§11. Power Balance Analysis in the Hydrogen and Helium Plasma of LHD High Ti Discharge

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The improvement of ion energy confinement was found in He rich plasma of ITB discharge in LHD[1]. Figure 1 shows the comparison of the achieved Te, Ti in the central region $(\rho < 0.3)$ and the deposition power per particle for the different H and He ratio. The set of the plasmas were heated by a 20 – 23 MW hydrogen neutral beam. The line averaged density was $1.35 \times 10^{19} \text{m}^{-3}$ and B_t was 2.75T. The fueling ratio of H and He $(R=n_H/(n_H+n_{He}))$ were controlled by the discharge cleaning and external gas fueling. As shown in Fig.1, T_i in the central region increases with the increase of He concentration, while Te almost stays constant. However, the deposition power per particle was almost constant. Figure 2 shows a comparison of profiles in the lowest and the highest achieved T_i in Fig.1. Due to the larger concentration of He2+ in He rich plasma, the total ion density becomes lower than that of H rich plasma. However, as shown in Fig.1, P_i/n_i is almost constant, thus, the increase of T_i is not due to the difference of P_i/n_i. As shown in Fig.2 (d), the L_{Ti}^{-1} are almost identical at $\rho < 0.9$ in both cases. This suggests the L_{Ti}^{-1} are limited by the ITG threshold [2] in both cases, while T_i in the edge pedestal region is clearly higher in He rich plasma. Higher T_i in the core region of He rich plasma is due to the higher T_i at $\rho > 0.9$ and is due to the relative improvement of thermal diffusivity in the plasma edge region as shown as follows.

Figure 3 (a) shows a comparison of the ion deposition power. With lower R, deposition becomes lower. Figure 3 (b) shows the spatial profile of effective ion thermal diffusivity ($\chi_{i \text{ eff}}$). $\chi_{i \text{ eff}}$ is defined as Q_i/n_i grad T_i , where Q_i is the sum of ion heat flux of H⁺, He²⁺ and C⁶⁺ ions. T_i was measured by CXRS from C⁶⁺ Doppler broadening and assumed to be same for these three ion species. Thus, $\chi_{i eff}$ was defined as a representative values of ion thermal diffusivity. According to the gyro-kinetic non linear simulation, non linear heat fluxes follow gyro-Bohn diffusivity[3], which is proportional to $m^{0.5}T^{1.5}/(aq^2B_t^2)$, where m_i is mass, T is temperature, q is the charge of each charged particle, a is minor radius, Bt is the magnetic field. discrepancy from the gyro-Bohm parameter The dependence is an indicator of improvement or deterioration of confinement. It should be noted that dependence on q is usually not considered in the conventional scaling, however, q shows strong effects in the gyro-Bohm diffusivity and is necessary to be taken into account in He plasma. Here, we defined the gyro-Bohm factor as $F_{GB}=m_{ieff}^{0.5}T_i^{1.5}/(q_{ieff}^2)$, where m_{ieff} is effective ion mass, q_{ieff} is effective ion charge. m_{ieff} and q_{ieff} were estimated from the ratio of the ions. Figure 3 (d) shows $\chi_{i eff}$ profile normalized by F_{GB}. As shown in Fig.3 (d), the difference of the $\chi_{i \text{ eff}}/F_{GB}$ becomes smaller. In particular, for R = 0.34 - 0.73, $\chi_{i \text{ eff}}/F_{GB}$ are almost identical at $\rho < 0.8$, while the difference is still large at $\rho > 0.8$. This result causes the higher edge pedestal T_i as shown in Fig.2 (d).

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- 2) Murakami, A., et al., Nucl. Fusion 46, S425 (2006).
- M. Nunami et al, Plasma Fusion Res. 6, 1403001 (2011)
- 4) Pusztai, I., et al, Phys. Plasma 18, 122501(2011)



Fig.1 Comparison of T_e , T_i at $\rho < 0.3$ and deposition power per particle for different He and H ratio. Deposition power is calculate by GNET[2]



Fig.2 (a) n_e (plain lines), n_i (dashed lines), (b) T_e , (c), T_i , (d) expanded view of T_i profiles (vertical axis is logarithmic scale) in He rich (R=0.34 and the highest T_i plasma) and H rich (R=0.78 and the lowest T_i discharge) plasmas.



Fig.3 (a) Ion deposition power, (b) ion kinetic energy, (c) χ_i $_{eff}$ and (d) normalized χ_i $_{eff}$