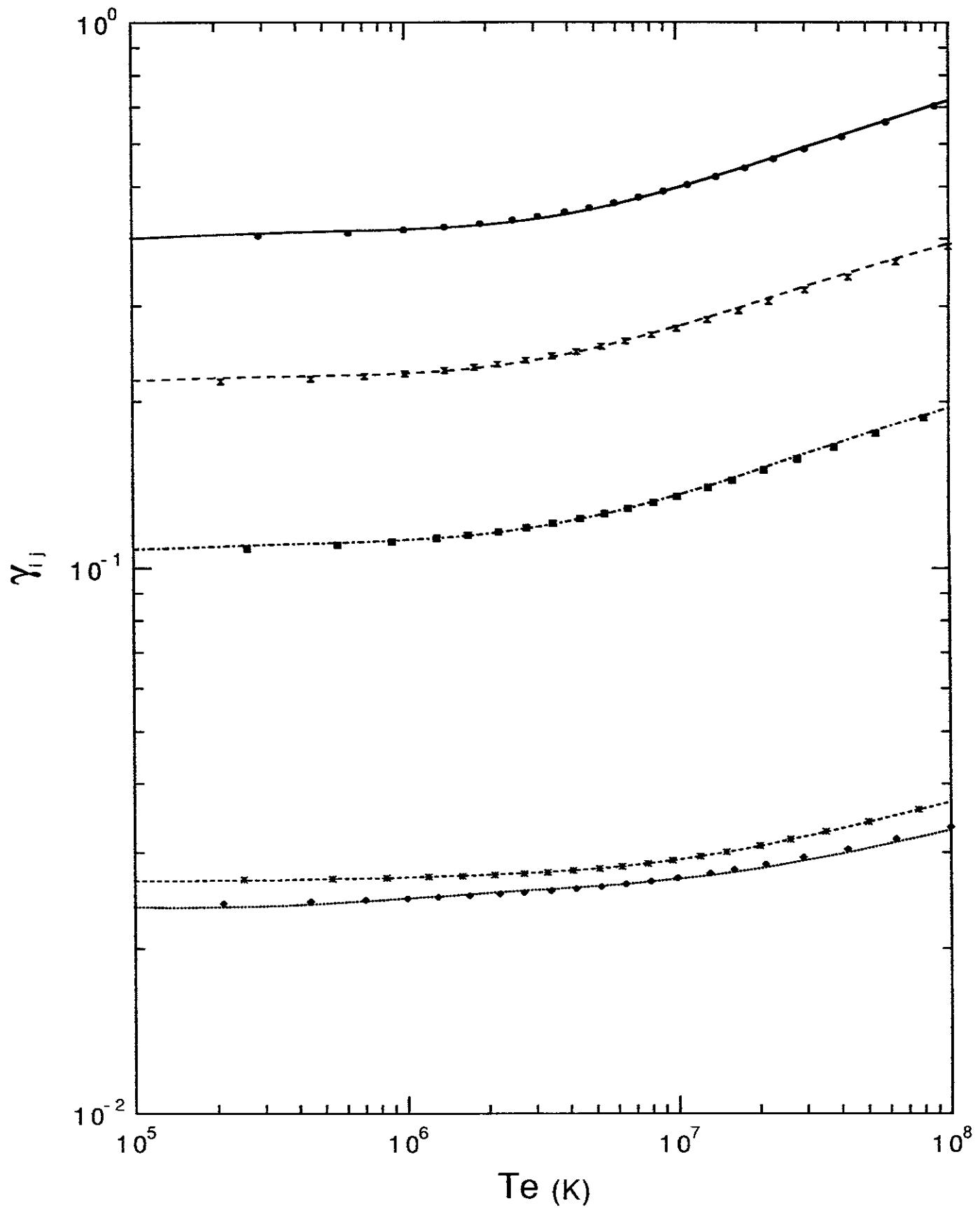


Fe^{19+} : $2s^2 2p^3$ $^2P_{1/2} \rightarrow 2s2p^4$ $^2S_{1/2}$ — * —
 $2s2p^4$ $^2P_{3/2}$ — ■ —
 $2s2p^4$ $^2D_{3/2}$ — ● —
 $2s2p^4$ $^2P_{1/2}$ — ▲ —
 $2s2p^4$ $^2D_{5/2}$ — ✕ —





Fe^{19+} : $2s^2 2p^3$ $^2P_{3/2} \rightarrow 2s 2p^4$ $^2P_{1/2}$ —●—
 $2s 2p^4$ $^2D_{5/2}$ -·-×-·-
 $2s 2p^4$ $^2P_{3/2}$ -·-■-·-
 $2s 2p^4$ $^2S_{1/2}$ -·-*-·-
 $2s 2p^4$ $^2D_{3/2}$ -·-◆-·-



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