§10. Core Electron Temperature Rise Due to Ar Gas-puff in EC-heated LHD Plasmas

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In order to investigate the electron heat transport in LHD plasmas closely, several temperature perturbation techniques, such as a TESPEL injection and a hydrogen ice pellet injection, have been applied. Under some conditions, an abrupt increase in core electron temperature $T_{\rm e}$ has been observed just after those perturbations. The response time of the core $T_{\rm e}$ rise to the edge perturbation is beyond the standard transport paradigm (local and diffusive). One of the candidates for explaining this phenomenon

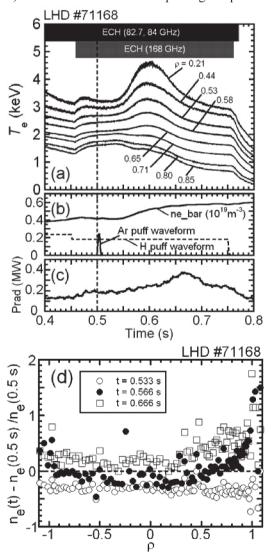


Fig. 1. Time evolution of (a) the $T_{\rm c}$ measured with the ECE radiometer at different normalized minor radii, (b) the line-averaged $n_{\rm e}$ and (c) the total radiated power. (d) Normalized radial profiles of the incremental $n_{\rm e}$ normalized by the base $n_{\rm e}$ (t = 0.5 s) measured with the Thomson scattering diagnostics. In (b), the waveforms of gas-puffs (H and Ar) are also plotted.

is non-locality in turbulence, such as turbulence spreading. In this regard, however, further experimental and theoretical investigations are necessary.

In 10th LHD campaign, after a slight Ar gas-puff, the similar core $T_{\rm e}$ rise has been observed as shown in Fig. 1(a). In this discharge, the plasma was sustained only by ECH (total injected power ~ 1.6 MW around the timing of the Ar gas-puff). The maximum increment of the core $T_{\rm e}$ is about 1.3 keV, which is comparable to the case with the TESPEL injection under the same plasma conditions. The line-averaged electron density $n_{\rm e_bar}$ is increased slightly about 30 ms after the Ar gas-puff. This increase in $n_{\rm e_bar}$ is attributed mainly to the increase in edge $n_{\rm e}$, as indicated in Fig. 1(d). And there is almost no change in the total radiated power before and after the Ar gas-puff. Thus the core $T_{\rm e}$ rise cannot be caused by the RI mode, which has characteristically a density peaking and a significant enhancement of radiated power.

A transient response analysis reveals a similarity of the core $T_{\rm e}$ rise between with the Ar gas-puff and with the TESPEL injection. In both cases, as shown in Fig. 2(a) and (c), the relation of the perturbed electron heat flux normalized by $n_{\rm e}$ to the gradient of perturbed $T_{\rm e}$ shows a similar hysteresis loop. And the maximum of normalized perturbed electron heat flux appears around $\rho \sim 0.5$ (see Fig. 2(b) and (d)). Thus the non-locality in the electron heat transport would be essential for the core $T_{\rm e}$ rise after the Ar gas-puff as well as for that due to the TESPEL injection.

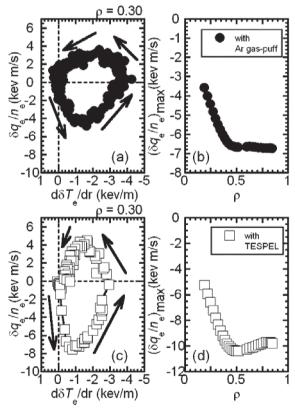


Fig. 2. The relation of the perturbed electron heat flux normalized by $n_{\rm e}$ to the gradient of perturbed $T_{\rm e}$ (left) and normalized radial profile of the maximum of normalized perturbed electron heat flux (right). The data with the Ar gas-puff is depicted in (a, b) and that with the TESPEL injection in (c, d).